Interference in Cellular Networks: Sum of Lognormals Modeling

by

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Abstract

We examine the existing methods for evaluating the distribution of the sum of lognormal random variables, focusing on closed-form results. We find that there are no results in literature that are both simple and accurate. We then derive a new closed-form expression for the lower tail of the distribution, and use it to construct a new method using a power-lognormal distribution. We apply both basic moment-matching and our new method the problem of the total interference power in a cellular system. For both methods, we derive equations that find the interference distribution essentially in closed form, using minimal numerical integration. We apply both methods to the uplink and downlink in systems with and without power control, for various cellular layouts, channel models and user activity probability. We compare distributions obtained by Monte-Carlo simulation directly with those obtained by our method, and find very good matches in many useful cases.

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– Soli Deo Gloria –

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List of Acronyms

AG Arithmetic-Geometric mean inequality method [46]

blf Best lognormal fit

BLN Bernoulli (times) Lognormal

BPLN Bernoulli (times) Power Lognormal

BR Beaulieu – Rajwani (distribution function) [27]

BS Base Station

BX Beaulieu – Xie (method for analysing the SLN problem) [35]

cdf Cumulative distribution function
 CDMA Code-Division Multiple Access
 cf Characteristic Function (of a RV)

DL Downlink

DMU Desired Mobile User (MU under analysis)

FH-SS Frequency Hopping (Spread Spectrum technique)

FS Fixed Station (i.e., a BS or RS)

FW Fenton – Wilkinson (method for analysing the SLN problem) [1]

h.t. Heavy-Tailed (property of distributions)

i.i.d. Independent, identically distributed (random variables)

i.n.i.d. Independent, but not necessarily identically distributed (random variables)

IFS Interfering Fixed Station

IMU Interfering Mobile User (MU whose link causes interference to the DMU)

LL Le-Ngoc – Lam (distribution function) [27]

LN Lognormal

LSG Log-Shifted Gamma (Distribution function used in the LL method) [27]

MU Mobile User/Unit

OFDMA Orthogonal Frequency-Division Multiple Access

PC Power Control

pdf Probability density function

PLN Power Lognormal
RS (fixed) Relay Station
RV Random Variable

s.e. Sub-Exponential (property of distributions)

SINR Signal to Interference and Noise (Power) Ratio

SIR Signal to Interference (Power) Ratio

SLN Sum of Lognormal (Random Variables)

SY Schwartz and Yeh (method for analysing the SLN problem) [37],[47]

UL Uplink

WMZ Wu-Mehta-Zhang (method for analysing the SLN problem) [35]

ZD Zhao – Ding (distribution function) [28]

ZS Zhang– Song (distribution function) [28]

List of Symbols

#{ <i>S</i> }	Cardinality of set S
a_I	μ -parameter of the BLN approximation to $F_I(x)$
$A_n, A_{DL}, \widetilde{A}_n$	A-coefficients describe the pathloss statistics for a given cellular layout.
b_I	σ -parameter of the BLN approximation to $F_I(x)$
B_n	$n^{ ext{th}}$ statistical moment of any Q_i
\widetilde{B}_n	$n^{ m th}$ statistical moment of any Q_i/L_i
C_0	Cell of interest
C_1	A nearest interfering cell
C_i	Interfering cell
$E[\mathcal{X}]$	Expected value of RV \mathcal{X}
$F_{\mathcal{X}}(x)$	cdf of RV ${\mathscr X}$
$f_{\mathcal{X}}(x)$	pdf of RV \mathcal{X}
$\tilde{F}(x)$	F(x) transformed under T
I	RV: Total received interference power
i, j	Index of interfering cell or lognormal summand
I_i	RV: Received interference power coming from cell C_i
K	Rician fading K -factor in linear scale: 0 – Rayleigh fading, ∞ – no fading
k_{DL}	Downlink option indicator: 0 – UL, 1 – DL
k_{PC}	Power control option indicator: 0 – no PC, 1 – PC
l	$+\infty$ or $-\infty$: indicates the upper and lower tail respectively on lognormal paper
L_i	RV: Pathloss between FS_i and MU_i
m	μ -parameter of the PLN distribution
m_I	μ -parameter of the BPLN approximation to $F_I(x)$
N	Number of interfering cells or number of terms in a sum of lognormals
n	$n^{\rm th}$ statistical moment
p	Probability of activity of an interferer from any given cell
$P[\mathcal{E}]$	Probability of event \mathcal{E}
P_i	RV: Normalised transmit power of interferer from cell C_i due to power control;

	without power control: constant = 1
p_{χ}	Probability of activity of at least one interferer
Q_i	RV: Product of all factors affecting the received power from cell C_i , except pathloss
R_i	RV: Inverse pathloss suffered by interfering signal from cell C_i
r_i	RV: Physical (normalised) distance from interferer in cell C_i
S	σ -parameter of the PLN distribution
S_I	σ -parameter of the BPLN approximation to $F_I(x)$
T	Transformation (mapping) from linear to lognormal probability paper
t	Power parameter of the PLN distribution
t_I	Power parameter of the BPLN approximation to $F_I(x)$
X	RV: sum of lognormal RV's Y_i
$x_i, y_i,$	Cartesian coordinates of FS_i (equivalently, the centre of cell C_i)
Y_i	RV: lognormal, with parameters μ_i , σ_i
Z_I	Auxiliary term in calculating the <i>BPLN</i> parameters to approximate $F_I(x)$
α_i	RV: Instantaneous fading envelope amplitude in linear scale of signal from cell C_i
β	Pathloss exponent
$\gamma(x,\alpha)$	Incomplete (lower) gamma function
$\Gamma(x)$	Gamma function
$\Phi(x)$	Standard normal <i>cdf</i>
ζ_i	RV: Lognormal-distributed shadowing in linear scale suffered by interfering signal
	from cell C_i
κ_i	Pathloss statistics for received power from cell C_i
$\widetilde{\mathcal{K}}_i$	Pathloss statistics for received power from cell C_i in the case of uplink, power control
\mathcal{K}_L	Constant factor multiplying the blf to the lower tail of the SLN distribution
κ_U	Constant factor multiplying the blf to the upper tail of the SLN distribution
λ	Constant = $\ln 10/10 \approx 0.23$, used to convert between dB and natural units
$\Lambda(x, y)$	Integral used to find the moments of the PLN distribution
μ_i	μ -parameter of lognormal RV Y_i
μ_L	μ -parameter of the blf to the lower tail of the SLN distribution
μ_U	μ -parameter of <i>blf</i> to the upper tail of the SLN distribution
ρ	Constant = 0.915, radius of circle that approximates well a unit hexagon

$ ho_{ij}$	Correlation coefficient between two Gaussian RV's with indices <i>i</i> , <i>j</i>
σ	Shadowing variance in natural units
σ_e	Power control error lognormal variance in natural units
σ_i	σ -parameter of lognormal RV Y_i
σ_L	σ -parameter of the blf to the lower tail of the SLN distribution
σ_U	σ -parameter of <i>blf</i> to the upper tail of the SLN distribution
χi	RV: Bernoulli-distributed indicator of activity of interferer from cell C_i

Many other symbols are used locally within sub-sections, sometimes overriding the above definitions. Hopefully, this will always be clear from context.

Chapter 1

Introduction: Total Interference in a Cellular Network

1.1 Motivation for this Work

While the problem of analysis of interference in cellular systems is at least as old as the first systems themselves, the topic still attracts much interest. This is essentially because the problem is not really solved yet, though there exist many methods that solve particular problems, with varying levels of complexity (numerical or closed-form), and various levels of accuracy. There does not seem to be a well-established systematic method to analyse this class of problems, and several papers [1]-[16] work on very similar problems without there being much incremental knowledge. Indeed, these papers have a substantial portion of overlap, particularly when it concerns two essential mathematical problems: calculating the effect of the mobile user's (MU) random distribution in space, and evaluating the sum of lognormal (SLN) random variables (RV). In this thesis we propose to make a synthesis of the common points of the work in these papers. We do this first by creating a mathematical and simulation model that is sufficiently detailed to encompass many particular problems, but simple enough for analytical tractability. We then evaluate the existing methods for calculating the SLN distribution, and develop our own method based on a synthesis of several known results. Finally we apply this method to the interference distribution problem, and obtain simple expressions to calculate an approximation to this distribution as a function of the parameters of a particular problem. In particular, we can separate the geometry of the cellular layout (equivalently, the user distribution) from the rest of the problem, which means that some geometrical parameters can be calculated once and for all and tabulated for future use. We believe that the method we propose is flexible enough to

accommodate many classical problems while retaining its simplicity, and also can be modified to solve new problems in emerging wireless systems.

The main underlying mathematical problem is the SLN distribution, which appears not only in wireless communications, but also in electronics [17], physics [18], optics [19], economics [20], and is of interest to statistical mathematicians [21], [22]. This versatility makes the problem even more important to solve.

Knowing the distribution of the interference in a system is the next best thing to actually knowing the interference at any given moment. Indeed, unless one uses adaptive channel mechanisms [1], [2] or cognitive radio techniques, one has to allocate scarce radio resources (spectrum, power) based on statistical knowledge. A precise knowledge of the interference cumulative distribution function (*cdf*) is necessary for optimal design: if the interference is overestimated, radio resources are wasted in trying to avoid outage situations that almost never happen, while if it is underestimated, the quality of service suffers from too frequent outages. In past years, not only cellular, but many other wireless technologies have been emerging (WiFi, WiMax, Bluetooth, UWB, ...), and much research is currently being done in ad-hoc, mesh and sensor networks. All these often use shared, unlicensed bands. As so many wireless technologies come to share the same spectrum, the systems become interference-limited. While it is not the purpose of this thesis to explore the interference between such systems, simply understanding the statistical behaviour of interference becomes more important than ever.

1.1.1 Some Previous Analytical Work on Interference

In this section we touch on two previous papers that analyse interference, we show why they are important and also what is lacking in them. It is these lacks that we want to fill in this thesis. We borrow the mathematical models in order to formulate the problem in Section 1.2.

Reference [3] is the work that is closest to ours and has inspired many of the points in this thesis. First of all, it is one of the rare papers that introduce the concept of geometrical coefficients, integrals that can be calculated independently of other system considerations, and tabulated for reference. Thus a particular cellular layout may be described by a collection of numbers. These can be substituted into closed-form expressions in order to obtain the interference statistics in a particular cellular system. A weakness of this paper is that many of the stochastic quantities have been replaced by their expected value, which is an over-simplification.

Reference [1] uses a numerical method by Schwartz and Yeh (SY) (see Section 2.3.1.1). The disadvantage of this paper is its heavy use of numerical integration at all stages of the process. We intend to show in our work that much of it is unnecessary, and that a tractable solution exists.

Both papers study interference in the uplink (UL). In [3], the downlink (DL) is also examined, but there is little analysis done and the method becomes mostly numerical. We intend to show that we may find closed-form expressions not only for the UL, but for the DL as well. The DL, though more challenging analytically, must be well understood because of the heavy traffic asymmetry leaning towards the DL in future systems [23]. The greatest weakness of both papers is the lack of a direct comparison between simulation results and the theoretically derived distributions. In fact, it is rare to find a comparison between a simulated and calculated distribution curve of the interference statistics. We will always support our theoretical results with Monte-Carlo simulations in this thesis.

1.1.2 Purpose and Motivation for a Closed-Form Solution

Reference [24] suggests that knowing the statistics of the interference coming from just one interferer is not enough, and that knowledge of the sum interference power is necessary for accurate interference modeling. We consider a cellular system with known layout, a propagation

model, user behaviour statistics, and potentially a power control (PC) mechanism. We can simulate interference to a particular desired mobile user (DMU) in such a system according to the model described in Section 1.2, using Monte-Carlo Simulation. This thesis has one main goal: to reproduce the Monte-Carlo interference *cdf* as closely as possible using, as much as possible, only analytical, closed-form expressions, with minimal numerical integration. This is useful for the following reasons:

- Monte-Carlo simulations can be very time-consuming, while numerical integrations, if well-behaved, are substantially faster. Also, analytical expressions programmed into a computer execute almost instantaneously.
- 2. A simple closed-form solution to a problem is always desirable: we try to aim for expressions that can be evaluated using a scientific calculator and that could be included in textbooks on basic wireless communications. Such a method is much more likely to gain popularity than one which takes a significant time to learn and implement.
- 3. A closed-form method to a simple problem forms a basis for further analysis of more complex ones, such as those in emerging next-generation wireless systems.
- 4. Closed-from expressions are useful in order to understand the exact effect of every system parameter on the final performance of the system. We can for example take derivatives with respect to the parameters and see how sensitive the system is to them.

1.1.3 Goals and Philosophy of this Work

Throughout our work, we have kept in mind the following criteria that our method should fulfill:

- 1. As much as possible, obtain closed-form expressions.
- 2. As much as possible, use only standard functions found on a scientific calculator.

- Obtain analytical cdf curves that are close (within 1dB) to the simulated interference cdf
 at all points of the curve.
- 4. As much as possible, simplify the expressions: if removing a term does not increase the error between the analytical and simulated curves by a significant percentage, then that term should be omitted.
- 5. As much as possible, avoid numerical methods.
- When numerical integration is necessary, let the integral be separable as much as possible from the closed-form part.
- 7. When numerical integration is necessary, let it be easily converging using Riemann sums with a moderate number of steps.
- 8. As much as possible, let the numerical integrations be general enough that they can be tabulated and later plugged into closed-form expressions.
- 9. The final distribution should be an analytical function with a fixed form and described by a few parameters. This way, only a few real numbers are sufficient to entirely describe the curve. Piecewise, implicit, or recursive solutions are undesirable.

1.2 Physical Model

We would like to propose a system model that is as general as possible, while retaining enough simplicity for analytical tractability. Our model is based closely on those used in [1]–[3], which apply to flat-fading channels, and is thus well suited for frequency-hopping spread-spectrum (FH-SS), as well as orthogonal frequency division multiple access (OFDMA). The wireless system is divided into cells, each served by a base station (BS). However, we may also study a wireless system with fixed relays [2], where each relay station (RS) defines a sub-cell around itself, and

the MU connects to the nearest RS or BS. In order to encompass both conventional and relayed networks, we call a BS or RS by the common name of fixed station (FS), the exact nature of which is irrelevant to the analysis at this level. We study systems where there can be at most one other user per FS using the same frequency channel at the same time, thus one interferer per FS. The various interfering signals can be considered incoherent amongst themselves, thus the total interference power received can be written as:

$$I = \sum_{i=1}^{N} I_{i} = \sum_{i=1}^{N} R_{i} \zeta_{i} \chi_{i} \alpha_{i}^{2} P_{i}.$$
 (1.1)

Here there are N potential interferers (equivalently, N FS's), and I_i is the interference due to FS i. In the UL, there are N interfering mobile users (IMU), each transmitting to their own FS a signal that is also received as interference by the desired FS. In the DL, it is the interfering fixed stations (IFS) that transmit to their own MU's a signal that is also received as interference by the DMU. R_i is the average path gain (inverse of pathloss) to the interferer. ζ_i and α_i^2 represent shadowing and fading respectively between the interferer i and the receiver. χ_i is the Bernoulli-distributed indicator function that is unity when there is an active user in interfering cell C_i on the channel under consideration, otherwise zero. P_i is the (normalised) transmit power coming from the interferer in cell C_i .

1.2.1 Wireless Channel

1.2.1.1 Fading

We consider a system such as OFDMA or FH-SS where flat fading occurs over any single channel. There are several fading models in existence. We consider a Rician model with parameter K. When K=0, the model becomes Rayleigh, and the fading power factor is exponentially distributed with parameter $\lambda=1$ such that $E\left[\alpha_i^2\right]=1$. When $K=\infty$, there is no

fading in the model. In general, for Rician fading, such that $E[\alpha_i^2] = 1$, we have the following probability density function (pdf) [25]:

$$f_{\alpha_i}(x) = 2x(K+1)\exp(-x^2(K+1) - K)I_0(2x\sqrt{K^2 + K}), x > 0,$$
 (1.2)

where I_0 is the modified Bessel function of the first kind. The fading power multiplied by 2(K+1) follows a non-central χ^2 distribution with 2 degrees of freedom and non-centrality parameter 2K:

$$\frac{\partial}{\partial x} P(\alpha_i^2 2(K+1) < x) = \frac{1}{2\sqrt{2x(K+1)}} f_{\alpha_i} \left(\sqrt{\frac{x}{2(K+1)}} \right)$$

$$= \frac{1}{2} \exp\left(-\frac{1}{2} x - K\right) I_0 \left(\sqrt{2xK}\right), x > 0.$$
(1.3)

1.2.1.2 Shadowing

The interfering signal also suffers long-term shadowing, which follows a lognormal (LN) law, with $\sigma(dB) = 6$ to 12dB spread. All paths are assumed to be independently faded and shadowed. In this thesis, we use the natural base for LN RV's. Thus we convert from decibel units (typically used for shadowing) to natural units by multiplying by a factor of $\lambda = \ln(10)/10 \approx 0.23$. In natural units, $\sigma = \lambda \sigma(dB)$.

1.2.1.3 Pathloss

The average distance-dependent pathloss model is [2]:

$$\frac{1}{R_i} = r_0^{-\beta} \left(\frac{4\pi r_0 f}{c}\right)^2 r_i^{\beta}, \qquad (1.4)$$

where r_0 is the close-in reference distance, f the carrier frequency, c the speed of light, β is the propagation exponent and r_i is the distance between the transmitter and receiver. Without loss of generality, we can ignore the constant factor, in which case R_i is equal to $r_i^{-\beta}$, where r_i is the

physical normalised distance between the interferer i and the receiver. In general, R_i 's have different distributions, and are not independent for different i when considering the DL.

1.2.2 Wireless System

We use the indicator k_{DL} to specify whether we are looking at the UL or the DL. In the UL (k_{DL} = 0), the desired FS can receive interference signals from MU's in other cells. In the DL (k_{DL} = 1), it is the MU under analysis, located in C_0 , that receives interference from other FS's, which are transmitting to their own MU's. The marginal distribution of the interference power from a given cell C_i is identical for UL and DL. However, only for the UL are the pathlosses independent for each interferer.

1.2.2.1 Cellular Network

We assume that all the cells are identical regular hexagons of unit side length, each with its FS in the center. A MU is assumed to be connected to its nearest FS, i.e., it lies in the corresponding hexagon. The position of the MU is assumed random, uniformly distributed over the cell. We call the cell under consideration C_0 , and all potentially interfering cells C_i .

Although we assume all cells to be hexagons, there is nothing in the analysis that prevents us from considering, for example, square cells in a Manhattan-style grid. The integrals and distributions that involve the cell geometry would need to be modified.

1.2.2.2 Power Control

We use the indicator k_{PC} to specify whether we are looking at a system with or without PC. Without loss of generality, we assume that, without PC ($k_{PC} = 0$), all interferers are transmitting at unit power. Under PC ($k_{PC} = 1$), the power P_i transmitted by the ith interferer has identical statistics for UL and DL. The PC mechanism is such that it attempts to compensate for shadowing

and pathloss is such a way that a FS receives the same power from each of its MU's, within some error margin [1], [3]. The transmit power P_i depends on the pathloss between the IMU and its IFS, the LN shadowing with spread σ and the PC LN error with spread $\sigma_e(dB) = 0$ to 1dB. The PC mechanism is assumed slow: it does not compensate for signal fading.

1.2.2.3 User Activity

We assume that each MU within a cell uses an orthogonal channel. Consequently, there is only intercellular interference, and there can be at most one interferer in each cell C_i . Thus, the presence of interference from a particular cell C_i can be modeled by a Bernoulli RV with parameter p (collision probability), which is a function of loading, number of channels, channel reuse, DTX [1], etc. It is an appropriate model for any channel in a FH-SS or OFDMA system, where there may or may not be activity on each given carrier, and the frequencies are assigned orthogonally within each cell.

We assume that the Bernoulli RV that models the activity is independent of all other quantities. In reality, the activity is generally dependent on the interference at a given point in time, because a new user may be admitted to the channel based on the activity level of other users on that channel. However, we do not model this effect here.

1.3 Main Contributions

The main analytical contributions of this thesis are:

- 1. A closed-form expression for the lower tail asymptote to the SLN distribution. It is given by (3.9), (3.11), (3.12).
- 2. Building on ideas from [3], and based on the Fenton-Wilkinson (FW) [26] method for SLN analysis, a closed-form expression (4.39), (4.40), (4.34), (4.35), for the interference in the

model described in Section 1.2. Our expression separates the cellular layout from other parameters, and we provide a lookup table from which the parameters describing the layout can be obtained. This method performs well in the body and tail, for lower probabilities of activity.

3. Using the two previous points, a novel method for solving the SLN problem using a power-lognormal (PLN) distribution. This method is described by (3.35), (3.48)-(3.51). We apply this method to obtain simple forms for the distribution parameters (4.41)-(4.46), (4.52), (4.54). Only one simple numerical integral is required (3.50). This method performs well over the entire range of values for high activity levels.

Furthermore, this thesis contains:

- 4. A survey of various properties of the SLN distribution.
- 5. A thorough survey of methods to solve the SLN problem, including the most recent findings (end of 2006).
- 6. A formal framework to study the tails of the SLN distribution, and a study of the tails of several recent SLN methods [27]–[30].
- 7. A four-parameter function that approximates the SLN *cdf* with reasonable accuracy over the entire range of values, and whose parameters are found in closed-form.
- 8. Several ideas for the continuation of this work with suggested approaches.

Chapter 2

Known Methods for Approximating the Distribution of the Sum of Lognormal Random Variables

The fundamental mathematical problem that lies behind the interference analysis is calculating the distribution of the SLN. Although the SLN distribution is not known in closed form, there are nevertheless many things that we do know about it. We give an overview of the properties of the SLN distribution, and of the methods used in literature to approximate it.

2.1 Properties of the SLN Distribution

2.1.1 Definition

The problem can be mathematically stated as follows: Let

$$X = \sum_{i=1}^{N} Y_i \ . \tag{2.1}$$

We assume in this thesis that the summands Y_i are independent. The pdf of each summand is LN:

$$Y_i \sim f_i(x; \mu_i, \sigma_i) = \frac{1}{\sqrt{2\pi}x\sigma_i} e^{-\frac{1}{2}\left(\frac{\ln x - \mu_i}{\sigma_i}\right)^2}, x > 0.$$
 (2.2)

Its moments are:

$$EY_{i}^{n} = e^{n\mu + \frac{1}{2}n^{2}\sigma_{i}^{2}}, n \in \Re.$$
(2.3)

The *cdf* is:

$$F_i(x; \mu_i, \sigma_i) = \Phi\left(\frac{\ln x - \mu_i}{\sigma_i}\right), x > 0,$$
(2.4)

where $\Phi(x)$ is the standard normal *cdf*. Then

$$X \sim f(x) = \frac{\partial}{\partial x} F(x), x > 0$$
 (2.5)

follows the SLN distribution, with no known closed form. When the summands Y_i are independent, f(x) is the convolution of all $f_i(x)$:

$$f(x) = f_1(x) \otimes f_2(x) \otimes \dots \otimes f_N(x). \tag{2.6}$$

Again, a closed-form method for performing even one of these convolutions is not known.

2.1.2 Properties

2.1.2.1 Analytically Smooth

By smooth we mean that all derivatives of the SLN *pdf* exist everywhere. This can be proved. From (2.5):

$$\left(\frac{\partial}{\partial x}\right)^n f(x) = \left[\left(\frac{\partial}{\partial x}\right)^n f_1(x)\right] \otimes f_2(x) \otimes \cdots \otimes f_N(x). \tag{2.7}$$

Since the LN $pdf f_1(x)$ has all derivatives, so does the SLN pdf. We would like the function approximating the SLN to also have this property, if possible.

2.1.2.2 Possibly Unimodal

It is not evident that the SLN *pdf* is always unimodal. In fact, a convolution of unimodal functions needs not be unimodal [31]. However, various researchers have worked on trying to estimate the

mode of the SLN *pdf* [18], [32]; and simulation results [27], [28] seem to confirm that the SLN *pdf* is indeed unimodal.

2.1.2.3 Heavy-Tailed and Sub-Exponential

Heavy-tailed (h.t.) and sub-exponential (s.e.) distributions are defined in [33]. The LN distribution is said to be h.t., which is formally defined as follows:

The distribution F_i is h.t. if and only if:

$$\lim_{x \to +\infty} \frac{1 - F_i(x + y)}{1 - F_i(x)} = 1, \forall y > 0.$$
 (2.8)

Furthermore, a subclass of *h.t. cdf* s is the set of *s.e. cdf* s. A distribution is said to be *s.e.*, if and only if:

$$\lim_{x \to +\infty} \frac{1 - F_i^{*n}(x)}{1 - F_i(x)} = n , \qquad (2.9)$$

where *n indicates the n-fold auto-convolution. By L'Hopital's rule, we may write this in terms of pdf's:

$$\lim_{x \to +\infty} \frac{f_i^{*n}(x)}{f_i(x)} = n, \qquad (2.10)$$

Now a set of s.e. distributions is not closed under convolution [33]. However, we can show that the SLN for independent terms, or equivalently, the convolution of N LN pdfs is s.e., and thus h.t..

Proof:

It is enough to show that the convolution of any two LN *pdf* s is *s.e.*, since Theorem 3 in [34] says that if all pair-wise convolutions of a set of *pdf* s are *s.e.*, then all manifold convolutions on this set are also *s.e.*.

If two LN pdf's are identical, then their convolution is s.e., by Proposition 2.8 in [33].

If two LN pdf's are different, then let us call them $f_1(x; \mu_1, \sigma_1)$ and $f_2(x; \mu_2, \sigma_2)$, such that $\sigma_1 > \sigma_2$; or, if $\sigma_1 = \sigma_2$, such that $\mu_1 > \mu_2$. Then:

$$\lim_{x \to +\infty} \frac{1 - F_2(x)}{1 - F_1(x)} = \lim_{x \to +\infty} \frac{f_2(x)}{f_1(x)}$$

$$= \lim_{x \to +\infty} \frac{\sigma_1}{\sigma_2} \exp \left[-\frac{1}{2} \left(\frac{1}{\sigma_2^2} - \frac{1}{\sigma_1^2} \right) \ln^2 x + \left(\frac{\mu_2}{\sigma_2^2} - \frac{\mu_1}{\sigma_1^2} \right) \ln x - \frac{1}{2} \left(\frac{\mu_2^2}{\sigma_2^2} - \frac{\mu_1^2}{\sigma_1^2} \right) \right]$$

$$= \begin{cases} \lim_{x \to +\infty} x^K = 0, K = \frac{\mu_2 - \mu_1}{\sigma_1^2} < 0, & \sigma_1 = \sigma_2 \\ \lim_{x \to +\infty} \frac{\sigma_1}{\sigma_2} \exp \left[K \ln^2 x \right] = 0, K = -\frac{1}{2} \left(\frac{1}{\sigma_2^2} - \frac{1}{\sigma_1^2} \right) < 0, & \sigma_1 > \sigma_2 \end{cases} = 0.$$
(2.11)

From Proposition 2.7 in [33], we find that the convolution $f_1 \otimes f_2$ is s.e..

This proves that the SLN is *s.e.* for independent summands. In [21], it was proved that the tail behaviour is the same for correlated summands. Thus, in general, the SLN distribution is *sub-exponential*, and thus *heavy-tailed*.

2.1.3 Location in a Function Space

The location of the SLN pdf in a function space is considered in [18]. The LN distribution has two parameters, $\sigma > 0$, μ , and as such can be considered to generate a half-plane in a function space, as shown in Figure 2-1. A LN pdf with small $(\sigma^2 << 1)$ variance is nearly Gaussian, because the two distributions are related by exponential function, which can be seen as locally

linear when small variances are involved. Also, since the LN distribution has finite moments (2.3), the Central Limit Theorem applies, and the successive convolutions of LN pdf's, properly normalised, converges to a Gaussian distribution. However, this convergence will be extremely slow for higher $\sigma^2 \gg 1$, which is our case of interest in cellular communications: $\sigma \geq 6dB = 1.38$ in linear scale. For these higher variances, the moments (2.3) are very large $(EY_i^2 \gg E^2Y_i)$, which makes convergence slow. This can also be seen as a consequence of the SLN distribution being heavy-tailed (see Section 2.1.2.3): because the Gaussian distribution is not h.t., while the sum of any number of LN RV's is, the SLN will not converge to a Gaussian distribution for any finite number N of summands. Several SLN methods, [26], [35]–[37] assume that the trajectory of successive convolutions lies near to the LN half-plane, effectively meaning that the SLN is approximately LN. However, this has proven not to be the case in many situations, and more recent work [27]–[30] suggests that new cdf forms need to be developed for approximating the SLN.

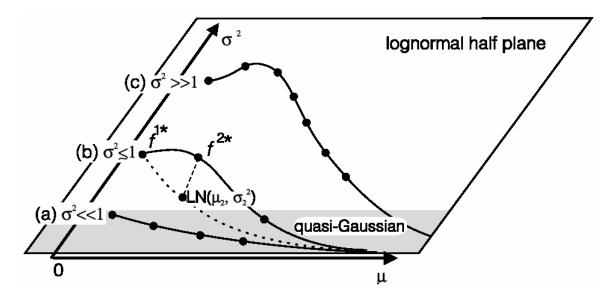


Figure 2-1 Locus of successive convolutions of LN pdfs in a function space. Taken from [18].

2.1.4 Lognormal Probability Paper

2.1.4.1 A Bijective Transformation

It is convenient to look at the *cdf* of the SLN on lognormal probability paper [35], where LN distributions are mapped onto straight lines by the following transformation *T*:

$$T: F(x) \mapsto \widetilde{F}(x) = \Phi^{-1}(F(e^x)). \tag{2.12}$$

Indeed this transformation maps:

$$T: \Phi\left(\frac{\ln x - \mu_i}{\sigma_i}\right) \mapsto \frac{x - \mu_i}{\sigma_i}, \tag{2.13}$$

which is a straight line of strictly positive slope. Conversely, each line of strictly positive slope corresponds to a unique LN distribution.

2.1.4.2 Order Conservation

Lemma 1:

For any supports I, J, such that $x \in I \Leftrightarrow e^x \in J$ we have:

$$\forall x \in I, \tilde{F}(x) > \frac{x - \mu}{\sigma} \iff \forall x \in J, F(x) > \Phi\left(\frac{\ln x - \mu}{\sigma}\right),$$
 (2.14)

and

$$\forall x \in I, \, \tilde{F}(x) < \frac{x - \mu}{\sigma} \iff \forall x \in J, \, F(x) < \Phi\left(\frac{\ln x - \mu}{\sigma}\right).$$
 (2.15)

Proof:

Let
$$g(x) = \tilde{F}(x) - \frac{x - \mu}{\sigma} > 0$$
 and

$$T: \Phi\left(\frac{\ln x - \mu}{\sigma} + g(\ln x)\right) = F(x) \mapsto \widetilde{F}(x) = \frac{x - \mu}{\sigma} + g(x).$$
 Since

 $\forall x \in I, g(x) > 0 \iff \forall x \in J, g(\ln x) > 0$ and $\Phi(x)$ is strictly increasing, we have:

$$\forall x \in J, F(x) > \Phi\left(\frac{\ln x - \mu}{\sigma}\right)$$
, which proves (2.14). The proof is analogous for (2.15).

2.1.4.3 Concavity

Based on the all the simulation and numerical curves in literature, it would seem that the SLN *cdf* is always concave down on LN paper. This assumption is made in [35]. However, we are not aware of any proof of this.

Because it is concave, the transformed cdf is no longer linear, and in fact becomes less and less so as the number of summands N increases, as seen in Figure 2-2 for independent, identically distributed (i.i.d.) RV's. This means that a LN distribution cannot, in general, be a good approximation to the SLN.

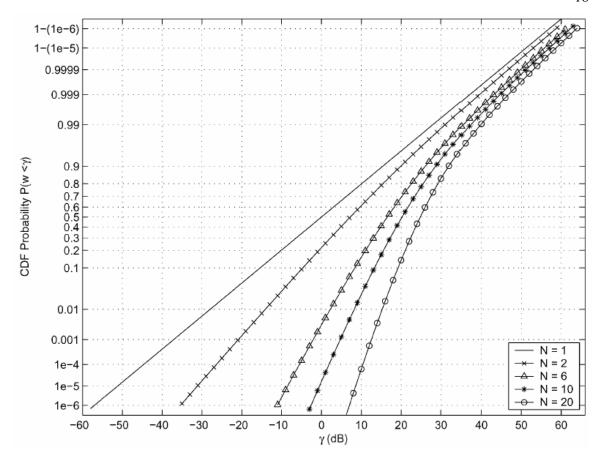


Figure 2-2 *cdf* on LN paper of the *i.i.d.* SLN with $\sigma = 12$ dB. Taken from [35].

2.1.4.4 Asymptotes

Again, looking at all the simulation and numerical curves in literature, it would appear that, on LN paper, the *cdf* has linear asymptotes at both limits. This would mean that the SLN distribution could be well-approximated by two (different) LN distributions at each tail. This fact is theoretically asserted for the upper tail [21], but there has been no mathematical evidence for there being an asymptote in the lower tail. Also, until now, most SLN methods have completely ignored the nature of the tails of the *cdf*. We explore these asymptotes formally in Chapter 3.

2.2 Closed-Form Methods

We are interested primarily in developing a closed-form solution to our interference problem, which means we are particularly interested in closed-form results on the SLN.

2.2.1 Positive Moment-Matching

The simplest approach to approximating the SLN is to assume that the SLN distribution is LN, and match its first two positive moments. This is the first method for solving the SLN problem, and is known as the Fenton-Wilkinson (FW) method [26]. The moments of a LN distribution are known (2.3). Let us use the problem statement in Section 2.1.1, and let us approximate the sum of the SLN RV *X* by a LN:

$$X \sim F(x) \approx \Phi\left(\frac{\ln x - \mu_{FW}}{\sigma_{FW}}\right).$$
 (2.16)

The first two moments of the SLN can be calculated as follows [38]:

$$E[X] = E\left[\sum_{i=1}^{N} Y_i\right] = \sum_{i=1}^{N} E[Y_i] = \sum_{i=1}^{N} e^{\mu_i + \frac{1}{2}\sigma_i^2}, \qquad (2.17)$$

$$E[X^{2}] = E\left[\left(\sum_{i=1}^{N} Y_{i}\right)^{2}\right] = E\left[\sum_{i=1}^{N} Y_{i}^{2} + 2\sum_{i=1}^{N} \sum_{j>i}^{N} Y_{i}Y_{j}\right] = \sum_{i=1}^{N} E[Y_{i}^{2}] + 2\sum_{i=1}^{N} \sum_{j>i}^{N} E[Y_{i}Y_{j}]$$

$$= \sum_{i=1}^{N} e^{2\mu_{i} + 2\sigma_{i}^{2}} + 2\sum_{i=1}^{N} \sum_{j>i}^{N} e^{\mu_{i} + \mu_{j} + \frac{1}{2}(\sigma_{i}^{2} + \sigma_{j}^{2}) + \rho_{ij}\sigma_{i}\sigma_{j}}.$$
(2.18)

Here ρ_{ij} is the correlation coefficient between the Gaussian RV's $\ln Y_i$ and $\ln Y_j$. These moments are then matched with those of (2.16):

$$\begin{cases} E[X] = e^{\mu_{FW} + \frac{1}{2}\sigma_{FW}^2} \\ E[X^2] = e^{2\mu_{FW} + 2\sigma_{FW}^2} \end{cases} \Leftrightarrow \begin{cases} \mu_{FW} = 2 \ln E[X] - \frac{1}{2} \ln E[X^2] \\ \sigma_{FW}^2 = -2 \ln E[X] + \ln E[X^2] \end{cases}$$
(2.19)

Matching higher moments has also been proposed [9], [26], [39]. In general, we find the moments of order n and n+1 in a manner similar to (2.17), (2.18), then match the moments according to:

$$\begin{cases}
E[X^n] = e^{n\mu_{FW} + \frac{1}{2}n^2\sigma_{FW}^2} \\
E[X^{n+1}] = e^{(n+1)\mu_{FW} + 2(n+1)^2\sigma_{FW}^2}
\end{cases} \Leftrightarrow
\begin{cases}
\mu_{FW} = \frac{n+1}{n} \ln E[X^n] - \frac{n}{n+1} \ln E[X^{n+1}] \\
\sigma_{FW}^2 = -\frac{2}{n} \ln E[X^n] + \frac{2}{n+1} \ln E[X^{n+1}]
\end{cases} .$$
(2.20)

Now matching higher and higher moments of positive RV's tends to give much more weight to the upper tail of the distribution. More specifically:

- Moment-matching tends to approximate the upper tail much better than the body or the lower tail [26], [35], [40], since moments essentially determine the tail of a distribution [43].
- 2. By matching higher moments as in (2.20), the overall curve match to the body gets worse and worse, and approximates the far upper tail more and more.
- 3. Based on the last point, it is possible to perform successive moment-matching for various *n* and construct a piecewise-LN method [39], [40]. However, we dislike a piecewise solution because it is cumbersome to use.
- 4. As the moment order *n* goes to infinity, the summand with the highest moments dominates, and the moment-matching method converges to a LN asymptote determined by the dominating term. This will be explored in more detail in Section 2.2.2.

In conclusion, moment-matching is a simple-closed form solution to the SLN problem. There exist known expressions for bounds on the error of this method [41].

2.2.2 Asymptotics

The behaviour of the SLN distribution F(x) as $x \to +\infty$ has been previously studied [21], [42]. It was shown in [42] that the moments $E[X^n]$ of a SLN are dominated by the heaviest summand(s) as n increases. In [40], it is suggested that matching higher and higher moments of the sum to those of a LN distribution gives a better and better LN asymptote to the upper tail behaviour. In [43], it is argued that moments of positive RV's (as is our case) determine the upper tail. All this information suggests that looking at the heaviest summand(s) alone will give us a good approximation of the tail.

Reference [21] gives a simple procedure to find the upper tail asymptote:

Choose the summands with the highest logarithmic standard deviation:

$$\sigma_U = \max_{i=1,\dots,N} \sigma_i. \tag{2.21}$$

Among these summands, chose those with the highest logarithmic mean:

$$\mu_U = \max_{i:\sigma_i = \sigma_U} \mu_i. \tag{2.22}$$

Count the number of these "heaviest" summands:

$$\kappa_U = \#\{(\mu_i, \sigma_i) = (\mu_U, \sigma_U)\}, \qquad (2.23)$$

where $\#\{S\}$ is the cardinality of set S.

The tail distribution is then that of the heaviest summand(s), multiplied by the number of these heaviest summands:

$$f(x) \stackrel{x \to +\infty}{\sim} f_U(x) = \frac{\kappa_U}{\sqrt{2\pi} x \sigma_U} e^{-\frac{1}{2} \left(\frac{\ln x - \mu_U}{\sigma_U}\right)^2}$$
 (2.24)

$$F(x) \stackrel{x \to +\infty}{\sim} F_U(x) = \kappa_U \Phi\left(\frac{\ln x - \mu_U}{\sigma_U}\right) + (1 - \kappa_U). \tag{2.25}$$

Note that [21] proves that this result is also valid when the summands with the same marginal distributions are correlated.

2.2.3 Bounds

2.2.3.1 Tighter, Non-Parametric Approximation

It is argued in [21] that if a summand is not heaviest, but nearly so, it also contributes significantly to the tail for large but bounded x. In [44] are given bounds that are tight in the upper tail. Thus, a better asymptote is the distribution of the maximum of the summands:

$$\sum_{i=1}^{N} Y_{i} \ge \max_{i} Y_{i}$$

$$P\left(\sum_{i=1}^{N} Y_{i} < x\right) = F(x)^{x \to +\infty} F_{UB}(x) = \prod_{i=1}^{N} \Phi\left(\frac{\ln x - \mu_{i}}{\sigma_{i}}\right) = P\left(\max_{i} Y_{i} < x\right) \ge F(x).$$
(2.26)

Indeed, when we take the sum of positive terms that have high variance, in most outcomes there will be one term that is much larger than the others, and the maximum of the RV's is a good lower bound to their sum. In the *i.i.d.* case this bound becomes a PLN distribution, which is called Farley's method in literature [45]:

$$F(x) \stackrel{x \to +\infty}{\sim} \Phi^{N} \left(\frac{\ln x - \mu}{\sigma} \right), \left[\mu_{i}, \sigma_{i} \right] = \left[\mu, \sigma \right] \forall i.$$
 (2.27)

We also have:

$$\sum_{i=1}^{N} Y_i \le N \max_{i}$$

$$P\left(\sum_{i=1}^{N} Y_i < x\right) = F(x)^{x \to +\infty} F_{LB}(x) = \prod_{i=1}^{N} \Phi\left(\frac{\ln x - \ln N - \mu_i}{\sigma_i}\right) = P\left(N \max_{i} Y_i < x\right) \le F(x).$$
(2.28)

Equations (2.26) and (2.28) are upper and lower bounds respectively to the SLN *cdf*, which are called lower and upper bounds respectively in [44], since they are used to refer either to the complementary *cdf* or to the value of the RV itself.

2.2.3.2 Arithmetic-Geometric Mean Inequality

While the sum of LN RV's is difficult to find, the product of LN is very simple. Let

$$Z = \prod_{i=1}^{N} Y_i \,. \tag{2.29}$$

Then Z is simply a LN RV with parameters:

$$\mu_Z = \sum_{i=1}^{N} \mu_i, \quad \sigma_Z^2 = \sum_{i=1}^{N} \sigma_i^2$$
 (2.30)

assuming Yi are independent.

The method proposed in [46] uses the geometric-arithmetic (AG) mean inequality: for any positive quantities, their geometric mean is always upper-bounded by their arithmetic mean:

$$\forall V \middle| V = \{x_i \ge 0\}_{i=1}^N : \frac{1}{N} \sum_{i=1}^N x_i \ge \left(\prod_{i=1}^N x_i\right)^{\frac{1}{N}}.$$
 (2.31)

In our SLN case,

$$X = \sum_{i=1}^{N} Y_i \ge N \left(\prod_{i=1}^{N} Y_i \right)^{\frac{1}{N}} = N Z^{\frac{1}{N}},$$
 (2.32)

from which we obtain a bound on the SLN distribution:

$$F(x) = P(X < x) \le P\left(NZ^{\frac{1}{N}} < x\right) = P\left(Z < \left(\frac{x}{N}\right)^{N}\right) = \Phi\left(\frac{N\ln x - \mu_{Z} - N\ln N}{\sigma_{Z}}\right). \quad (2.33)$$

Now let us examine the inequality (2.31) with some numerical examples. Take $V = \{4,5,6\}$, where all the terms are of the same order of magnitude. Then the arithmetic mean is 5, and the geometric mean is 4.932.... Not only is the bound correct, it is also very tight in this case. Let us take another example: $V = \{1,10,100\}$, with different orders of magnitude. Then the arithmetic mean is 37, and the geometric mean is 10, which is a very loose lower bound.

From this we come to the conclusion that if our LN summands are of different μ_i or of high σ_i , they will take values of very different orders of magnitude. This is indeed the case in our problem, where there are interferers located both near and far, and the LN shadowing is of at least 6dB. Hence the bound in (2.33) will be very loose, and essentially useless.

This bound is however very useful in solving another problem: calculating the total interference power from in-cell users in a power-controlled Code Division Multiple Access (CDMA) system. In this case, all receive powers are essentially the same, but vary slightly in a LN manner, say by 1dB, due to imperfect PC. In this case the method, which was only tested for up to 4dB, is good for 4dB and excellent for 1dB [46].

2.3 Numerical Methods

The methods in this class consist in calculating numerical integrals, perhaps iteratively, and perhaps in several dimensions.

2.3.1 Parametric Methods

Here are methods that use numerical integration in order to find the parameters of some function that approximates the SLN *cdf*.

2.3.1.1 Logarithmic Moment-Matching: Schwartz - Yeh

This is one of the earliest and most important methods, originally published in [37], it was used in [1], [47] to analyse outage probability in cellular systems. The method is described in [38], so we only give a brief description here:

The SY method approximates the SLN by a LN distribution. Moment-matching is performed in the logarithmic, instead of the linear, domain. The sum of the first two LN RV's is approximated by a partial sum X_2 , which is LN with parameters μ_{X_2} , σ_{X_2} which are found by solving:

$$E[\ln X_2] = \mu_{X_2} = E[\ln(Y_1 + Y_2)]$$
 (2.34)

$$E[\ln X_2^2] = \mu_{X_2}^2 + \sigma_{X_2}^2 = E[\ln^2(Y_1 + Y_2)]$$
 (2.35)

Equation (2.34) is reformulated as follows:

$$\mu_{X_2} = E[\ln Y_1] + E\left[\ln\left(1 + \frac{Y_2}{Y_1}\right)\right]$$

$$= \mu_1 + \int_0^{+\infty} \ln(1+x) \frac{1}{\sqrt{2\pi}x(\sigma_1^2 + \sigma_2^2)} \exp\left[-\frac{1}{2}\left(\frac{\ln x - (\mu_2 - \mu_1)}{\sqrt{\sigma_1^2 + \sigma_2^2}}\right)^2\right] dx$$
(2.36)

This formulation is due to the fact that Y_2/Y_1 is also LN. The integral cannot be calculated in closed form and is thus left to numerical integration. A similar formulation is used to find σ_{X_2} .

Once the sum of two LN RV's has been approximated by another LN, we proceed iteratively by solving in the same manner:

$$E[\ln X_n] = \mu_{X_n} = E[\ln(X_{n-1} + Y_n)], \qquad (2.37)$$

$$E\left[\ln X_n^2\right] = \mu_{X_n}^2 + \sigma_{X_n}^2 = E\left[\ln^2(X_{n-1} + Y_n)\right]. \tag{2.38}$$

This is a recursive method with respect to the number of summands. In total we perform N-1 iterations, each involving a numerical integral. In order to calculate this integral more efficiently and accurately, modified methods have been proposed [48], [49].

2.3.1.2 Log-Shifted Gamma Function: Le-Ngoc - Lam

The LL method was proposed in [27], and is inspired by the SY method. A three-parameter function: a log-shifted Gamma (LSG) function (see Section 3.2.1.1), is used as an approximating *cdf*. The sum of two LN RV's is approximated by a LSG, after which each successive partial sum is computed as the sum of a LN and a LSG RV. The three parameters of each successive LSG distribution are found through matching the first two logarithmic moments of the sum (as in SY), and also the first linear moment (as in FW). This method also requires N-1 iterations, and numerical integrations at every step. The method is expanded to include the sum of correlated LN RV's in [50].

2.3.1.3 Pearson Type IV Function: Zhang - Song

A four-parameter Pearson Type IV (see Section 3.2.1.4) *cdf* is used by Zhang – Song (ZS) in [28] as an approximating distribution to the logarithm of a SLN. These parameters are found by matching for the first four moments of the logarithm of a sum of LN RV's, given by:

$$E[\ln X^n] = E\left[\ln^n\left(\sum_{i=1}^N Y_i\right)\right], n = 1,2,3,4.$$
 (2.39)

It is not specified how these moments are calculated, but it is probably through numerical integration.

2.3.1.4 Negative Moment-Matching

Rather than matching positive moments of the SLN as in Section 2.2.1, [21] proposed to match the first two negative moments (powers -2 and -1). The advantage of this method is that it gives less weight to the higher tails, and more to the lower portion of the cdf. The disadvantage is that there is no closed-form expression for the negative moments of a sum of RV's, and an N-dimensional numerical integration is necessary.

This method was extended in [40] so as to use (2.20) for both positive and negative n except 0 and -1. This gives us a piecewise solution that is accurate in both tails, but again requires numerical integration for negative n, which give the lower tail. The method does not yield a very exact match in the body (centre) of the cdf though, and produces a piecewise solution.

2.3.1.5 Flexible Lognormal Approximation: Wu – Mehta – Zhang

The WMZ method presented in [36] approximates the SLN *cdf* by a LN. However, the LN *cdf* can be matched to any part of the SLN curve, by selecting parameters correctly. The method uses the LN distribution's moment-generating function. Since its closed form is not known, a short

Gauss-Hermite expansion of less than 20 terms is used. This method has the advantages of providing a piecewise LN solution of arbitrarily good accuracy given enough pieces, and this without recursion. However a non-linear equation still needs to be solved numerically.

The method is expanded in [51] to include correlated LN and LN-Rice RV's.

2.3.2 Characteristic Function Evaluation

Another numerical approach to the SLN problem for independent summands is to take (2.6) and consider it in the Fourier domain: the characteristic function (*cf*) of a sum of independent RV's is the product of their individual *cf*'s. Now the *cf* of a LN RV is not known in closed form, but can be found numerically. This results in an integral with a very long oscillating tail, and the integral converges very slowly. The first numerical study of the LN *cf* seems to be [19]. The modified Clenshaw-Curtis method is used in [35], while [52] uses a new method where the integral is transformed so that in no longer oscillates. Reference [53] proposes yet another integral. The *cf* method is also treated mathematically in [22], where it is shown that the SLN *cdf* for independent terms can be expressed as a sum of products of LN terms.

2.4 Curve-Fitting Methods

The following methods could be called "empirical", because they consist in already knowing the SLN cdf as a vector of values, and then fitting the best possible analytical curve to it. The SLN cdf is obtained numerically by methods from Section 2.3.2. These methods essentially correspond to a tabulation of curve parameters against the input parameters of the summed LN RV's. The weakness of this approach is that there exists an almost infinite variety of cases for the set of input parameters $\{\mu_i, \sigma_i\}_{i=1}^N$. Thus these methods only examine i.i.d. cases, where there are only two parameters: N and σ , while μ can be set to zero without loss of generality.

2.4.1 Minimax Lognormal Fitting: Beaulieu – Xie

After using the modified Clenshaw-Curtis method to find the cf, the BX method [35] assumes that the SLN cdf is concave on LN paper (see Section 2.1.4.3), and fits the best possible straight line on the LN paper to the SLN cdf using the minimax criterion over the range [10^{-6} , $1-10^{-6}$] of probabilities. The result is a LN approximation to the SLN distribution that does not privilege any particular region of the curve.

2.4.2 Three-Parameter Function: Beaulieu – Rajwani

The previous method (or any other method approximating the SLN by a LN) cannot be accurate everywhere, because in most cases the SLN cdf is not approximately LN (see Section 2.1.4.3). Another cdf is proposed by Beaulieu – Rajwani [30] (BR), which will be analysed in detail in Section 3.2.1.2. Here we will only say that the distribution has three parameters, which are found by a non-linear least-squares fit to the cdf obtained numerically. These parameters are tabulated for N = 2 to 20, and $\sigma = 6dB$ and 12dB.

2.4.3 Three-Parameter Function: Zhao – Ding

Another approach, similar to the previous one, uses also a three-parameter cdf that is parabolic under the transformation T. This is the ZD method [29], whose cdf will be examined in detail in Section 3.2.1.3. Beaulieu and Rajwani mention that they have also tried and discarded a parabolic function [30]. Our simulation results in Figure 3-2 show that the BR method greatly outperforms the ZD method. However, the ZD method has the advantage of being defined for a whole range of N and σ .

2.5 Validation of Methods

Given so many methods to choose from, we would like to analyse and compare them in order to choose the ones most suited to our purpose.

2.5.1 Comparisons in Literature

There are several papers in literature dealing with comparing FW and SY methods: [45] argues that FW is better in several cases, while [38] prefers SY. In fact which method is better depends on the choice of LN parameters [48]. Also, FW tends to approximate the upper tail better, while SY approximates the body, or middle portion of the *cdf*, as can be seen in [27], [35], [45]. Also Farley's method is always asymptotic to the SLN *cdf* in the upper tail, but is not very good in the lower tail, sometimes better and sometimes worse than FW [27], [35], [45].

2.5.2 Comparison and Classification of Methods

There exist various approximate solutions to the SLN, which generally consist of two distinct parts: choosing a general form for the SLN distribution, and obtaining its parameters for particular cases. One approach has been to approximate the SLN distribution by a LN, and various methods were used to find the best-fitting parameters: FW [26], BX [35], WMZ [36], SY [37], negative moment-matching [40], [54], and LN bound [46]. However, it has become clear that as the number of summands increases, the SLN distribution appears less and less LN (see Figure 2-2). There also exist other methods, which use more complex distribution functions with three or more free parameters: LL [27], ZS [28], ZD [29], and BR [30]. Nonparametric solutions have also been proposed: piecewise-lognormal [26], [39], [40], or product-of-lognormal (Farley's method [45], other bounds [44]). Each of these methods has its drawbacks: the first class of methods, where a LN approximation is used, cannot possibly be accurate over the entire range of *x*. Many methods require numerical integration to find the required parameters: [27], [28], [37],

[40], [54]. While the behaviour of the upper tail, i.e., as $x \to \infty$ is known [21], [42], that of the lower tail, i.e., as $x \to 0$ has no simple expression. Some methods have only been verified for the case of *i.i.d.* summands [29], [30]. Also, we will show that methods that match the body of the distribution well [27]–[30], are not necessarily accurate in the far tails. On the other hand, the piecewise-lognormal method proposed in [40] does give a good behaviour at the tails, but again requires numerical integration in N dimensions and does not perform as well in the body. Hence, there is no single method that really gives a complete solution to the problem, and there is no closed form for the lower tail.

Figure 2-3 shows the various methods compared according to both accuracy (precision) of the approximation (which is not entirely an objective measure, as different parts of the curve may be better approximated by different methods, and for different parameter choices.) and also by complexity. Complexity is again a partly subjective measure. We tried to estimate the difficulty of understanding and implementing a specific method, while also considering the level of integration required (iterative, multi-variate?). We also included our two new methods, PLN and 4P, introduced in Chapter 3, for comparison. We have indicated by an arrow that the advantageous region for a SLN method is both high precision and low complexity. As can be seen from the Figure, there are no known methods in that region, which indicates that the SLN problem remains essentially open.

Table 2-1 shows a history of the principal developments in the field of the SLN problem in the electrical engineering community. Some methods only apply to *i.i.d.* or independent but not necessarily identically distributed (*i.n.i.d.*) RV's. It does not include work on *cf*'s or comparative work.

Method	Form	Statistics Matched	Complexity	Scope
FW ('60) [26]	LN	$E[X], E[X^2]$	Closed-form	any
Schleher ('76) [39]	Piece-LN	Cumulants: $E[(X - E[X])^n]$	Closed-form	i.i.d.
SY ('82) [37], [38]	LN	$E[\ln X], E[\ln^2 X]$	Num. integrals, <i>N</i> –1 recursion	any
Farley (?) [45]	PLN	$\max_{i} Y_{i} \leq \sum_{i=1}^{N} Y_{i}$	Closed-form	any
Bounds ('01) [44]	Product of LN	$\max_{i} Y_{i} \leq \sum_{i=1}^{N} Y_{i}$ $\leq N \max_{i} Y_{i}$	Closed-form	i.n.i.d.
BX ('04) [35]	LN	Minimax fit	Numerically computed <i>cdf</i>	any
AG ('04) [46]	LN	$\frac{1}{N} \sum_{i=1}^{N} Y_i \ge \left(\prod_{i=1}^{N} Y_i \right)^{\frac{1}{N}}$	Closed-form	any
BR ('04) [30]	(3p.)	Non-linear least-squares fit	Numerically computed <i>cdf</i>	i.i.d.
LL ('05) [27], [50]	LSG (3p.)	$E[X], E[\ln X],$ $E[\ln^2 X]$	Num. integrals, <i>N</i> –1 recursion	any
WMZ ('05) [36], [51]	LN	Moment-generating function	Num. equations, Non-recursive	any
Negative Moment Match. ('05) [40], [54]	Piece-LN	$E[X^n], E[X^{n+1}],$ n = 3, -2, 1, 2, 3,	Num. integrals in <i>N</i> dimensions	i.n.i.d.
ZD ('06) [29]	(3p.)	Quadratic least-squares fit	Numerically computed <i>cdf</i>	i.i.d.
ZS ('06) [28]	Pearson IV (4p.)	$E[\ln^n X], n = 1,,4$	Closed-form *	any

^{*} It is not explained how the moments of the logarithm of the SLN are found. Once these are found, the method is closed-form.

Table 2-1 Summary of properties of SLN methods.

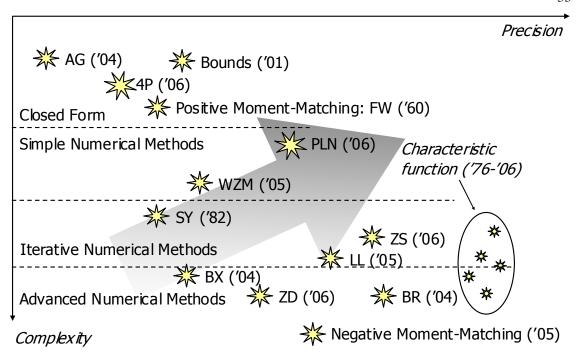


Figure 2-3 Comparison of SLN methods by complexity and precision.

2.5.3 Monte-Carlo Simulation Algorithm

Though various numerical methods exist to find the exact SLN *cdf* (see Section 2.3.2), we preferred to use Monte-Carlo simulations due to its simplicity in concept and implementation. The algorithm flowchart is given in Figure 2-4. The strongest drawback of this method is the long computational time. We mitigate this problem by letting the simulations run overnight.

We use the simulated curve as a benchmark for evaluating existing SLN methods, or, in Chapter 3, methods of our own device. We will also give an algorithm that efficiently simulates the far lower tail of the SLN *cdf* in Section 3.1.3.5.

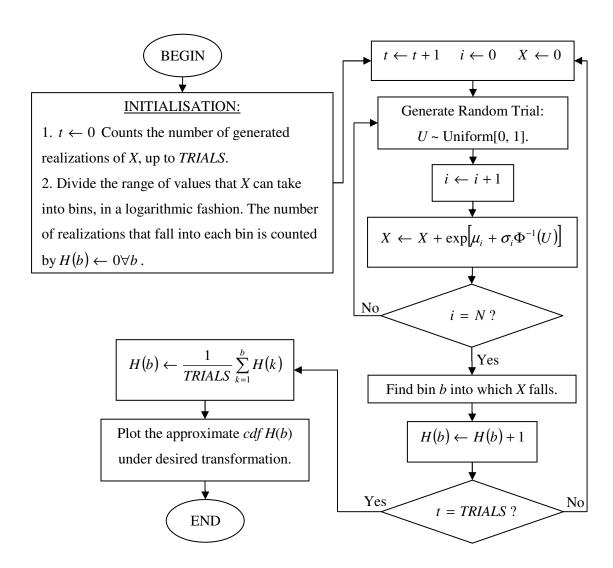


Figure 2-4 Algorithm that simulates the SLN cdf.

Chapter 3

Towards a Good, Simple, Closed-Form Approximation of the Distribution of the Sum of Lognormal Random Variables

In Chapter 2, we have made a thorough search of existing SLN methods and found that there is still no closed-form method that approximates the SLN *cdf* well. In particular, the closed-form methods gave us some good results in the upper tail, but failed in the lower tail. In this chapter, we develop a closed-form approximation to the lower tail, and combine it with other closed-form results to produce new, more exact closed-form or nearly-closed-form methods.

While it can be argued that knowledge of the tails of the distribution is less important than that of the body, we will show that theoretical results on the tails provide information on the body shape also, because the distribution is smooth (see section 2.1.2.1). Also, in order to apply moment matching on positive RV's, it is important to understand the upper tail [43], since a small error in the distribution at high values may translate into a large error in the moment.

3.1 Best Lognormal Fit to Either Tail

In Section 2.1.4.4 we have seen that, on lognormal paper, the tails of the SLN distribution seem to have linear asymptotes, which would indicate that the distribution is approximately LN in either tail. We will now examine this formally by defining a property of tails of distributions, the *best lognormal fit (blf)*, and showing what it means in terms of convergence. We then examine the properties of both tail of the SLN distribution using this framework.

3.1.1 Definition

Let $T: F(x) \mapsto \tilde{F}(x)$. Then $\frac{x-\mu}{\sigma}$ is an asymptote to $\tilde{F}(x)$ at $l=\pm\infty$ if and only if:

$$\lim_{x \to l} \tilde{F}(x) - \frac{x - \mu}{\sigma} = 0. \tag{3.1}$$

We say that F(x) has a *best lognormal fit* $\Phi\left(\frac{\ln x - \mu}{\sigma}\right)$ at e^l . At any given tail, a *cdf* may have at most one *blf*.

3.1.2 Best Lognormal Fit in the Linear Domain

We show that if, on lognormal paper, a line L is not the blf to a curve C at the limit l, then the corresponding distributions have tails in the corresponding limit e^l that are not comparable at any order of magnitude, i.e., their ratio will diverge. In other words, a necessary (but not sufficient) condition for approximating a distribution by a LN at a given tail is that the two distributions be asymptotes to each other on lognormal paper.

Lemma 2:

If
$$\lim_{x \to -\infty} \tilde{F}(x) - \frac{x - \mu}{\sigma} \neq 0$$
 or does not exist, then $\lim_{x \to 0} \frac{F(x)}{\Phi\left(\frac{\ln x - \mu}{\sigma}\right)} = 0, +\infty$ or does not exist.

Likewise, if
$$\lim_{x \to +\infty} \tilde{F}(x) - \frac{x - \mu}{\sigma} \neq 0$$
 or does not exist, then $\lim_{x \to +\infty} \frac{1 - F(x)}{1 - \Phi\left(\frac{\ln x - \mu}{\sigma}\right)} = 0, +\infty$ or

does not exist.

Proof:

$$\lim_{x \to -\infty} \tilde{F}(x) - \frac{x - \mu}{\sigma} \neq 0 \text{ or does not exist}$$

$$\Rightarrow \neg \left(\lim_{x \to -\infty} \tilde{F}(x) - \frac{x - \mu}{\sigma} = 0\right) \Leftrightarrow \neg \left(\forall \varepsilon > 0, \exists \delta, \forall x < \delta, \left| \tilde{F}(x) - \frac{x - \mu}{\sigma} \right| < \varepsilon\right)$$

$$\Leftrightarrow \exists \varepsilon > 0, \forall \delta, \exists x < \delta, \left| \tilde{F}(x) - \frac{x - \mu}{\sigma} \right| \ge \varepsilon$$

$$\Leftrightarrow \exists \varepsilon > 0, \forall \delta, \exists x < \delta, \tilde{F}(x) \ge \frac{x - (\mu - \sigma\varepsilon)}{\sigma} \lor \tilde{F}(x) \le \frac{x - (\mu + \sigma\varepsilon)}{\sigma}.$$

$$(3.2)$$

Applying Lemma 1:

$$\exists \varepsilon > 0, \forall \delta, \exists x < \delta, \widetilde{F}(x) \ge \frac{x - (\mu - \sigma \varepsilon)}{\sigma}$$

$$\Leftrightarrow \exists \varepsilon > 0, \forall \delta, \exists x : 0 < x < e^{\delta}, F(x) \ge \Phi\left(\frac{\ln x - (\mu - \sigma \varepsilon)}{\sigma}\right)$$

$$\Leftrightarrow \exists \varepsilon > 0, \forall \delta, \exists x : 0 < x < e^{\delta}, \frac{F(x)}{\Phi\left(\frac{\ln x - \mu}{\sigma}\right)} \ge \frac{\Phi\left(\frac{\ln x - (\mu - \sigma \varepsilon)}{\sigma}\right)}{\Phi\left(\frac{\ln x - \mu}{\sigma}\right)}$$

$$\Leftrightarrow \neg\left(\forall \varepsilon > 0, \exists \delta, \forall x : 0 < x < e^{\delta}, \frac{F(x)}{\Phi\left(\frac{\ln x - \mu}{\sigma}\right)} < \frac{\Phi\left(\frac{\ln x - (\mu - \sigma \varepsilon)}{\sigma}\right)}{\Phi\left(\frac{\ln x - \mu}{\sigma}\right)}\right)$$

$$\Rightarrow \neg\left(\forall \varepsilon > 0, \lim_{x \to 0} \frac{F(x)}{\Phi\left(\frac{\ln x - \mu}{\sigma}\right)} < \lim_{x \to 0} \frac{\Phi\left(\frac{\ln x - (\mu - \sigma \varepsilon)}{\sigma}\right)}{\Phi\left(\frac{\ln x - \mu}{\sigma}\right)} = \lim_{x \to 0} \frac{\partial \Phi\left(\frac{\ln x - (\mu - \sigma \varepsilon)}{\sigma}\right)}{\partial \Phi\left(\frac{\ln x - \mu}{\sigma}\right)}\right)$$

$$\Leftrightarrow \forall \varepsilon > 0, \lim_{x \to 0} \frac{F(x)}{\Phi\left(\frac{\ln x - \mu}{\sigma}\right)} \ge \lim_{x \to 0} \exp\left(-\frac{1}{2}\varepsilon^2 - \varepsilon\frac{(\ln x - \mu)}{\sigma}\right) = +\infty \Rightarrow \lim_{x \to 0} \frac{F(x)}{\Phi\left(\frac{\ln x - \mu}{\sigma}\right)} = +\infty.$$
(3.3)

Likewise:

$$\exists \varepsilon > 0, \forall \delta, \exists x < \delta, \widetilde{F}(x) \le \frac{x - (\mu + \sigma \varepsilon)}{\sigma}$$

$$\Leftrightarrow \exists \varepsilon > 0, \forall \delta, \exists x < e^{\delta}, F(x) \le \Phi\left(\frac{\ln x - (\mu + \sigma \varepsilon)}{\sigma}\right)$$

$$\Rightarrow \lim_{x \to 0} \frac{F(x)}{\Phi\left(\frac{\ln x - \mu}{\sigma}\right)} \le \lim_{x \to 0} \frac{\Phi\left(\frac{\ln x - (\mu + \sigma \varepsilon)}{\sigma}\right)}{\Phi\left(\frac{\ln x - \mu}{\sigma}\right)} = \lim_{x \to 0} \exp\left(-\frac{1}{2}\varepsilon^{2} + \varepsilon\frac{(\ln x - \mu)}{\sigma}\right) = 0.$$
(3.4)

But
$$\forall x, \frac{F(x)}{\Phi\left(\frac{\ln x - \mu}{\sigma}\right)} \ge 0$$
, thus:

$$\exists \varepsilon > 0, \forall \delta, \exists x < \delta, \tilde{F}(x) < \frac{x - (\mu + \sigma \varepsilon)}{\sigma} \Rightarrow \lim_{x \to 0} \frac{F(x)}{\Phi\left(\frac{\ln x - \mu}{\sigma}\right)} = 0 \text{ or does not exist. } (3.5)$$

Thus, if
$$\lim_{x \to -\infty} \tilde{F}(x) - \frac{x - \mu}{\sigma} \neq 0$$
 or does not exist, then $\lim_{x \to 0} \frac{F(x)}{\Phi\left(\frac{\ln x - \mu}{\sigma}\right)}$ is $0, +\infty$, or does not

exist. The proof for $x \to +\infty$ is analogous. This proves Lemma 2.

Lemma 2 has several important implications:

- 1. If a distribution has no linear asymptote in the LN domain at limit *l*, then it cannot possibly be approximated by a LN function in the linear domain at the corresponding tail.
- 2. If the ratio of two *pdf*'s converges to a non-zero constant at limit *l*, and one of the two distributions has a *blf*, then the other must have the same *blf*.

3.1.3 Lower Tail (for Independent Summands)

To our knowledge, there is no work dealing directly with the lower tail behaviour of the SLN. There is some literature on the lower tail of the distribution of the sum of positive RV's, but it does not readily apply to the SLN. The polynomial conditions required in [55] are not fulfilled by the LN distribution.

3.1.3.1 Discrete Convolution Approximation

In order to find the lower tail of the distribution of the sum of two independent LN RV's with pdf's $f_1(x)$ and $f_2(x)$, we need to find the convolution near 0 of two LN pdf's. We proceed as follows: let us sample the cdf's at 0, x, 2x, ..., where x is small. Then, we perform a discrete convolution and look at the first non-zero sample:

$$f(2x) \stackrel{x \to 0}{\approx} x f_1(x) f_2(x). \tag{3.6}$$

Proceeding likewise for *N* functions:

$$f(Nx) \stackrel{x \to 0}{\approx} x^{N-1} \prod_{i=1}^{N} f_i(x) . \tag{3.7}$$

$$f(Nx) \stackrel{x \to 0}{\approx} \frac{1}{x} \sqrt{2\pi}^{-N} \prod_{i=1}^{N} \sigma_{i}^{-1} e^{-\frac{1}{2} \sum_{i=1}^{N} \frac{(\ln x - \mu_{i})^{2}}{\sigma_{i}^{2}}} = \frac{1}{x} \sqrt{2\pi}^{-N} \prod_{i=1}^{N} \sigma_{i}^{-1} e^{-\frac{1}{2} \ln^{2} x \sum_{i=1}^{N} \frac{1}{\sigma_{i}^{2}} + \ln x \sum_{i=1}^{N} \frac{\mu_{i}}{\sigma_{i}^{2}} - \frac{1}{2} \sum_{i=1}^{N} \frac{\mu_{i}^{2}}{\sigma_{i}^{2}}} . \quad (3.8)$$

Let

$$\sigma_L = \left(\sum_{i=1}^N \frac{1}{\sigma_i^2}\right)^{-\frac{1}{2}},\tag{3.9}$$

then:

$$f(Nx) \stackrel{x \to 0}{\approx} = \frac{1}{x} \sqrt{2\pi}^{-N} \prod_{i=1}^{N} \sigma_{i}^{-1} e^{-\frac{1}{2} \frac{1}{\sigma_{L}^{2}} \left(\ln^{2} x - 2 \ln x \sigma_{L}^{2} \sum_{i=1}^{N} \frac{\mu_{i}}{\sigma_{i}^{2}} + \sigma_{L}^{2} \sum_{i=1}^{N} \frac{\mu_{i}^{2}}{\sigma_{i}^{2}} \right)}$$

$$= \frac{1}{x} \sqrt{2\pi}^{-N} \prod_{i=1}^{N} \sigma_{i}^{-1} e^{-\frac{1}{2} \frac{1}{\sigma_{L}^{2}} \left(\ln x - \sigma_{L}^{2} \sum_{i=1}^{N} \frac{\mu_{i}}{\sigma_{i}^{2}} \right)^{2} - \left(\sigma_{L}^{2} \sum_{i=1}^{N} \frac{\mu_{i}}{\sigma_{i}^{2}} \right)^{2} + \sigma_{L}^{2} \sum_{i=1}^{N} \frac{\mu_{i}^{2}}{\sigma_{i}^{2}} \right)}$$

$$= \frac{1}{x} \sqrt{2\pi}^{-N} \prod_{i=1}^{N} \sigma_{i}^{-1} e^{-\frac{1}{2} \left(\ln x - \sigma_{L}^{2} \sum_{i=1}^{N} \frac{\mu_{i}}{\sigma_{i}^{2}} \right)^{2} + \frac{1}{2} \sigma_{L}^{2} \left(\sum_{i=1}^{N} \frac{\mu_{i}}{\sigma_{i}^{2}} \right)^{2} - \frac{1}{2} \sum_{i=1}^{N} \frac{\mu_{i}^{2}}{\sigma_{i}^{2}}}{\sigma_{i}^{2}}$$

$$= \frac{1}{x} \sqrt{2\pi}^{-N} \prod_{i=1}^{N} \sigma_{i}^{-1} e^{-\frac{1}{2} \left(\ln x - \sigma_{L}^{2} \sum_{i=1}^{N} \frac{\mu_{i}}{\sigma_{i}^{2}} \right)^{2} - \frac{1}{2} \sum_{i=1}^{N} \frac{\mu_{i}^{2}}{\sigma_{i}^{2}}}{\sigma_{i}^{2}}$$

$$= \frac{1}{x} \sqrt{2\pi}^{-N} \prod_{i=1}^{N} \sigma_{i}^{-1} e^{-\frac{1}{2} \left(\ln x - \sigma_{L}^{2} \sum_{i=1}^{N} \frac{\mu_{i}}{\sigma_{i}^{2}} \right)^{2} - \frac{1}{2} \sum_{i=1}^{N} \frac{\mu_{i}^{2}}{\sigma_{i}^{2}}}{\sigma_{i}^{2}}$$

$$= \frac{1}{x} \sqrt{2\pi}^{-N} \prod_{i=1}^{N} \sigma_{i}^{-1} e^{-\frac{1}{2} \left(\ln x - \sigma_{L}^{2} \sum_{i=1}^{N} \frac{\mu_{i}}{\sigma_{i}^{2}} \right)^{2} - \frac{1}{2} \sum_{i=1}^{N} \frac{\mu_{i}^{2}}{\sigma_{i}^{2}}}{\sigma_{i}^{2}}$$

$$= \frac{1}{x} \sqrt{2\pi}^{-N} \prod_{i=1}^{N} \sigma_{i}^{-1} e^{-\frac{1}{2} \left(\ln x - \sigma_{L}^{2} \sum_{i=1}^{N} \frac{\mu_{i}}{\sigma_{i}^{2}} \right)^{2} - \frac{1}{2} \sum_{i=1}^{N} \frac{\mu_{i}}{\sigma_{i}^{2}}}{\sigma_{i}^{2}}$$

3.1.3.2 Scaled Lognormal Approximation

From the above approximation, it would seem that we can express the lower tail of the SLN distribution as a LN times constant distribution:

$$f(x) \stackrel{x \to 0}{\approx} \frac{\kappa_L}{\sqrt{2\pi}\sigma_L x} e^{-\frac{1}{2}\left(\frac{\ln x - \mu_L}{\sigma_L}\right)^2} \qquad F(x) \stackrel{x \to 0}{\approx} \kappa_L \Phi\left(\frac{\ln x - \mu_L}{\sigma_L}\right), \tag{3.11}$$

where:

$$\mu_{L} = \ln N + \sigma_{L}^{2} \sum_{i=1}^{N} \frac{\mu_{i}}{\sigma_{i}^{2}}, \quad \kappa_{L} = N \sigma_{L} \left(\prod_{i=1}^{N} \sigma_{i} \right)^{-1} \sqrt{2\pi^{1-N}} e^{\frac{1}{2} \sigma_{L}^{2} \left(\sum_{i=1}^{N} \frac{\mu_{i}}{\sigma_{i}^{2}} \right)^{2} - \frac{1}{2} \sum_{i=1}^{N} \frac{\mu_{i}^{2}}{\sigma_{i}^{2}}}.$$
 (3.12)

3.1.3.3 Best Lognormal Fit

Since the expression in (3.11) is a constant factor of a LN *pdf*, and assuming the development in Section 3.1.3.1 is valid, we present the following important result:

Hypothesis:

The SLN has a $blf \Phi\left(\frac{\ln x - \mu_L}{\sigma_L}\right)$ at 0, where the parameters are given by (3.9), (3.12).

3.1.3.4 Difficulties in Approximating the Lower Tail

It is important to note that, the LN distribution is very flat near zero; in fact, $\lim_{x\to 0} \frac{\partial^n}{\partial x^n} \Phi\left(\frac{\ln x - \mu}{\sigma}\right) = 0 \forall n = 1,2,.... \text{ The convolution is thus difficult to analyse, as it does not}$

fit in the framework of [55], nor is there a guarantee that the discrete convolution approximation performed in Section 3.1.3.1 is very good. Indeed, we find by simulations that the value of κ_L does not improve the approximation much, as will be seen in the next section.

3.1.3.5 Simulation of lower tail values

We find that in order to validate our asymptote, it is necessary to simulate the SLN cdf at very low values, which correspond to probabilities as low as 10^{-20} or 10^{-30} , which would require an inversely long simulation length, which is simply prohibitive. We can however make this simulation much more efficient by noticing that the sum of positive RV's is always larger than (or equal to) any of the individual RV's. In other words, we only need to generate low values of Y_i in order to obtain low values of the sum X. More specifically, to generate the portion of the cdf below a value SimLowVal, we need only generate values of Y_i from the distribution F_i , truncated on the interval [0, SimLowVal] and rescaled properly. We can thus obtain the sum distribution using conditional probabilities:

$$F(x) = P\left(\sum_{i=1}^{N} Y_i < x\right) = P\left(\sum_{i=1}^{N} Y_i < x \cap \{Y_i < x \forall i\}\right) = P\left(\sum_{i=1}^{N} (Y_i | Y_i < x) < x\right) \prod_{i=1}^{N} P(Y_i < x) \quad (3.13)$$

The algorithm for this simulation is presented as a flowchart in Figure 3-1. There is unfortunately no such simple algorithm for simulating high values of X.

In Section 3.1.3.6, we perform some such simulations. The simulations can be done in reasonable time if the number of summands N is not too high, which is why we focused on cases

of N = 2, and cases where the summands have different distributions and thus one tends to dominate. In Figures 3-2 – 3-4 we see broken pieces of a curve. Each piece corresponds to one simulation with a different value of SimLowVal. Tail probabilities of the order of 10^{-20} can be obtained using only millions of TRIALS (compare with 10^{20}). However, as N increases (say for 6 i.i.d. terms), the algorithm is no longer very efficient.

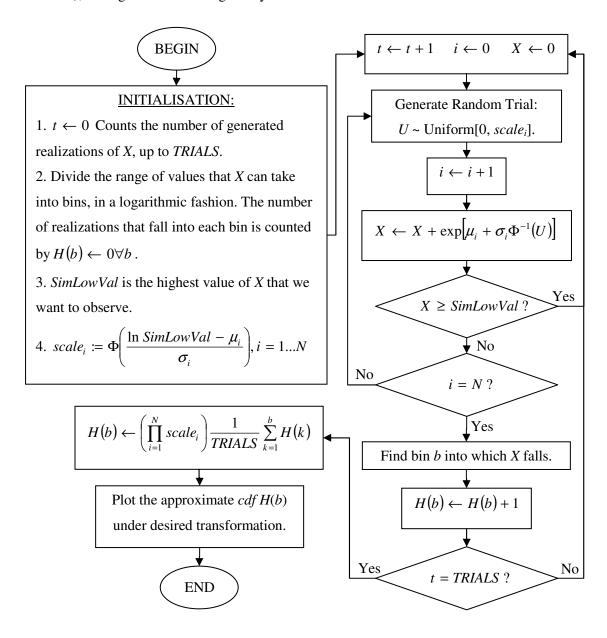


Figure 3-1 Algorithm that efficiently simulates the far lower tail of the SLN cdf.

3.1.3.6 Comparison with Simulations and Recent Methods

Using the method described in Section 3.1.3.5, we perform Monte-Carlo simulations of the lower tail of the SLN cdf curves, along with our theoretical asymptotes, and we also reproduce some curves from the LL, ZD and BR methods. We plot all these on LN probability paper. The simulations are done using values from Table 3-1 and are shown on Figures 3-2-3-4.

We use the values from Table I in [30] to obtain curves for the BR method for several *i.i.d.* cases. For the LL curves, we used the values from Table I in [27]. Note that in the fifth row of that Table, the log-variance should be 12dB, not 6dB. For the ZD curves, we used Equations (18)–(20) in [29], which are applicable to *i.i.d.* cases only, and where σ is in dB. The ZS curves are difficult to reproduce because the *cdf* is not known in closed form [28].

We also plot the theoretical asymptotes: one straight asymptote, the *blf* $\Phi\left(\frac{\ln x - \mu_U}{\sigma_U}\right)$, and one LN-times-constant asymptote for each tail (3.11).

We see that in all cases the BR and LL methods, and for most cases, the ZD method, give accurate results in the body of the distribution. However, it is evident from the Figures that each method gives a curve that tends to diverge from the *blf* of the SLN at either tail. We will show in Section 3.2.1 that this must always be the case. In particular, with the LL method, it could already be seen that there is this divergence in the tails, on LN paper [27].

The SLN converges instead to our proposed asymptotes. We see that the simulated curve tends to be sandwiched between these two curves. By virtue of Lemma 1, this would imply that the SLN cdf is bounded by $\Phi\left(\frac{\ln x - \mu_U}{\sigma_U}\right)$ from above and by (3.11) from below. Since these two

functions are constant multiples of each other, it would imply that there exists κ'_L , $\kappa_L \le \kappa'_L \le 1$ such that:

$$\lim_{x \to 0} \frac{\kappa_L'}{F(x)} \Phi\left(\frac{\ln x - \mu_L}{\sigma_L}\right) = 1. \tag{3.14}$$

This would be a much stronger result than simply a blf. However, we cannot prove it.

Figure	N	μ_i (dB)	$\sigma_i(dB)$
3-2 (A)	2	0	6
3-2 (B)	2	0	12
3-3	6	0	6, 8, 9, 10, 11, 12
3-4	6	-25, -15, -5, 5, 15, 25	12

Table 3-1 Simulation parameters for SLN cdfs in the lower tail.

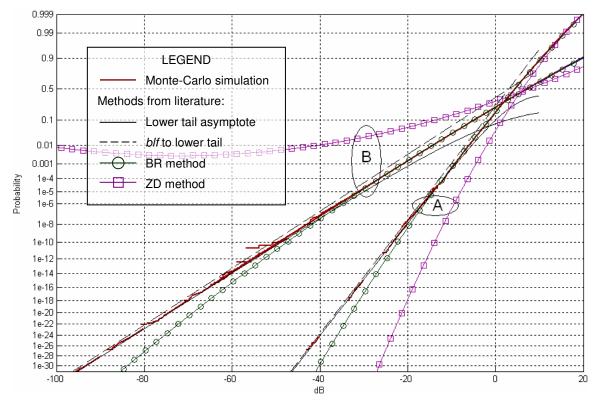


Figure 3-2 The *cdf* of the sum of two *i.i.d.* LN RV's – A: σ = 6dB and B: σ = 12dB.

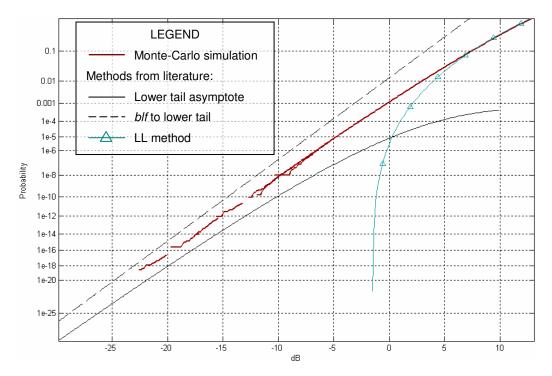


Figure 3-3 The cdf of the sum of lognormals with different variances.

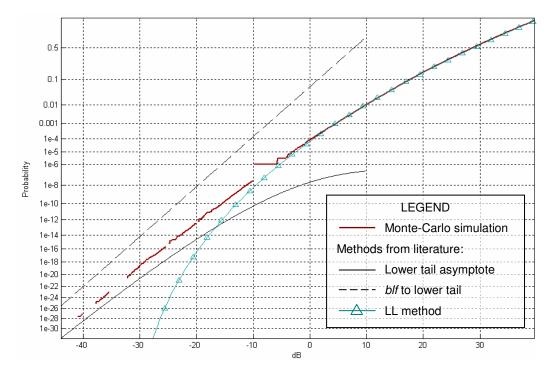


Figure 3-4 The cdf of the sum of lognormals with different means.

3.1.4 Upper Tail

Since the distribution in (2.24), (2.25) is a finite constant multiple of a LN, and Lemma 2 says that if a function is not the *blf* at a given tail of another function, then the ratio at that tail of those two functions must diverge. Consequently, the SLN distribution has $blf \Phi\left(\frac{\ln x - \mu_U}{\sigma_U}\right)$ at $+\infty$.

3.2 Potential Functions for Approximating the SLN Distribution

Having examined the existing methods for estimating the SLN distribution in Chapter 2, and in light of the results on the tail behaviours of the SLN distribution from Section 3.1, we now show that all the known parametric methods are inadequate in the sense that they necessarily diverge from the true SLN distribution in the tails. We then show how we can construct functions that do respect the tail constraints presented in Sections 3.1.3 and 3.1.4.

3.2.1 Inadequacy of Existing Functions in the Tails

asymptote at $l = \pm \infty$, except in some trivial cases when $F_{XX}(x)$ is LN.

Our simulation results suggest that the existing methods for finding the SLN distribution are not accurate in the far tails. The SLN distribution is shown to have a *blf* in the upper tail, and probably in the lower tail also. In this section, we prove that recent available methods all use distribution functions that cannot have a *blf* at either tail (except perhaps in the trivial LN case). Consequently, all these methods must inevitably give tail values that have infinite relative error. We examine four distributions that have recently been proposed as approximations to the SLN. For each of these we show that they cannot possibly give a good approximation at the tails. We do this formally by showing for all four methods that for the proposed *cdf* $F_{XX}(x)$, $x \ge 0$, for any $\sigma > 0$, μ , the result of the transformation $T: F_{XX}(x) \mapsto \tilde{F}_{XX}(x)$ cannot have $\frac{x - \mu}{\sigma}$ as an

We also show that only some of the proposed functions have the LN distribution as a special case, which should be a characteristic of an SLN approximation.

We use the constant $\lambda = 0.1 \ln 10$ to convert from decibel units.

3.2.1.1 Log-Shifted Gamma Function: Le-Ngoc - Lam (LL)

This approximating *cdf* has three parameters (α, β, δ) [27]:

$$F_{LL}(x) = \frac{1}{\Gamma(\alpha)} \gamma \left(\frac{\ln x}{\lambda \beta} - \frac{\delta}{\beta}, \alpha \right), x > e^{\delta \lambda}, \beta > 0, \alpha > 0,$$
 (3.15)

where $\gamma(x, \alpha) = \int_0^x e^{-t} t^{\alpha - 1} dt$ is the (lower) incomplete gamma function, and $\Gamma(\alpha)$ is the gamma function.

$$T: F_{LL}(x) \mapsto \widetilde{F}_{LL}(x) = \Phi^{-1}\left(\frac{1}{\Gamma(\alpha)} \gamma \left(\frac{x - \delta \lambda}{\lambda \beta}, \alpha\right)\right), x > \delta \lambda.$$
 (3.16)

The function is not even defined for arbitrarily low values.

In the upper tail, we use an approximation of the incomplete gamma function [56]:

$$\Gamma(\alpha) - \gamma(x, \alpha) \stackrel{x \to \infty}{\sim} x^{\alpha - 1} e^{-x} (1 + O(x^{-1})), \tag{3.17}$$

and an approximation to the inverse standard normal *cdf* [57]:

$$\Phi^{-1}(x) \stackrel{x \to 0}{\sim} - \sqrt{-2 \ln(2\sqrt{\pi}x\sqrt{-\ln(2x)})}. \tag{3.18}$$

Then:

$$\lim_{x \to +\infty} \tilde{F}_{LL}(x) - \frac{x - \mu}{\sigma} = \frac{\mu - \lambda \delta}{\sigma} - \lim_{x \to +\infty} \Phi^{-1} \left(\frac{1}{\Gamma(\alpha)} x^{\alpha - 1} e^{-x} \right) - \frac{\lambda \beta}{\sigma} x$$

$$= \lim_{x \to +\infty} \sqrt{2x} - \frac{\lambda \beta}{\sigma} x + \frac{\mu - \lambda \delta}{\sigma} = -\infty.$$
(3.19)

Hence, the LL distribution does not have a blf at either tail.

3.2.1.2 Three-Parameter Function: Beaulieu – Rajwani (BR)

This approximating cdf has three parameters (a_0, a_1, a_2) [30]:

$$F_{BR}(x) = \Phi\left(a_0 - a_1 x^{-\frac{a_2}{\lambda}}\right), x \ge 0, a_0 > 0, a_1 > 0, a_2 > 0,$$
(3.20)

$$T: F_{BR}(x) \mapsto \tilde{F}_{BR}(x) = a_0 - a_1 e^{-\frac{a_2}{\lambda}x}.$$
 (3.21)

Then:

$$\lim_{x \to l} \tilde{F}_{BR}(x) - \frac{x - \mu}{\sigma} = a_0 + \frac{\mu}{\sigma} - \lim_{x \to l} a_1 e^{-\frac{a_2}{\lambda}x} + \frac{x}{\sigma} = \pm \infty.$$
 (3.22)

Hence, the BR distribution does not have a blf at either tail.

In [29] it is mentioned that this form does not include the LN distribution as a special case. We can, however, obtain convergence of $F_{BR}(x) \to \Phi\left(\frac{\ln x - \mu}{\sigma}\right) \forall x$:

Let us consider μ , σ as constants, such that:

$$\mu = \frac{a_1 - a_0}{a_1 a_2}, \sigma = \frac{1}{a_1 a_2}.$$
 (3.23)

Then:

$$F_{BR}(x) = \Phi\left(\frac{-\frac{a_1 - a_0}{a_1 a_2} - \left(\frac{x^{-a_2} - 1}{a_2}\right)}{\frac{1}{a_1 a_2}}\right) = \Phi\left(\frac{\left(\frac{1 - x^{-a_2}}{a_2}\right) - \mu}{\sigma}\right). \tag{3.24}$$

Using an alternative definition of the natural logarithm:

$$\ln x = \lim_{a \to 0} \frac{1 - x^{-a}}{a}, x \ge 0 \tag{3.25}$$

we obtain:

$$F_{BR}(x) \xrightarrow{a_2 \to 0} \Phi\left(\frac{\ln x - \mu}{\sigma}\right), x \ge 0.$$
 (3.26)

Thus, in the limiting case:

$$a_2 \to 0, a_1 = (a_2 \sigma)^{-1}, a_0 = a_1 - \mu / \sigma,$$
 (3.27)

the BR distribution is LN.

3.2.1.3 Three-Parameter Function: Zhao - Ding (ZD)

This approximating *cdf* has three parameters (c_0, c_1, c_2) [29]:

$$F_{ZD}(x) = \Phi\left(c_0 + c_1 \frac{\ln x}{\lambda} + c_2 \left(\frac{\ln x}{\lambda}\right)^2\right), x \ge 0,$$
(3.28)

$$T: F_{ZD}(x) \mapsto \tilde{F}_{ZD}(x) = c_0 + \frac{c_1}{\lambda} x + \frac{c_2}{\lambda^2} x^2.$$
 (3.29)

Then, if $c_2 \neq 0$:

$$\lim_{x \to l} \tilde{F}_{ZD}(x) - \frac{x - \mu}{\sigma} = \lim_{x \to l} \frac{c_2}{\lambda^2} x^2 + \left(\frac{c_1}{\lambda} - \frac{1}{\sigma}\right) x + c_0 + \frac{\mu}{\sigma} = \pm \infty$$
 (3.30)

Hence, the ZD distribution does not have a \emph{blf} at either tail, except in the LN case $\,c_2=0\,.$

3.2.1.4 Pearson Type IV Function: Zhang – Song (ZS)

The proposed distribution for the SLN in dB is Pearson type IV, which has four parameters (u, m, d, v) [28]. Consequently, the *pdf* of the SLN in linear scale has distribution:

$$f_{ZS}(x) = \frac{k}{\lambda x} \left(1 + y^2 \right)^{-m} e^{-v \arctan y},$$

$$y = \frac{\ln x}{\lambda d} + \frac{u}{d}, m > \frac{1}{2}, x \ge 0,$$
(3.31)

and k is the constant chosen so that the pdf be valid and is found by numerical integration for each particular case. It would seem that the cdf is not readily available in closed form. However, we can still study its tail behaviour:

$$f_{ZS}(x) \stackrel{x \to 0}{\sim} k |d|^{2m} \lambda^{2m-1} x^{-1} |\ln x|^{-2m} e^{v\frac{\pi}{2} \operatorname{sgn}(d)}$$

$$= C_1 \frac{|\ln x|^{-2m}}{x}, C_1 > 0,$$
(3.32)

$$F_{ZS}(x) \stackrel{x \to 0}{\sim} = \int_0^x C_1 \frac{\left| \ln y \right|^{-2m}}{y} dy = C_1 \frac{\left| \ln x \right|^{1-2m}}{2m-1}. \tag{3.33}$$

Again, using the result (3.18):

$$T: F_{ZS}(x) \mapsto \tilde{F}_{ZS}(x) \xrightarrow{x \to -\infty} \Phi^{-1} \left(\frac{C_1}{2m-1} |x|^{1-2m} \right)$$

$$\xrightarrow{x \to -\infty} -\sqrt{(4m-2)} \sqrt{\ln(-x)} \Rightarrow \lim_{x \to -\infty} \tilde{F}_{ZD}(x) - \frac{x-\mu}{\sigma}$$

$$= -\lim_{x \to -\infty} \sqrt{(4m-2)} \sqrt{\ln(-x)} + \frac{x-\mu}{\sigma} = +\infty.$$
(3.34)

A similar analysis applies for the upper tail.

Hence, the ZS distribution does not have a *blf* at either tail.

3.2.2 Potential Function Candidates

If the existing methods use *cdf*'s that cannot behave properly in the tails, what then would functions that do look like? We give here a collection of three functions that verify at least some of the tail requirements given in Sections 3.1.3.2 and 2.2.2.

3.2.2.1 Power Lognormal

A PLN *cdf* has general form:

$$F_{PLN}(x) = \Phi^t \left(\frac{\ln x - m}{s}\right), x \ge 0 \quad t \ge 1, s > 0$$
 (3.35)

It has pdf:

$$f_{PLN}(x) = \frac{t}{\sqrt{2\pi}xs} \exp\left[-\frac{1}{2} \left(\frac{\ln x - m}{s}\right)^2\right] \Phi^{t-1}\left(\frac{\ln x - m}{s}\right), x \ge 0$$
 (3.36)

It has *blf* $\Phi\left(\frac{\ln x - m}{s}\right)$ at $+\infty$. Proof:

$$\lim_{x \to +\infty} \frac{1 - F_{PLN}(x)}{1 - \Phi\left(\frac{\ln x - m}{s}\right)} = \lim_{x \to +\infty} \frac{f_{PLN}(x)}{\frac{1}{\sqrt{2\pi}xs} \exp\left[-\frac{1}{2}\left(\frac{\ln x - m}{s}\right)^{2}\right]} = \lim_{x \to +\infty} t\Phi^{t-1}\left(\frac{\ln x - m}{s}\right) = t \quad (3.37)$$

By lemma 2, since the limit is finite (non-zero), the function is indeed the *blf*. This completes the proof.

The lower tail poses more of a problem, because it does not have a *blf* unless t = 1.

Proof:

Consider

$$T: F_{PLN}(x) \mapsto \widetilde{F}_{PLN}(x) = \Phi^{-1}\left(\Phi^{t}\left(\frac{x-m}{s}\right)\right). \tag{3.38}$$

Using the following approximation:

$$\Phi(x)^{x \to -\infty} \frac{1}{-\sqrt{2\pi}x} e^{-\frac{1}{2}x^2}$$
(3.39)

we find:

$$\widetilde{F}_{PLN}(x) \stackrel{x \to -\infty}{\sim} \Phi^{-1} \left(\frac{s^t e^{-\frac{t}{2} \left(\frac{x-m}{s} \right)^2}}{\sqrt{2\pi}^t (m-x)^t} \right). \tag{3.40}$$

We use the tail approximation (3.18) from [57]:

$$\widetilde{F}_{PLN}(x) \stackrel{x \to -\infty}{\sim} - \sqrt{-2 \ln \left(2\sqrt{\pi} \frac{s^t e^{-\frac{t}{2} \left(\frac{x-m}{s} \right)^2}}{\sqrt{2\pi}^t (m-x)^t} \sqrt{-\ln \left(2\frac{s^t e^{-\frac{t}{2} \left(\frac{x-m}{s} \right)^2}}{\sqrt{2\pi}^t (m-x)^t} \right)} \right)} \\
= - \sqrt{-2t \ln \left(2\sqrt{\pi} \frac{s}{\sqrt{2\pi}} \right) + t \left(\frac{x-m}{s} \right)^2 + 2t \ln(m-x)} \\
= - \sqrt{-2 \ln \left(\sqrt{-t \ln \left(2\frac{s}{\sqrt{2\pi}} \right) + t \ln(m-x) + \frac{t}{2} \left(\frac{x-m}{s} \right)^2} \right)} \\
= - \sqrt{-2t \ln \left(2\frac{s}{\sqrt{2\pi}} \right) + t \left(\frac{x-m}{s} \right)^2 - \ln \left(\frac{-t \ln \left(2\frac{s}{\sqrt{2\pi}} \right) + t \ln(m-x) + \frac{t(m-x)^{2(1-t)}}{2s^2} \right)}{(m-x)^{2t}} \right)} \\
\sim - \sqrt{t \left(\frac{x-m}{s} \right)^2 - 2 \ln \left(\frac{\sqrt{t} (m-x)^{(1-t)}}{\sqrt{2s}} \right)} \sim - \sqrt{t \left(\frac{x-m}{s} \right)^2 + 2(t-1) \ln(-x)}. \tag{3.41}$$

Now if $\Phi\left(\frac{\ln x - \mu}{\sigma}\right)$ is the *blf* at 0, then the following equation must hold:

$$\exists \mu, \exists \sigma : \lim_{x \to -\infty} \frac{x - \mu}{\sigma} + \sqrt{t \left(\frac{x - m}{s}\right)^2 + 2(t - 1)\ln(-x)} = 0$$

$$\Rightarrow \lim_{x \to -\infty} \left(\frac{x - \mu}{\sigma}\right)^2 + t \left(\frac{x - m}{s}\right)^2 + 2(t - 1)\ln(-x) + 2\left(\frac{x - \mu}{\sigma}\right)\sqrt{t \left(\frac{x - m}{s}\right)^2 + 2(t - 1)\ln(-x)} = 0$$

$$\Rightarrow \lim_{x \to -\infty} \left(\frac{x - \mu}{\sigma}\right)^2 + t \left(\frac{x - m}{s}\right)^2 + 2(t - 1)\ln(-x) - 2\sqrt{t} \left(\frac{x - \mu}{\sigma}\right)\left(\frac{x - m}{s}\right) = 0$$

$$\Rightarrow \lim_{x \to -\infty} Ax^2 + Bx + C + 2(t - 1)\ln(-x) = 0.$$
(3.42)

This has no solutions, for any constants A, B, C, unless t = 1. This proves that the PLN cdf does not have a blf at $-\infty$. However the lower tail does have a limiting slope:

$$\lim_{x \to -\infty} \frac{\partial}{\partial x} \widetilde{F}_{PLN}(x) = -\lim_{x \to -\infty} \frac{\partial}{\partial x} \sqrt{t \left(\frac{x-m}{s}\right)^2 + 2(t-1)\ln(-x)}$$

$$= -\lim_{x \to -\infty} \frac{2\frac{t}{s} \left(\frac{x-m}{s}\right) - \frac{2(t-1)}{x}}{2\sqrt{t \left(\frac{x-m}{s}\right)^2 + 2(t-1)\ln(-x)}} = \lim_{x \to -\infty} \frac{2\frac{t}{s} \left(\frac{x-m}{s}\right) - \frac{2(t-1)}{x}}{2x\sqrt{t \left(\frac{1}{s} - \frac{m}{sx}\right)^2 + 2(t-1)\frac{\ln(-x)}{x^2}}}$$

$$= \lim_{x \to -\infty} \frac{2\frac{t}{s} \left(\frac{x-m}{s}\right)}{2x\sqrt{t \left(\frac{x-m}{sx}\right)^2}} = \lim_{x \to -\infty} \frac{t \left(\frac{x-m}{s}\right)}{-s\sqrt{t} \left(\frac{m-x}{s}\right)} = \frac{\sqrt{t}}{s}.$$
(3.43)

We are thus interested in the PLN distribution because it has a *blf* in the upper tail, and at least a limiting slope in the lower tail. Also, [22] suggested that the SLN distribution could be written as the sum of products of LN distributions. A PLN distribution is a particular case of this, which makes it an interesting candidate.

3.2.2.2 A Four-Parameter Function

Since we have seen that the SLN has blf^*s at both tails, it would be interesting to create a function that also has the same blf^*s at both tails. To do this, we first construct a function r(x), such that:

1. r(x) is analytically smooth.

$$2. \quad \lim_{x \to -\infty} r(x) = 0.$$

$$3. \quad \lim_{x \to +\infty} r(x) - x = k \ .$$

One such function is:

$$r(x,k) = \ln(1 + e^k e^x)$$
 (3.44)

Now, in order to construct an approximation to the SLN *cdf*, we construct the following four-parameter (4P) function:

$$\tilde{F}_{4P}(x) = ax - r\left(cx, \frac{d}{b}\right) - \ln b = ax - \ln(b + de^{cx}), a > c > 0,$$
(3.45)

which has tail behaviour:

$$\tilde{F}_{4P}(x) \stackrel{x \to +\infty}{\approx} (a-c)x - \ln d \quad \text{and} \quad \tilde{F}_{4P}(x) \stackrel{x \to -\infty}{\approx} ax - \ln b.$$
 (3.46)

Consequently,

$$T^{-1}: \tilde{F}_{4P}(x) \mapsto F_{4P}(x) = ax - r\left(cx, \frac{d}{b}\right) - \ln b = \Phi\left(\ln\left(\frac{x^a}{b + dx^c}\right)\right), a > c > 0$$
 (3.47)

is our proposed function 4P for approximating SLN. It has $blf \Phi((a-c) \ln x - \ln d)$ at $+\infty$, and $blf \Phi(a \ln x - \ln b)$ at $-\infty$.

3.3 Obtaining the Parameters of the Approximating Functions

3.3.1 Power Lognormal Distribution

We wish to fit the PLN *cdf* to that of the SLN as well as possible. We know the *blf* of both the upper and lower tail of the SLN, and we can also find its moments analytically, whereas the PLN

distribution has three free parameters only. We choose to match the slopes of both tails first, because an error in the slope of the tail asymptote will become infinitely more significant than an error in the offset as we approach the limit. We also found that matching the first moment of the PLN cdf with that of the SLN was a best-effort method to get a good curve match. It is fortunate that we do not have to solve for all three parameters simultaneously, but rather we first find s, then, using s, find t, finally we find m using a numerical integral involving s and t.

3.3.1.1 Matching the Tail Slopes

We first find the value of s. The upper tail slope is given by the blf at $+\infty$ (see Sections 3.1.4 and (2.21)) and is matched with the blf at $+\infty$ of the PLN, given in Section 3.2.2.1:

$$s = \sigma_U = \max_i \sigma_i . (3.48)$$

From the value of s we can find t. The lower tail slope is given by the blf at $-\infty$ (3.9) and is matched with (3.43):

$$\frac{\sqrt{t}}{s} = \frac{\sqrt{t}}{\sigma_U} = \frac{1}{\sigma_L} = \left(\sum_{i=1}^N \frac{1}{\sigma_i^2}\right)^{\frac{1}{2}} \Rightarrow t = \sigma_U^2 \sum_{i=1}^N \frac{1}{\sigma_i^2} = \sum_{i=1}^N \frac{\max_i \sigma_i^2}{\sigma_i^2} \ge N,$$
 (3.49)

and t = N if and only if all variances are equal.

3.3.1.2 Moment-Matching

The moments of the of the PLN distribution with pdf (3.36) are:

$$E[X^{n}] = \frac{t}{\sqrt{2\pi}s} \int_{0}^{+\infty} x^{n} \exp\left[-\frac{1}{2} \left(\frac{\ln x - m}{s}\right)^{2}\right] \Phi^{t-1}\left(\frac{\ln x - m}{s}\right) \frac{dx}{x}$$

$$= \frac{t}{\sqrt{2\pi}s} \int_{-\infty}^{+\infty} e^{xn} \exp\left[-\frac{1}{2} \left(\frac{x - m}{s}\right)^{2}\right] \Phi^{t-1}\left(\frac{x - m}{s}\right) dx$$

$$= \frac{t}{\sqrt{2\pi}s} e^{mn} \int_{-\infty}^{+\infty} e^{xn} \exp\left[-\frac{1}{2} \left(\frac{x}{s}\right)^{2}\right] \Phi^{t-1}\left(\frac{x}{s}\right) dx$$

$$= e^{mn} \Lambda(sn, t), \quad \Lambda(sn, t) = \frac{t}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} e^{xsn - \frac{1}{2}x^{2}} \Phi^{t-1}(x) dx,$$

$$(3.50)$$

where the integral Λ needs to be calculated numerically as a function of s and t found in Section 3.3.1.1. This is essentially the same expression as found in [58], where the PLN distribution is defined slightly differently. Notice that the parameter m comes out of the integral, which allows us to solve for m, using the first moment:

$$m = \ln E[X] - \ln \Lambda(s,t) = \ln \sum_{i=1}^{N} e^{\mu_i + \frac{1}{2}\sigma_i^2} - \ln \Lambda(s,t)$$
 (3.51)

Although our method requires numerical integration at this point, it still remains much simpler than other numerical methods, for the following reasons:

- The integration needs only be performed once, not iteratively for all summands as in SY
 [37] or LL [27].
- 2. The integral is very well-behaved and can be calculated directly using a Riemann sum, unlike some of the numerical methods used in other SLN methods [35].

We calculate the integral by summing it over the interval [-25, 25] with step 0.01.

We can partially verify our method since we do have a closed-form expression for the moments of the PLN distribution when t = 2:

$$\Lambda(sn,2) = \frac{2}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} e^{xsn - \frac{1}{2}x^2} \Phi(x) dx$$

$$= \frac{2}{2\pi} \int_{-\infty}^{+\infty} e^{xsn - \frac{1}{2}x^2} \int_{-\infty}^{x} e^{-\frac{1}{2}t^2} dt dx = \frac{1}{\pi} \int_{-\infty}^{+\infty} \int_{-\infty}^{x} e^{xsn - \frac{1}{2}x^2 - \frac{1}{2}t^2} dt dx.$$
(3.52)

Now make the change of variables:

$$\begin{cases} a = x + t \\ b = x - t \end{cases} \Rightarrow \begin{cases} x = \frac{a + b}{2} \\ t = \frac{a - b}{2} \end{cases} \Rightarrow J = \frac{\partial(x, t)}{\partial(a, b)} = -\frac{1}{2} \Rightarrow |J| = \frac{1}{2}.$$
 (3.53)

$$E\left[X^{n} \middle| t = 2\right] = e^{mn} \Lambda(sn,2) = \frac{e^{mn}}{2\pi} \int_{-\infty}^{+\infty} \int_{0}^{+\infty} e^{-\frac{1}{4}(a^{2} + b^{2}) + \frac{sn}{2}(a+b)} dbda$$

$$= \frac{e^{mn}}{2\pi} \int_{-\infty}^{+\infty} e^{-\frac{1}{4}x^{2} + \frac{sn}{2}x} dx \int_{0}^{+\infty} e^{-\frac{1}{4}x^{2} + \frac{sn}{2}x} dx = \frac{e^{mn}}{\sqrt{\pi}} \sqrt{2\pi} \int_{0}^{+\infty} e^{-\frac{1}{2}x^{2} + \frac{sn}{\sqrt{2}}x} dx = 2\Phi\left(\frac{sn}{\sqrt{2}}\right) e^{mn + \frac{1}{2}s^{2}n^{2}}.$$
(3.54)

We verify this by comparing with numerical values in Table 3-2, where we see that the Riemann sum and the theoretical values coincide perfectly. However, we were not able to do a separation of variables for t higher than 2, so the general closed-form expression for the moments of a PLN distribution is to our knowledge unknown.

sn	$E[X^n t = 2, m = 0] = 2\Phi\left(\frac{sn}{\sqrt{2}}\right)e^{\frac{1}{2}s^2n^2}$	Numerical $\Lambda(sn,2)$
1.38155	4.3405163407083 <u>4</u>	4.3405163407083 <u>2</u>
2.76310	88.661408234589 <u>91</u>	88.661408234589 <u>83</u>
5.52620	8559626.6550080 <u>20</u>	8559626.6550080 <u>11</u>
11.0524	$6.7113558551748\underline{30}\cdot10^{26}$	$6.7113558551748\underline{11}\cdot10^{26}$
22.1048	$2.53\underline{6010695502967} \cdot 10^{106}$	$2.53\underline{1281918731208} \cdot 10^{106}$

Table 3-2 Analytical and numerical calculations of moments of the PLN distribution when the power is t = 2, with differences underlined.

3.3.2 Four-Parameter Function

Based on the *blf*'s given in Section 3.2.2.2, the system of equations for matching asymptotes is:

$$\frac{1}{\sigma_U} = a - c, \quad \frac{\mu_U}{\sigma_U} = \ln d, \quad \frac{1}{\sigma_L} = a, \quad \frac{\mu_L}{\sigma_L} = \ln b. \tag{3.55}$$

This is directly solvable:

$$a = \frac{1}{\sigma_L}, \quad b = e^{\frac{\mu_L}{\sigma_L}}, \quad c = \frac{1}{\sigma_L} - \frac{1}{\sigma_U}, \quad d = e^{\frac{\mu_U}{\sigma_U}},$$
 (3.56)

where the parameters are given by equations (2.21), (2.22), (3.9), and (3.12). After substitutions, we obtain the following closed-form expression in terms of the parameters of the summands:

$$F_{4P}(x) = \Phi\left(\frac{\ln x - \mu_U}{\sigma_U} - \ln\left(1 + e^{\frac{\ln x - \mu_U}{\sigma_U} - \frac{\ln x - \mu_L}{\sigma_L}}\right)\right). \tag{3.57}$$

3.3.3 Evaluation of Closed-Form Methods against Simulations

In this section, we evaluate the performance of the two SLN methods we developed against all other known closed-form methods. The known closed-form methods are:

- 1. Moment-Matching by the FW method using (2.16)–(2.19). We only consider matching the first and second moment, as these give the best results overall [40].
- 2. The lower bound expression in (2.26), which is an upper bound to the *cdf*.
- 3. The upper bound expression in (2.28), which is a lower bound to the *cdf*.
- 4. The AG method given by (2.30), (2.33).

We also plot our two new methods: PLN and 4P. While our PLN method is not strictly closed-form, the numerical integration required is minimal (and there is hope that the numerical integral in (3.50) can some day be found in closed form, at least for integer *t*). We run Monte-Carlo

simulations to validate all these methods. We have chosen four cases, described in Table 3-3.

The results are plotted on LN paper in Figures 3-5-3-8.

Figure	N	μ_i (natural units)	$\sigma_i\left(\mathbf{dB}\right)$
3-5	6	0	12
3-6	30	0	6
3-7	12	0, 1, 2, 3, 4, 5 6, 7, 8, 9, 10, 11	8
3-8	36	5, 5, 5, 5, 5, 5, 4, 4, 4, 4, 4, 4, 3, 3, 3, 3, 3, 3, 2, 2, 2, 2, 2, 2, 1, 1, 1, 1, 1, 1, 0, 0, 0, 0, 0, 0	12, 12, 12, 12, 12, 12, 11, 11, 11, 11, 11, 11, 10, 10, 10, 10, 10, 10, 9, 9, 9, 9, 9, 9, 8, 8, 8, 8, 8, 8, 7, 7, 7, 7, 7, 7

Table 3-3 Simulation parameters for sums of lognormals.

From the simulation results, we make the following observations:

- 1. The FW method performs well in approximating the upper tail.
- 2. The true SLN *cdf* is sandwiched by the two bounds (2.26) and (2.28).
- 3. The AG method gives a rather loose bound, as predicted in Section 2.2.3.2.
- 4. Our PLN method tends to perform better than all other methods, particularly in the complicated case of Figure 3-8. In particular, it performs better than the upper and lower bounds, which are also of the form of a product of *N* LN *cdf* s.
- 5. Our 4P method does not give a very close match to the real SLN *cdf*, particularly in the body. However, with additional degrees of freedom, we believe that we can make the curve move in the body so as to fit the entire curve well. Also, we conjecture that the 4P *cdf* is an upper bound to the SLN *cdf*.

In light of these observations, we choose both the FW method and our PLN method as candidates to approximate the interference distribution. The FW method is commonly used in similar interference analysis work [9], [13], [16], and is thus a good benchmark for comparison. Both methods are simple to implement, as will be seen in Chapter 4.

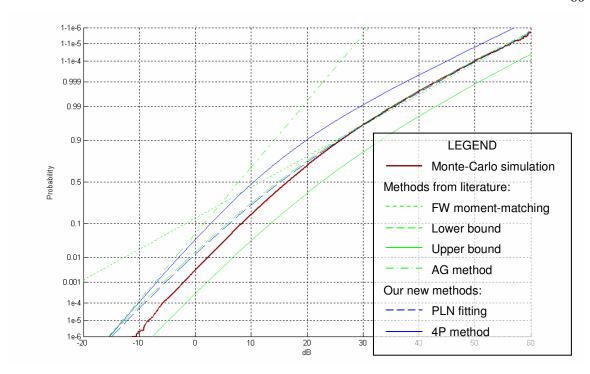


Figure 3-5 The *cdf* of the sum of 6 *i.i.d.* lognormals, and closed-form methods.

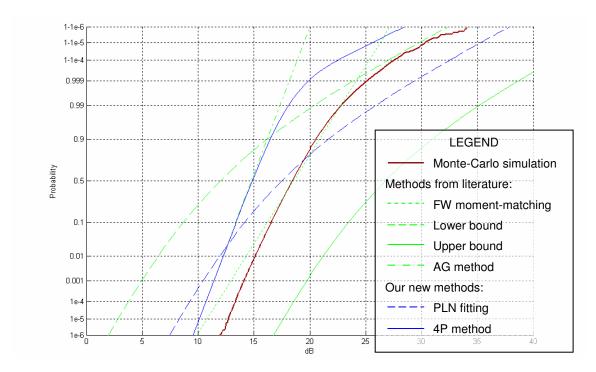


Figure 3-6 The *cdf* of the sum of 30 *i.i.d.* lognormals, and closed-form methods.

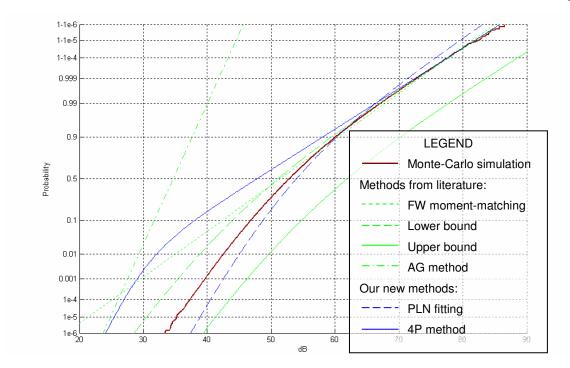


Figure 3-7 The *cdf* of the sum of 12 lognormals with different means, and closed-form methods.

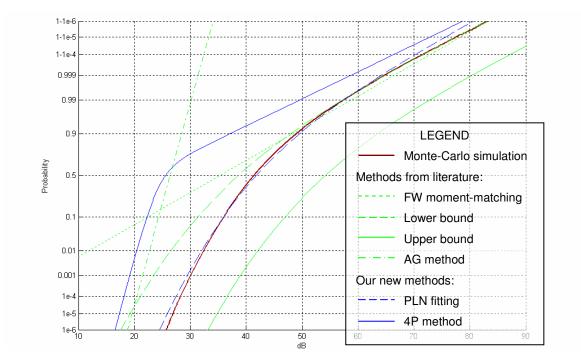


Figure 3-8 The cdf of the sum of 36 lognormals with different means and variances, and closed-form methods.

Chapter 4

Statistical Analysis of Total Interference Power

In cellular communications, the purpose of the mathematical problem of the SLN is to find the distribution of the total interference power *I* experienced by the DMU. We would like to do this without simulations, preferably with as little numerical methods as possible, ideally in simple closed form. We will use two different methods to find this distribution: one is the well-known FW moment-matching method [26], which is the simplest one and is entirely closed form. We found that no other closed-from method really existed that we could use with good results, which is the reason for our work in chapter 3. The second method we use is the PLN approximation, where we apply the method developed in Section 3.3.1 to the cellular interference problem.

In this chapter, we first show that the interference I_j coming from one cell can be well-modelled by a Bernoulli Lognormal (BLN) RV. The sum of BLN is a more difficult problem, so we need to modify our SLN methods slightly. For both methods, we devise an exact method that is a bit long, as well as handy approximations that turn out to be good and can be essentially calculated with a scientific calculator and the Tables provided in this chapter. In order to give credibility to our results we verify our solutions by Monte-Carlo simulations, which is sometimes lacking in literature [1], [3]. We find that our methods tend to predict the distribution well in most cases, but not always, clearly indicating that more work needs to be done in the SLN problem. We believe that a better method can be found by pursuing the direction of Sections 3.2.2 and 3.3, which looks promising in that both tails could be fitted using closed-form equations.

4.1 Simulation Algorithm

Figures 4-2 and 4-3 give a flowchart of the algorithm that simulates interference in a cellular system according to the model given above. We used this algorithm to create our simulated

curves. It is necessary to generate MU positions (both DMU and IMU) according to a spatially uniform distribution within a hexagonal cell of unit side. Also the positions of all the IFS's are needed. We can represent these as complex numbers, and then distances can be calculated as magnitudes of differences of positions. Figure 4-1 gives an example of a layout with normalised distances and mapped onto a complex plane.

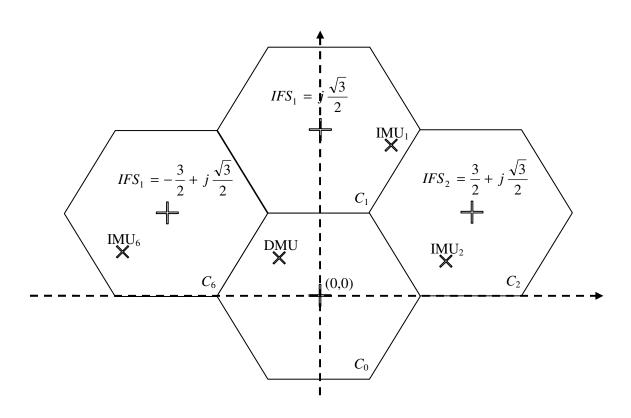


Figure 4-1 Normalised MU and FS positions in a complex plane.

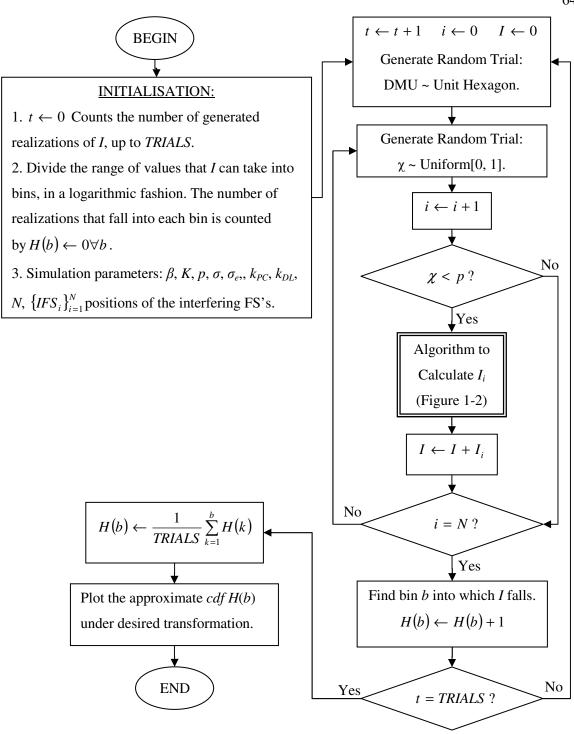


Figure 4-2 Algorithm for simulating interference in cellular systems.

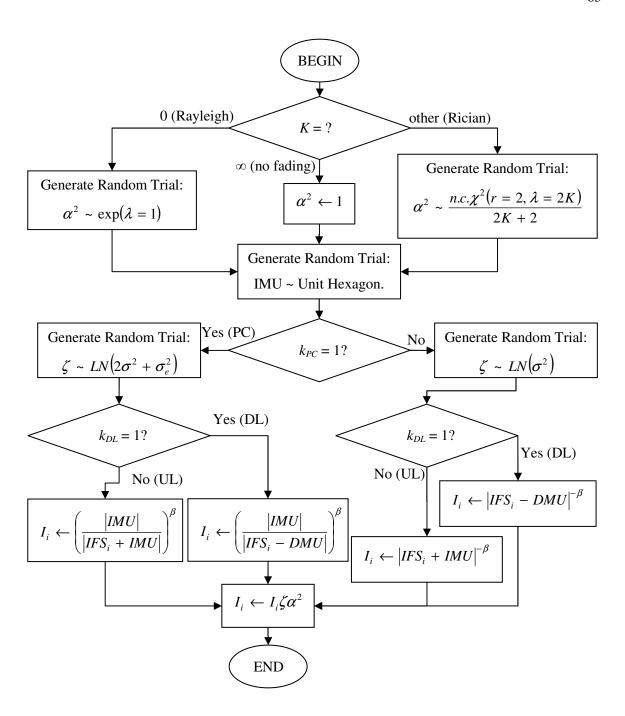


Figure 4-3 Algorithm to calculate I_i .

4.2 Distribution of the Interference Power from One Interferer

We first want to see what is the (approximate) distribution of the individual terms I_i .

4.2.1 Separating the Pathloss from all Other Factors

An important insight is that the pathloss is independent of shadowing, fading, activity, and in the DL, of transmit power due to PC. We can rewrite (1.1) as:

$$Q_i = \zeta_i \chi_i \alpha_i^2 P_i \Rightarrow I = \sum_{i=1}^N R_i Q_i . \tag{4.1}$$

The case of UL – PC requires a slight modification, as the transmit power and pathloss both depend simultaneously on the IMU position. Hence we can essentially separate the geometry of the problem from all other considerations, and can tabulate coefficients that are a function only of the location of the cell and the pathloss exponent. The coefficients are found by numerical integration, and here is the beauty of the method, need not be recalculated ever again, and do not change with other system parameters. This is the approach proposed in [3]. We found in literature [1], [6], [7] that what is often attempted is to integrate the pathloss distribution together with shadowing, fading, etc., but we feel this is "reinventing the wheel" at each time, if we only want the moments of the resulting RV. It is much more convenient to be able to treat the geometry of the problem separately, tabulate it, and then work on the problem at hand without going into any layout considerations. This is a core idea of this chapter, and leads naturally to our goal of obtaining simple, closed-form solutions to the interference distribution.

This separation of geometry is also important because all other factors: shadowing, fading, transmit power and activity are *i.i.d.* across the cells, whereas the pathlosses are both statistically different, and also, in DL, correlated. For all these reasons, we believe that the separation is a judicious one.

4.2.2 Statistical Moments of the Inverse Pathloss

The moments of the individual inverse pathlosses need to be calculated by numerical integration, though analytical approximations exist [3], [5], [10]:

$$\kappa_i(n\beta) = E[R_i^n] = \frac{2}{3\sqrt{3}} \iint_{(x,y)\in C_0} \sqrt{(x-x_i)^2 + (y-y_i)^2} dxdy.$$
 (4.2)

Here C_0 is the cell under consideration, enclosed by a unit hexagon centered at the origin, and (x_i, y_i) are the coordinates of IFS i. Figure 4-1 illustrates a possible layout. We tabulate these values for all cells in a traditional three-tier cellular layout (Figure 4-4). We see in the Figure that there are five different classes of cells, where cells in the same class are located in the same manner with respect to the cell under consideration C_0 . Thus we need to calculate κ coefficients for only five different cells (Tables 4-1 and 4-2).

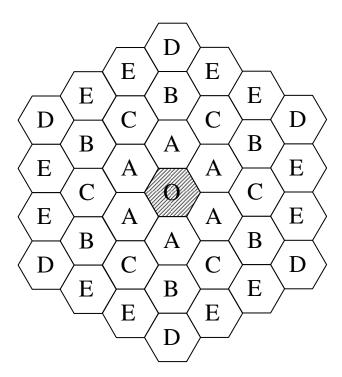


Figure 4-4 Three tiers of cells with five different cell positions (classes A-E).

β	Class A	Class B	Class C	Class D	Class E
3	0.2798	0.02609	0.04130	0.007384	0.01088
3.2	0.2645	0.02058	0.03365	0.005336	0.008072
3.4	0.2508	0.01624	0.02745	0.003858	0.005993
3.5	0.2446	0.01443	0.02480	0.003281	0.005165
3.6	0.2387	0.01283	0.02242	0.002790	0.004452
3.8	0.2279	0.01014	0.01832	0.002018	0.003308
4	0.2184	0.008023	0.01499	0.001461	0.002459
4.2	0.2099	0.006351	0.01227	0.001057	0.001829
4.4	0.2024	0.005031	0.01006	0.0007655	0.001361
4.5	0.1990	0.004479	0.009108	0.0006515	0.001174
4.6	0.1958	0.003988	0.008250	0.0005545	0.001013
4.8	0.1900	0.003164	0.006774	0.0004018	0.0007540
5	0.1849	0.002512	0.005567	0.0002912	0.0005617

Table 4-1 $\kappa_i(\beta)$ function tabulated for all cell classes and pathloss exponents.

β	Class A	Class B	Class C	Class D	Class E
3	0.1680	0.0007999	0.002113	0.00005849	0.0001296
3.2	0.1645	0.0005084	0.001443	0.00003085	0.00007224
3.4	0.1624	0.0003240	0.0009882	0.00001629	0.00004035
3.5	0.1618	0.0002590	0.0008187	0.00001184	0.00003017
3.6	0.1616	0.0002070	0.0006788	0.000008608	0.00002257
3.8	0.1620	0.0001325	0.0004676	0.000004555	0.00001264
4	0.1634	0.00008508	0.0003230	0.000002413	0.000007090
4.2	0.1659	0.00005473	0.0002237	0.000001280	0.000003982
4.4	0.1693	0.00003529	0.0001553	0.0000006796	0.000002240
4.5	0.1714	0.00002836	0.0001295	0.0000004954	0.000001681
4.6	0.1737	0.00002280	0.0001081	0.0000003612	0.000001262
4.8	0.1790	0.00001476	0.00007536	0.0000001922	0.0000007116
5	0.1853	0.000009576	0.00005267	0.0000001024	0.0000004019

Table 4-2 $\kappa_i(2\beta)$ function tabulated for all cell classes and pathloss exponents.

4.2.3 Statistical Moments of all Other Factors

Since the components of Q_i are independent, and all Q_i are i.i.d., we may define the *B*-coefficients as follows:

$$B_n = E[Q_i^n] = E[\zeta_i^n] E[\alpha_i^{2n}] E[\chi_i^n] E[P_i^n], \forall i.$$
(4.3)

If there is no PC: $P_i = 1, \forall i$.

With PC, the power transmitted will be proportional to the pathloss between the IMU i and its FS. The position of the IMU has uniform distribution over its cell area C_i . Now the cell area is a regular hexagon of unit side, which it is conventional to approximate by a circle of radius $\rho \leq 1$. We define the mean pathloss between an IMU i and its own IFS to be L_i :

$$L_i \sim f_L(x) = \frac{2}{\beta \rho^2} x^{\frac{2}{\beta} - 1}, x \in [0, \rho^{\beta}],$$
 (4.4)

$$E[L_i^n] = \frac{2}{2 + n\beta} \rho^{n\beta}. \tag{4.5}$$

In [1], [3], ρ is assumed 1. We found, through numerical integration, that $\rho = 0.915$ gives a good approximation for the first two moments and use this value in our analysis. A value of $3^{\frac{3}{4}}/\sqrt{2\pi} \cong 0.9094$ is used in [5] and [10].

The transmit power will also be affected by LN shadowing, for which the PC mechanism will attempt to compensate. We model this by a LN variable with statistics $(0, \sigma^2 + \sigma_e^2)$, as done in [1]. Finally, with PC, the moments of the transmit power are:

$$E[P_i^n] = \frac{2}{2 + n\beta} \rho^{n\beta} e^{\frac{n^2}{2} \left(\sigma^2 + \sigma_e^2\right)}. \tag{4.6}$$

Consequently:

$$B_{n} = pE\left[\alpha_{i}^{2n}\right]e^{\frac{n^{2}}{2}\sigma^{2}}\left[\frac{2}{2+n\beta}\rho^{n\beta}e^{\frac{n^{2}}{2}\left(\sigma^{2}+\sigma_{e}^{2}\right)}\right]^{k_{PC}}.$$
(4.7)

We normalise the fading power such that the power mean $E[\alpha_i^2] = 1$. Then the second moment of the power is known [59]:

$$E\left[\alpha_i^4\right] = 2 - \left(1 + K^{-1}\right)^{-2}. (4.8)$$

Note that when there is no fading: $K = \infty$, $E[\alpha_i^4] = 1$. For Rayleigh fading, $K = 0 \Rightarrow E[\alpha_i^4] = 2$. We then obtain:

$$B_{1} = p e^{\frac{1}{2}\sigma^{2}} \left[\frac{2}{2+\beta} \rho^{\beta} e^{\frac{1}{2}(\sigma^{2} + \sigma_{\epsilon}^{2})} \right]^{k_{PC}}, \tag{4.9}$$

$$B_2 = p \left(2 - \left(1 + K^{-1} \right)^{-2} \right) e^{2\sigma^2} \left[\frac{1}{1 + \beta} \rho^{2\beta} e^{2\sigma^2 + 2\sigma_e^2} \right]^{k_{PC}}. \tag{4.10}$$

We can write, for all cases except UL - PC, that the pathloss is independent of all other factors, which leads to:

$$E[I_i^n] = E[R_i^n Q_i^n] = E[R_i^n] E[Q_i^n] = \kappa_i(n\beta) B_n \qquad [k_{DL}, k_{PC}] \neq [0,1].$$
 (4.11)

However in the case of UL-PC, R_i and Q_i are not independent, and a modified method needs to be developed.

β	Class A	Class B	Class C	Class D	Class E
3	0.2799	0.02609	0.04130	0.007384	0.01088
3.2	0.2645	0.02058	0.03365	0.005336	0.008072
3.4	0.2508	0.01624	0.02745	0.003858	0.005993
3.5	0.2446	0.01443	0.02480	0.003281	0.005165
3.6	0.2387	0.01283	0.02242	0.002790	0.004452
3.8	0.2279	0.01014	0.01832	0.002018	0.003308
4	0.2184	0.008023	0.01499	0.001461	0.002459
4.2	0.2099	0.006351	0.01227	0.001057	0.001829
4.4	0.2024	0.005031	0.01006	0.0007655	0.001361
4.5	0.1990	0.004479	0.009108	0.0006515	0.001174
4.6	0.1958	0.003988	0.008250	0.0005545	0.001013
4.8	0.1900	0.003164	0.006774	0.0004018	0.000754
5	0.1849	0.002512	0.005567	0.0002912	0.0005617

Table 4-3 $\tilde{\kappa}_i(\beta)$ function tabulated for all cell classes and pathloss exponents.

β	Class A	Class B	Class C	Class D	Class E
3	0.1680	0.0007999	0.002113	0.00005849	0.0001296
3.2	0.1645	0.0005084	0.001443	0.00003085	0.00007224
3.4	0.1624	0.0003240	0.0009882	0.00001629	0.00004035
3.5	0.1618	0.0002589	0.0008187	0.00001184	0.00003017
3.6	0.1616	0.0002070	0.0006788	0.000008608	0.00002257
3.8	0.1620	0.0001325	0.0004676	0.000004555	0.00001264
4	0.1634	0.00008508	0.0003230	0.000002413	0.00000709
4.2	0.1659	0.00005473	0.0002237	0.00000128	0.000003982
4.4	0.1693	0.00003529	0.0001553	0.0000006796	0.00000224
4.5	0.1714	0.00002836	0.0001295	0.0000004954	0.000001681
4.6	0.1737	0.00002280	0.0001081	0.0000003612	0.000001262
4.8	0.1790	0.00001476	0.00007536	0.0000001922	0.0000007116
5	0.1853	0.000009576	0.00005267	0.0000001024	0.0000004019

Table 4-4 $\tilde{\kappa}_i(2\beta)$ function tabulated for all cell classes and pathloss exponents.

4.2.4 The Special Case of Uplink with Power Control

In the UL – PC case, the pathloss between the IMU and the FS of cell C_0 and the pathloss between the IMU and its own FS are two correlated RV's. We modify our method as follows:

$$\widetilde{\kappa}_{i}(n\beta) = E\left[R_{i}^{n}L_{i}^{n}\right] = \frac{2}{3\sqrt{3}} \iint_{(x,y)\in C_{0}} \sqrt{\left((x-x_{i})^{2} + (y-y_{i})^{2}\right)\left(x_{i}^{2} + y_{i}^{2}\right)^{-n\beta}} dxdy . \tag{4.12}$$

Here the integral is performed in the same manner as (4.2). These values are tabulated in Tables 4-3 and 4-4. We also define:

$$\widetilde{B}_{n} = E \left[\frac{Q_{i}^{n}}{L_{i}^{n}} \right] = E \left[\zeta_{i}^{n} \right] E \left[\alpha_{i}^{2n} \right] E \left[\chi_{i}^{n} \right] E \left[\frac{P_{i}^{n}}{L_{i}^{n}} \right], \forall i.$$
(4.13)

$$\widetilde{B}_n = pE\left[\alpha_i^{2n}\right] e^{\frac{n^2}{2}\left(2\sigma^2 + \sigma_e^2\right)}. \tag{4.14}$$

$$\widetilde{B}_1 = pe^{\sigma^2 + \frac{1}{2}\sigma_e^2}. \tag{4.15}$$

$$\widetilde{B}_2 = p \left(2 - \left(1 + K^{-1} \right)^{-2} \right) e^{4\sigma^2 + 2\sigma_e^2}. \tag{4.16}$$

Then the moments can be expressed as:

$$E\left[I_i^n\right] = E\left[R_i^n Q_i^n\right] = E\left[R_i^n L_i^n\right] E\left[\frac{Q_i^n}{L_i^n}\right] = \tilde{\kappa}_i(n\beta)\tilde{B}_n \qquad \left[k_{DL}, k_{PC}\right] = \left[0,1\right]. \tag{4.17}$$

4.2.5 Bernoulli Lognormal Approximation to Single Interferer Power Distribution

The signal power coming from a single interferer, I_j , can be modeled approximately by a BLN RV, with cdf:

$$F_i(x) \cong \Phi\left(\frac{\ln x - a_i}{b_i}\right) p + (1 - p), x > 0.$$
 (4.18)

4.2.5.1 Support in Literature

We model the interference from one user as a product of a LN shadowing component, a fading (Rayleigh, Rician or none) component, a pathloss component that assumes some spatial distribution of the MU. Also, if there is PC, the transmit power is not constant, but follows a distribution that is LN (compensating for shadowing and also because of LN PC error), multiplied by a quantity proportional to the pathloss (so as to compensate for it) between a randomly located IMU and its FS. The product of all these is approximated as a LN RV. Furthermore, user activity is modelled as Bernoulli, thus the interference power from a given interferer is modelled as BLN. This is the same approximation as in [1].

4.2.5.2 Theoretical Basis

It is not accidental that the distribution of the product of many RV's approaches a LN distribution, particularly if some of the terms are already LN, with high variance. First of all, the product of LN RV's is also LN. Secondly, by virtue of the Central Limit Theorem for products, the (properly normalised) product of many RV's with square-integrable *pdf* s will approach a LN RV in distribution [60]. The Bernoulli RV does not satisfy the condition and thus cannot be absorbed in the LN approximation. Indeed, multiplying any number of Bernoulli RV's can only result in another Bernoulli RV. The Central Limit Theorem for products works quite well in our case: though only a few RV's are multiplied, there are already among them LN RV's with high variance: the shadowing (6-12 dB) and, if there is PC, another shadowing factor due to the power compensation. Thus we expect the approximation to be good (and get better, with higher shadowing).

4.2.5.3 Fitting the Distribution by Moment-Matching

The moments of I_i can be matched with those of the approximating distribution $F_i(x)$:

$$E[I_i^n] = p \cdot e^{na_i + \frac{n^2}{2}b_i^2}. \tag{4.19}$$

Inverting (4.19), we obtain:

$$a_{i} = 2 \ln E[I_{i}] - \frac{1}{2} \ln E[I_{i}^{2}] - \frac{3}{2} \ln p$$

$$= \begin{cases} 2 \ln \kappa_{i}(\beta) - \frac{1}{2} \ln \kappa_{i}(2\beta) - \frac{3}{2} \ln p - \frac{1}{2} \ln(2 - (1 + K^{-1})^{-2}) \\ + k_{PC} [\beta \ln \rho - 2 \ln(1 + \frac{1}{2}\beta) + \frac{1}{2} \ln(1 + \beta)] & [k_{DL}, k_{PC}] \neq [0,1] \end{cases}$$

$$= \begin{cases} 2 \ln \kappa_{i}(\beta) - \frac{1}{2} \ln \kappa_{i}(2\beta) - \frac{3}{2} \ln p - \frac{1}{2} \ln(2 - (1 + K^{-1})^{-2}) & [k_{DL}, k_{PC}] = [0,1] \end{cases},$$

$$(4.20)$$

$$b_{i}^{2} = -2 \ln E[I_{i}] + \ln E[I_{i}^{2}] + \ln p$$

$$= \begin{cases}
\sigma^{2} - 2 \ln \kappa_{i}(\beta) + \ln \kappa_{i}(2\beta) + \ln p + \ln(2 - (1 + K^{-1})^{-2}) \\
+ k_{PC}[\sigma^{2} + \sigma_{e}^{2} + 2 \ln(1 + \frac{1}{2}\beta) - \ln(1 + \beta)] \quad [k_{DL}, k_{PC}] \neq [0,1]
\end{cases}$$

$$(4.21)$$

$$2\sigma^{2} + \sigma_{e}^{2} - 2 \ln \tilde{\kappa}_{i}(\beta) + \ln \tilde{\kappa}_{i}(2\beta) + \ln p + \ln(2 - (1 + K^{-1})^{-2}) \quad [k_{DL}, k_{PC}] = [0,1].$$

4.2.5.4 Validation through Simulation

We simulate the distribution of the interference coming from one nearby interferer (class A), as well as from a distant interferer (class D) and compare the results with the theoretical LN distributions with parameters given by (4.20), (4.21). Note that we set p = 1, since the pdf of a RV multiplied by a Bernoulli RV is just the scaled original pdf with additional mass at zero, and there is no need to simulate this. Also, when there is no PC, both UL and DL interference are the same for one user, and thus the simulated results apply equally to both cases.

The simulations are done according to the parameters in Table 4-5, while $\sigma_e = 1 \text{dB} \approx 0.23$. The results are plotted on Figures 4-5 to 4-10. We observe an excellent match in the higher values of

the distributions, which can be expected, since moment-matching should give a good match to the upper tail [43], but the lower tails tend to diverge in most cases. The lower tail is well-matched only when there is no fading, and the interfering cell is far, meaning that the pathloss will not vary overmuch, and the power is essentially LN (Figure 4-8). The results are not encouraging in the lower tail. However, we will see in subsequent sections that although this affects the results of the problem of many interferers, the error in the lower tail does not necessarily accumulate, and may be in fact be quite small.

Figure		4-5			4-6			4-7	
β		3			4			3.5	
σ [dB]	6	9	12	6	9	12	6	9	12
K		0			∞			10	
Cell					Class A				
Link		UL/DL			DL			UL	
PC		No			Yes			Yes	
a_i		-2.002			-3.884			-3.567	
$oldsymbol{b}_i$	1.837	2.400	3.017	2.312	3.181	4.099	2.290	3.165	4.086
Figure		4-8			4-9			4-10	
β		3.5			3			4	
σ [dB]	6	9	12	6	9	12	6	9	12
K		∞			10			0	
Cell					Class D				
Link		UL/DL			DL			UL	
PC		No			Yes			Yes	
a_i		-5.767			-6.430			-8.685	
b_i	1.416	2.095	2.780	2.069	3.009	3.967	2.218	3.113	4.046

Table 4-5 Simulation parameters for single interferer power.

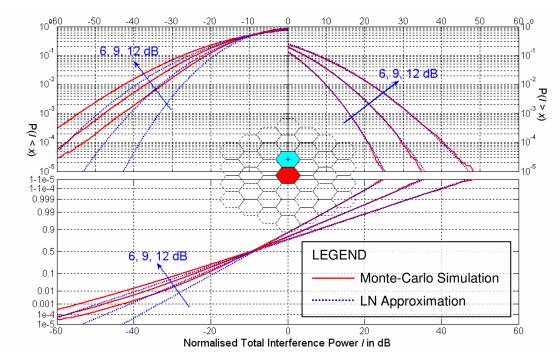


Figure 4-5 Single nearby interferer, Rayleigh fading, without power control.

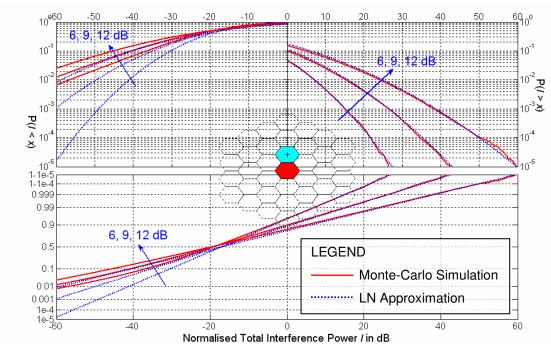


Figure 4-6 Single nearby interferer, no fading, downlink with power control.

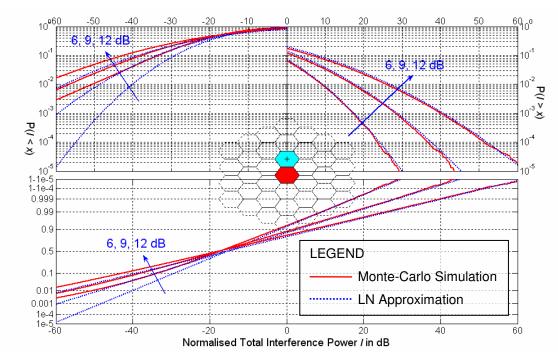


Figure 4-7 Single nearby interferer, Rician fading, uplink with power control.

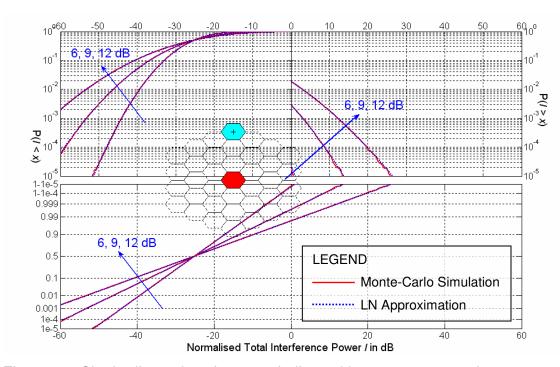


Figure 4-8 Single distant interferer, no fading, without power control.

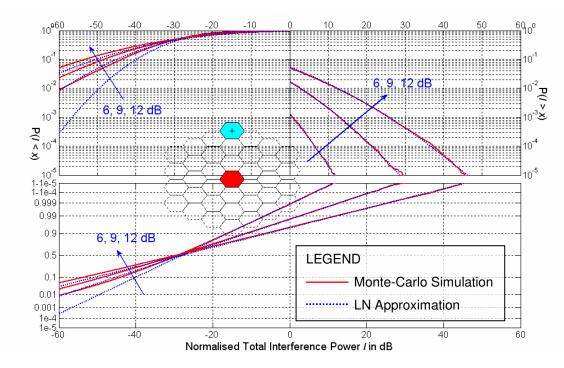


Figure 4-9 Single distant interferer, Rician fading, downlink with power control.

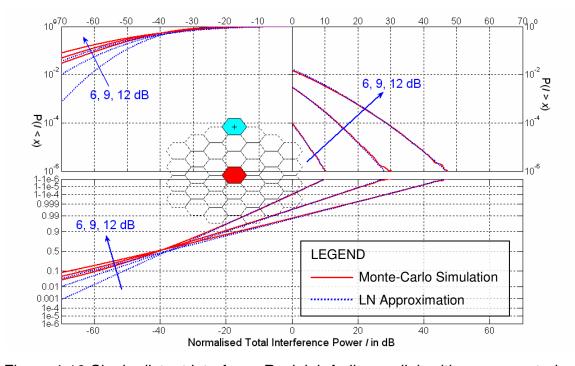


Figure 4-10 Single distant interferer, Rayleigh fading, uplink with power control.

4.3 Statistical Moments of Total Interference Power

Although the moments (of natural order) do not entirely characterise a distribution [43], they nevertheless offer much useful information about it. In this section, we present both an exact expression for the first two moments of the total interference power I, as well as good approximations, which hold under the assumption that the shadowing variance is large enough, which we find to be the case in realistic scenarios (6 – 12 dB).

4.3.1 First Moment

4.3.1.1 Uplink without Power Control, and Downlink

Let us define A_n coefficients, which are a function of the cellular layout and the propagation constant, as follows:

$$A_{n} = E\left[\sum_{i=1}^{N} R_{i}^{n}\right] = \sum_{i=1}^{N} E\left[R_{i}^{n}\right] = \sum_{i=1}^{N} \kappa_{i}(n\beta) . \tag{4.22}$$

In the DL, the DMU receives interference from all FS's. The power transmitted by these is independent of the position of other users in the interfering cells. Similarly, in the UL – no PC case, the transmit power of each IMU is independent of its position. We thus conclude that each pathloss is independent of all other factors, and we may write:

$$E[I] = E\left[\sum_{i=1}^{N} R_{i} Q_{i}\right] = \sum_{i=1}^{N} E[R_{i}] E[Q_{i}] = E\left[\sum_{i=1}^{N} R_{i}\right] E[Q_{i}] = A_{1} B_{1} \qquad [k_{DL}, k_{PC}] \neq [0,1].$$
 (4.23)

4.3.1.2 Uplink with Power Control

The method needs to be modified slightly for the UL – PC case. We define:

$$\widetilde{A}_n = E\left[\sum_{i=1}^N R_i^n L_i^n\right] = \sum_{i=1}^N E\left[R_i^n L_i^n\right] = \sum_{i=1}^N \widetilde{\kappa}_i(n\beta) . \tag{4.24}$$

Then:

$$E[I] = \sum_{i=1}^{N} E[R_i L_i] E\left[\frac{Q_i}{L_i}\right] = E\left[\sum_{i=1}^{N} R_i L_i\right] E\left[\frac{Q_i}{L_i}\right] = \tilde{A}_1 \tilde{B}_1 \qquad [k_{DL}, k_{PC}] = [0,1] . \tag{4.25}$$

4.3.2 Second Moment

4.3.2.1 Uplink without Power Control

$$E[I^{2}] = E\left[\left(\sum_{i=1}^{N} I_{i}\right)^{2}\right] = \sum_{i=1}^{N} \sum_{j=1}^{N} E[I_{i}I_{j}] = \sum_{i=1}^{N} E[I_{i}^{2}] + \sum_{i=1}^{N} \sum_{j=1}^{N} E[I_{i}I_{j}]$$

$$= \sum_{i=1}^{N} E[R_{i}^{2}] E[Q_{i}^{2}] + \sum_{i=1}^{N} \sum_{j=1}^{N} E[R_{i}R_{j}] E[Q_{i}] E[Q_{j}] = A_{2}B_{2} + B_{1}^{2} \sum_{i=1}^{N} \sum_{j=1}^{N} E[R_{i}R_{j}].$$

$$(4.26)$$

In the UL, the pathlosses are independent, hence:

$$E[I^{2}] = A_{2}B_{2} + \left(\sum_{i=1}^{N} \sum_{j=1}^{N} E[R_{i}]E[R_{j}] - \sum_{i=1}^{N} E[R_{i}^{2}]\right)B_{1}^{2}$$

$$= A_{2}B_{2} + \left(\left(\sum_{i=1}^{N} \kappa_{i}(\beta)\right)^{2} - \sum_{i=1}^{N} \kappa_{i}(2\beta)\right)B_{1}^{2} = A_{2}B_{2} + (A_{1}^{2} - A_{2})B_{1}^{2} \qquad k_{DL} = k_{PC} = 0.$$

$$(4.27)$$

4.3.2.2 Uplink with Power Control

Analogically, we can find the moment in UL-PC:

$$E[I^{2}] = \tilde{A}_{2}\tilde{B}_{2} + (\tilde{A}_{1}^{2} - \tilde{A}_{2})\tilde{B}_{1}^{2} \quad [k_{DL}, k_{PC}] = [0,1]. \tag{4.28}$$

4.3.2.3 Downlink

Define:

$$A_{DL} = \sum_{i=1}^{N} \sum_{j=1}^{N} E[R_i R_j] = E\left[\left(\sum_{i=1}^{N} R_i\right)^2\right]. \tag{4.29}$$

In the DL, the distances between the DMU and the IFS's are not independent and A_{DL} cannot be separated into a sum of terms, but must be integrated for the entire geometry as follows:

$$A_{DL} = \frac{2}{3\sqrt{3}} \iint_{(x,y)\in C_0} \left(\sum_{i=1}^N \sqrt{(x-x_i)^2 + (y-y_i)^2} \right)^2 dx dy,$$
 (4.30)

where x_i , y_i and C_0 are the same as in (4.2). Then:

$$E[I^{2}] = A_{2}B_{2} + \left(E\left[\sum_{i=1}^{N}\sum_{j=1}^{N}R_{i}R_{j}\right] - \sum_{i=1}^{N}E[R_{i}^{2}]\right)B_{1}^{2} = A_{2}B_{2} + (A_{DL} - A_{2})B_{1}^{2} \qquad k_{DL} = 1. \quad (4.31)$$

4.3.3 A Simple Approximation to Moments

We found in our numerical examples that only the first term in the second moment expressions is significant, and that the following approximation is usually exact to several digits:

$$E[I^2] \cong A_2 B_2 \qquad \forall k_{DL}, \forall k_{PC} \tag{4.32}$$

We conjecture that, in general,

$$E[I^n] \cong A_n B_n \qquad n = 1, 2, 3 \dots \qquad \forall k_{DL}, \forall k_{PC}$$

$$(4.33)$$

We believe the reason for this is that since B_n is dominated by a LN RV factor with high variance, which has moments of the order of e^{n^2} , therefore $B_n >>> B_{n-1}$. If write the multinomial expansion of (4.33), we will find that there will be one term containing B_n , which is the one given by the equation, which will dominate.

Notice that the expressions for the moments are then identical for UL and DL, which is suggested by [47].

4.4 Modeling the Total Interference Power as Bernoulli Lognormal

We approximate I by a LN RV of parameters (a_l, b_l^2) multiplied by an independent Bernoulli RV with parameter p_{χ} , giving a BLN distribution:

$$F_I(x) \cong \Phi\left(\frac{\ln x - a_I}{b_I}\right) p_{\chi} + (1 - p_{\chi}), x > 0,$$
 (4.34)

$$p_{\gamma} = P[I \neq 0] = 1 - (1 - p)^{N}. \tag{4.35}$$

4.4.1 Arguments for the Choice of Distribution

We have, according to the FW method, chosen to model a sum of (approximately) LN RV's by another LN RV. It is not self-evident that a linear combination of LN RV's with random weights (namely the Bernoulli RV's χ_i), will again be LN. Nevertheless, in [1] this assumption is made. We can argue that for any realisation of the vector χ_i , we have a sum of $\sum_{i=1}^{N} \chi_i$ interferers, which is itself approximately LN. The resulting pdf is then the weighed sum of the pdf's for each realisation of the vector χ_i , with the weights being the probability associated with each realisation. This is called a mixture of LN distributions; it is not known whether it is approximately LN. However we do know that because the individual distributions are smooth (LN), their sum will also be smooth, so a LN approximation to their sum is not inconceivable. We notice, however, that if the probability of activity p is low, there is a chance of there being absolutely no interference if all $\chi_i = 0$. This creates a discontinuity in the cdf of I, which modifies the LN

distribution into a BLN one in (4.34). In most cases, for many interferers and moderate to high p, we have $p_{\chi} \approx 1$ and the model becomes LN as in [1].

4.4.2 Moment-Matching

We match the moments of the interference with those of the approximating distribution:

$$E[I^n] = p_{\chi} \cdot e^{na_I + \frac{n^2}{2}b_I^2}. \tag{4.36}$$

Inverting (4.36), we obtain:

$$a_I = 2 \ln E[I] - \frac{1}{2} \ln E[I^2] - \frac{3}{2} \ln p_{\chi},$$
 (4.37)

$$b_I^2 = -2 \ln E[I] + \ln E[I^2] + \ln p_{\chi}. \tag{4.38}$$

4.4.3 Simple Approximations to the Distribution Parameters

Applying the approximation (4.32) to equations (4.37) and (4.38), we obtain simple closed-form equations that are applicable to all cases:

$$a_{I} \approx 2 \ln A_{1} B_{1} - \frac{1}{2} \ln A_{2} B_{2} - \frac{3}{2} \ln(p_{\chi})$$

$$= 2 \ln A_{1} - \frac{1}{2} \ln A_{2} + \frac{3}{2} \ln \frac{p}{p_{\chi}} - \frac{1}{2} \ln(2 - (1 + K^{-1})^{-2}) ,$$

$$+ k_{PC} \left[\beta \ln \rho - 2 \ln(1 + \frac{1}{2}\beta) + \frac{1}{2} \ln(1 + \beta)\right]$$
(4.39)

$$b_I^2 \cong -2 \ln A_1 B_1 + \ln A_2 B_2 + \ln p_{\chi}$$

$$= \sigma^2 - 2 \ln A_1 + \ln A_2 - \ln \frac{p}{p_{\chi}} + \ln \left(2 - \left(1 + K^{-1} \right)^{-2} \right) .$$

$$+ k_{PC} \left[\sigma^2 + \sigma_e^2 + 2 \ln \left(1 + \frac{1}{2} \beta \right) - \ln \left(1 + \beta \right) \right]$$

$$(4.40)$$

For the UL – PC case, this approximation simply treats R_j and L_j as independent, but with the same marginal distributions. We notice that the expressions are identical for UL and DL. A

similar result was obtained in [47]. This is not entirely surprising, since the interfering power from each individual cell has identical distribution in UL and DL. Only in the DL case, the pathlosses R_j are correlated. However, since these are further multiplied by i.i.d. LN RV's with high variance, the correlation between the individual powers becomes very small. In other words, because every path suffers from independent shadowing with large variance, this effectively "drowns out" the correlation between the pathlosses and makes them very nearly independent. Also, in the UL – PC case, the independent shadowing terms make the correlation between R_j and L_j insignificant.

4.5 Modeling the Total Interference Power as Bernoulli Power Lognormal

We approximate I by a PLN RV of parameters (m_I, s_I^2, t_I) multiplied by an independent Bernoulli RV with the same parameter p_{χ} as given by (4.35), resulting in a Bernoulli Power Lognormal (BPLN) distribution:

$$F_I(x) \cong \Phi^{t_I} \left(\frac{\ln x - m_I}{s_I} \right) p_{\chi} + (1 - p_{\chi}), x > 0.$$
 (4.41)

4.5.1 Arguments for the Choice of Distribution

Following the same reasoning as in Section 4.4.1, we modify the PLN approximation to the SLN by multiplying it by the same Bernoulli RV, resulting in the discontinuous model (4.41).

4.5.2 Matching the Upper Tail Slope

We have seen in Section 3.3.1.1 that we can match the slope of the upper tail (on LN paper) of the SLN and the PLN distributions, by taking the slope of the component with the highest mean from among the components of the highest variance. A few numerical examples can be used to verify that interference from the nearest cells has not only the highest mean power in dB (which

is natural), but also the highest dB spread. This is expected since the distance from a point in cell C_0 to a point in a faraway cell cannot vary overmuch for different points in the cells. The variability is much stronger for nearby cells.

In order to match the upper tail slope, we find the cdf of the interference coming from one of the nearest cells, given that the user activity is constant (p = 1), since the result is only valid for a sum of LN RV's, and not of BLN RV's. For p < 1, we argue that a BLN cdf is just a scaled version of a corresponding LN cdf in the upper tail, and thus, by Lemma 2, these have the same blf, since their ratio is finite in the tail.

4.5.2.1 Uplink without Power Control, and Downlink

We first define the convenient intermediary quantity:

$$z_{I} = \ln \frac{B_{2}}{p} - 2 \ln \frac{B_{1}}{p} = \sigma^{2} + \ln \left(2 - \left(1 + K^{-1} \right)^{-2} \right) + k_{PC} \left[\sigma^{2} + \sigma_{e}^{2} + 2 \ln \left(1 + \frac{1}{2} \beta \right) - \ln \left(1 + \beta \right) \right].$$
(4.42)

We call C_1 a nearest cell to C_0 , and we can find the logarithmic variance of the interference component coming from this nearest cell by applying (4.21) to (3.48):

$$s_I^2 = b_1^2 = \ln \kappa_1(2\beta) - 2\ln \kappa_1(\beta) + z_I \quad [k_{DL}, k_{PC}] \neq [0,1].$$
 (4.43)

4.5.2.2 Uplink with Power Control

Analogically, for UL – PC:

$$\tilde{z}_I = \ln \frac{\tilde{B}_2}{p} - 2 \ln \frac{\tilde{B}_1}{p} = 2\sigma^2 + \sigma_e^2 + \ln \left(2 - \left(1 + K^{-1} \right)^{-2} \right),$$
 (4.44)

$$s_L^2 = \ln \tilde{\kappa}_1(2\beta) - 2\ln \tilde{\kappa}_1(\beta) + \tilde{z}_L \quad [k_{DL}, k_{PC}] = [0,1].$$
 (4.45)

4.5.3 Matching the Lower Tail Slope

4.5.3.1 A Simple Approximation for All Cases

If all components are *i.i.d.*, the PLN approximation to the SLN uses N as its power parameter ,as in [44]. Even if the components have different distribution, we find that this choice of parameter gives a good fit to the lower tail slope. However, if the individual components are multiplied by Bernoulli RV's with parameter p, then there are, on average, only pN interferers, which is the idea in [3]. We can thus write:

$$t_I \approx pN, \quad t_I \ge pN.$$
 (4.46)

We proceed to show in the next section that this approximation is good.

4.5.3.2 Uplink without Power Control, and Downlink

A more exact approach to finding the power parameter of the PLN is to use (3.49), which equates the power parameter with $b_1^2 \sum_{i=1}^{N} \frac{1}{b_i^2}$. By applying (4.21) and multiplying by p, as was done in (4.46), we obtain:

$$t_{I} = ps_{I}^{2} \sum_{i=1}^{N} \frac{1}{\ln \kappa_{i}(2\beta) - 2\ln \kappa_{i}(\beta) + z_{I}} \quad [k_{DL}, k_{PC}] \neq [0,1],$$
 (4.47)

where s_I^2 is given by (4.43).

We rewrite (4.47) as

$$t_{I} = p \sum_{i=1}^{N} \frac{\ln \kappa_{1}(2\beta) - 2\ln \kappa_{1}(\beta) + z_{I}}{\ln \kappa_{i}(2\beta) - 2\ln \kappa_{i}(\beta) + z_{I}} \quad [k_{DL}, k_{PC}] \neq [0,1], \tag{4.48}$$

and we observe, from numerical calculations, that:

$$\forall i : \ln \kappa_i(2\beta) - 2\ln \kappa_i(\beta) \le \ln \kappa_1(2\beta) - 2\ln \kappa_1(\beta) << z_I. \tag{4.49}$$

Applying (4.49) to (4.48), we obtain the approximation in (4.46).

4.5.3.3 Uplink with Power Control

Likewise, for the UL – PC case, we have:

$$t_{I} = p s_{I}^{2} \sum_{i=1}^{N} \frac{1}{\ln \tilde{\kappa}_{i}(2\beta) - 2\ln \tilde{\kappa}_{i}(\beta) + \tilde{z}_{I}} \quad [k_{DL}, k_{PC}] = [0,1], \tag{4.50}$$

where s_I^2 is given by (4.45).

4.5.4 Moment-Matching

Having matched the slopes at both tails, we have only one more degree of freedom. We match the first moment of the PLN approximation with that of the interference, given by (4.25).

4.5.4.1 Uplink without Power Control, and Downlink

Matching the moments gives:

$$EI = e^{m_I} \Lambda(s_I, t_I) p_{\chi} = A_1 B_1 \quad [k_{DL}, k_{PC}] \neq [0,1],$$
 (4.51)

where, the Λ function is given by (3.50). We proceed as in Section 3.3.1.2 to find m_l :

$$m_{I} = \ln A_{1}B_{1} - \ln \Lambda(s_{I}, t_{I})p_{\chi} = \frac{1}{2}\sigma^{2} + \ln A_{1} - \ln \Lambda(s_{I}, t_{I})$$

$$+ \ln \frac{p}{p_{\chi}} + k_{PC} \left[\frac{1}{2}\sigma^{2} + \frac{1}{2}\sigma_{e}^{2} + \beta \ln \rho - \ln(1 + \frac{1}{2}\beta) \right] \quad [k_{DL}, k_{PC}] \neq [0,1].$$
(4.52)

4.5.4.2 Uplink with Power Control

Likewise, for UL – PC:

$$EI = e^{m_I} \Lambda(s_I, t_I) p_{\gamma} = \tilde{A}_1 \tilde{B}_1 \quad [k_{DL}, k_{PC}] = [0,1].$$
 (4.53)

$$m_{I} = \ln \tilde{A}_{1} \tilde{B}_{1} - \ln \Lambda(s_{I}, t_{I}) p_{\chi}$$

$$= \sigma^{2} + \frac{1}{2} \sigma_{e}^{2} + \ln \tilde{A}_{1} - \ln \Lambda(s_{I}, t_{I}) + \ln \frac{p}{p_{\chi}} \quad [k_{DL}, k_{PC}] = [0,1].$$
(4.54)

4.6 Comparison of Simulation, BLN and BPLN Methods

4.6.1 Cellular Layouts and their Coefficients

Tables 4-6 to 4-10 have the geometry coefficients tabulated against β for five different cellular layouts: a single tier of nearest interferers (IT), three tiers of interferers in a reuse 1 cellular network (3T), a reuse 3 cellular network (R3), a reuse 7 cellular network (R7), and a random layout (RD) to test the robustness of our method. Only the A_{DL} coefficients are really new information, and need to be integrated as per (4.30). All other coefficients can be found from Tables 4-1 to 4-4 by summing over the individual cells using (4.22) and (4.24).

β	A_1	A ₂	$/A_{ m DL}$	\widetilde{A}_1	\widetilde{A}_2
3	1.679	1.008	.2.978	0.6042	0.3169
3.2	1.587	0.9868	2.701	0.5548	0.3107
3.4	1.505	0.9743	2.474	0.5130	0.3079
3.5	1.467	0.9710	2.377	0.4945	0.3076
3.6	1.432	0.9695	2.289	0.4775	0.3079
3.8	1.368	0.9718	2.138	0.4473	0.3105
4	1.310	0.9805	2.016	0.4215	0.3153
4.2	1.260,	0.9953	1.920	0.3994	0.3224
4.4	1.215	1.016	1.845	0.3805	0.3314
4.5	1.194	1.028	(1.815 \)	(0.3721 >-	0.3367
4.6	1.175	1.042	1.789	0.3643	0.3426
4.8	1.140	1.074	(1.749)	0.3504	0.3557
5	1.109	1.112	1.725	0.3386	0.3709
β	A_1	A ₂	A _{DL}	$\widetilde{\mathrm{A}}_{\mathrm{l}}$	\widetilde{A}_2
β 3	A ₁ 2.258	A ₂	A _{DL} + 5.280	Ã ₁	Ã ₂ 0.3205
	•		<u> </u>		-
3	2.258	1.027	5.280	0.7874	0.3205
3 3.2	2.258	1.027	5.280	0.7874	0.3205 0.3130
3 3.2 3.4	2.258 2.041 1.862	1.027 0.9995 0.9827	5.280 4.370 3.697	0.7874 0.6915 0.6155	0.3205 0.3130 0.3093
3 3.2 3.4 3.5	2.258 2.041 1.862 1.785	1.027 0.9995 0.9827 0.9779	5.280 4.370 3.697 3.427	0.7874 0.6915 0.6155 0.5835	0.3205 0.3130 0.3093 0.3087
3 3.2 3.4 3.5 3.6	2.258 2.041 1.862 1.785	1.027 0.9995 0.9827 0.9779 0.9752	5.280 4.370 3.697 3.427 3.194	0.7874 0.6915 0.6155 0.5835 0.5548	0.3205 0.3130 0.3093 0.3087 0.3089
3 3.2 3.4 3.5 3.6 3.8	2.258 2.041 1.862 1.785 1.714 1.590	1.027 0.9995 0.9827 0.9779 0.9752 0.9756	5.280 4.370 3.697 3.427 3.194 2.814	0.7874 0.6915 0.6155 0.5835 0.5548 0.5057	0.3205 0.3130 0.3093 0.3087 0.3089 0.3111
3 3.2 3.4 3.5 3.6 3.8 4	2.258 2.041 1.862 1.785 1.714 1.590 1.487	1.027 0.9995 0.9827 0.9779 0.9752 0.9756 0.9830 0.9970	5.280 4.370 3.697 3.427 3.194 2.814 2.526 2.307	0.7874 0.6915 0.6155 0.5835 0.5548 0.5057 0.4659	0.3205 0.3130 0.3093 0.3087 0.3089 0.3111 0.3157
3 3.2 3.4 3.5 3.6 3.8 4 4.2	2.258 2.041 1.862 1.785 1.714 1.590 1.487 1.400	1.027 0.9995 0.9827 0.9779 0.9752 0.9756 0.9830 0.9970	5.280 4.370 3.697 3.427 3.194 2.814 2.526 2.307	0.7874 0.6915 0.6155 0.5835 0.5548 0.5057 0.4659 0.4333	0.3205 0.3130 0.3093 0.3087 0.3089 0.3111 0.3157 0.3226
3 3.2 3.4 3.5 3.6 3.8 4 4.2 4.4	2.258 2.041 1.862 1.785 1.714 1.590 1.487 1.400 1.326	1.027 0.9995 0.9827 0.9779 0.9752 0.9756 0.9830 0.9970 1.017	5.280 4.370 3.697 3.427 3.194 2.814 2.526 2.307 2.142 2.076	0.7874 0.6915 0.6155 0.5835 0.5548 0.5057 0.4659 0.4333 0,4065	0.3205 0.3130 0.3093 0.3087 0.3089 0.3111 0.3157 0.3226 0.3316
3 3.2 3.4 3.5 3.6 3.8 4 4.2 4.4 4.5	2.258 2.041 1.862 1.785 1.714 1.590 1.487 1.400 1.326 1.294	1.027 0.9995 0.9827 0.9779 0.9752 0.9756 0.9830 0.9970 1.017 1,029	5.280 4.370 3.697 3.427 3.194 2.814 2.526 2.307 2.142 2.076	0.7874 0.6915 0.6155 0.5835 0.5548 0.5057 0.4659 0.4333 0.4065 0.3949	0.3205 0.3130 0.3093 0.3087 0.3089 0.3111 0.3157 0.3226 0.3316 0.3369

Table 4-6 Geometry parameters for various cellular layouts.

β	A_1	A ₂	A _{DL}	\widetilde{A}_1	$\widetilde{\mathbf{A}}_2$
3	0.2921	0.01303	0.08559	0,09301	0.002510
3.2	0.2339	0.008843	0.05497	0.07097	0.001626
3.4	0.1879	0.006027	0.03549	0.05441	0.001062
3.5	0.1685	0.004983	0.02857	0.04773	0.0008605
3.6	0.1512	0.004124	0.02303	0.04191/	0.0006983
3.8	0.1220	0.002833	0.01502	0.03243	0:0004621
4	0.09868	0.001952	0.009844	0.02519	0.0003075
4.2	0.07997	0.001350	0.006480	(0.01965)	0:0002057
4.4	0.06494 +	0,0009358 +	0.004284	+ 0.01539	+ 0.0001382
4.5	0.05856	0.0007799	0.003489	0.01363	0.0001134
4.6	0.05283	0.0006505	0.002844	0.01209	0.00009318
4.8	0.04305	0.0004533	(0.001895)	0.009531	0.00006308
5	0.03515	0.0003166	0.001268	0.007537	0.00004284
β	\mathbf{A}_1	A_2	A_{DL}	\widetilde{A}_1	$\mathbf{\tilde{A}}_2$
$\frac{\beta}{3}$	A ₁ 0.06525	0.0007766	0.004260	0,02032	$\frac{\tilde{A}_2}{0.0001320}$
	<u> </u>		<u>`</u> }		
3	0.06525	0.0007766	0.004260	0,02032	0.0001320
3 3.2	0.06525	0.0007766	0.004260	0,02032	0.0001320 0.00006918
3 3.2 3.4	0.06525 0.04843 0.03595	0.0007766 0.0004330 0.0002417	0.004260 0.002347 0.001294	0,02032 0.01430 0,01009	0.0001320 0.00006918 0.00003643
3 3.2 3.4 3.5	0.06525 0.04843 0.03595 0.03099	0.0007766 0.0004330 0.0002417 0.0001807	0.004260 0.002347 0.001294 0.0009614	0,02032 0.01430 0,01009 0.008480	0.0001320 0.00006918 0.00003643 0.00002648
3 3.2 3.4 3.5 3.6	0.06525 0.04843 0.03595 0.03099 0.02671	0.0007766 0.0004330 0.0002417 0.0001807 0.0001352	0.004260 0.002347 0.001294 0.0009614 0.0007142	0,02032 0.01430 0,01009 0.008480 0.007132	0.0001320 0.00006918 0.00003643 0.00002648 0.00001927
3 3.2 3.4 3.5 3.6 3.8	0.06525 0.04843 0.03595 0.03099 0.02671 0.01984	0.0007766 0.0004330 0.0002417 0.0001807 0.0001352 0.00007568	0.004260 0.002347 0.001294 0.0009614 0.0007142 0.0003945	0,02032 0.01430 0,01009 0.008480 0.007132 0.005053	0.0001320 0.00006918 0.00003643 0.00002648 0.00001927 0.00001024
3 3.2 3.4 3.5 3.6 3.8 4	0.06525 0.04843 0.03595 0.03099 0.02671 0.01984 0.01475	0.0007766 0.0004330 0.0002417 0.0001807 0.0001352 0.00007568	0.004260 0.002347, 0.001294 0.0009614 0.0007142 0.0003945 0.0002181	0,02032 0.01430 0,01009 0.008480 0.007132 0.005053 0.003587	0.0001320 0.00006918 0.00003643 0.00002648 0.00001927 0.00001024 +0.000005467
3 3.2 3.4 3.5 3.6 3.8 4 4.2	0.06525 0.04843 0.03595 0.03099 0.02671 0.01984 0.01475 0.01097	0.0007766 0.0004330 0.0002417 0.0001807 0.0001352 0.00007568 0.00004243 0.00002382	0.004260 0.002347, 0.001294 0.0009614 0.0007142 0.0003945 0.0002181 0.0001207	0,02032 0.01430 0,01009 0.008480 0.007132 0.005053 0.003587 0.002551	0.0001320 0.00006918 0.00003643 0.00002648 0.00001927 0.000001024 0.000005467 0.000002930
3 3.2 3.4 3.5 3.6 3.8 4 4.2 4.4	0.06525 0.04843 0.03595 0.03099 0.02671 0.01984 0.01475 0.01097 0.008161	0.0007766 0.0004330 0.0002417 0.0001807 0.0001352 0.00007568 0.00004243 0.00002382 0:00001339	0.004260 0.002347 0.001294 0.0009614 0.0007142 0.0003945 0.0002181 0.0001207 0.00006682	0,02032 0.01430 0,01009 0.008480 0.007132 0.005053 0.003587 0.002551 0.001818	0.0001320 0.00006918 0.00003643 0.00002648 0.00001927 0.000001024 0.000005467 0.000002930 0.000001576
3 3.2 3.4 3.5 3.6 3.8 4 4.2 4.4 4.5	0.06525 0.04843 0.03595 0.03099 0.02671 0.01984 0.01475 0.01097 0.008161 0.007040	0.0007766 0.0004330 0.0002417 0.0001807 0.0001352 0.00004243 0.00004243 0.00001339 0.00001005	0.004260 0.002347 0.001294 0.0009614 0.0007142 0.0003945 0.0002181 0.0001207 0.00006682 0.00004974	0,02032 0.01430 0,01009 0.008480 0.007132 0.005053 0.003587 0.002551 0.001818 0.001536	0.0001320 0.00006918 0.00003643 0.00002648 0.00001927 0.00001024 0.000005467 0.000002930 0.000001576 0000001157

Table 4-6 (continued)

β	A_1	A ₂	A _{DL}	\widetilde{A}_{1}	$ ilde{ ilde{A}}_2$
3	1.017	0.5087	1.196	0.3578	0.1593
3.2	0.9300	0.4964	1.043	0.3183	0.1559
3.4	0.8582	0.4891	0.9308	0.2867	0.1543
3.5	0.8268	0.4870	0.8862	0.2732	0.1540
3.6	0.7981	0.4860	0.8479	0.2611	0.1542
3.8	0.7476	0.4867	0.7871	0.2403	0.1554
4	0.7049 +	0.4908	0.7433	0,2231	_ 0.1578
4.2	0.6686	0.4980	0.7127	0.2090	0.1612
4.4	0.6378 +	0.5082	0.6926	0.1973	0.1658
4.5	0.6241	0.5143	0.6860	(0.1921	Ø1684
4.6	0.6114	0.5212	0.6813	0,1875	0.1713
4.8	0.5888	0.5371	(0.6773)	⁶ 0.1792	0.1779
5	0.5695	0.5559	0.6796	0.1724	0.1855

Table 4-6 (continued)

4.6.2 Simulated and Analytical Parameters and Graphs

We run simulations according to the algorithm in Figures 4-2 and 4-3. The input parameters are given in Table 4-7 while $\sigma_e = 1 \text{dB} \approx 0.23$. Figures 4-11 to 4-19 show the results of the simulations, as well as the analytical approximations.

- 1. The exact BLN method uses (4.27), (4.28), or (4.31) in (4.37) and (4.38) to obtain the exact LN parameters a_I and b_I .
- 2. The approximate BLN method uses (4.39) and (4.40) to obtain approximate values of these parameters.
- 3. Both the exact and approximate BPLN methods use (4.42) and (4.43), or (4.44) and (4.45) to find the first parameter s_I .

- 4. The approximate BPLN method uses (4.46) to find t_I .
- 5. The exact BPLN method uses (4.42) and (4.47), or (4.44) and (4.50) to find t_I .
- 6. Both the exact and approximate methods use (4.52) or (4.54) to find m_l , where the integral $\Lambda(s_I, t_I)$ is calculated numerically as per (3.50).

Figure		4-11		4-	12		4-13	
β		4		۷	1		3	
Σ [dB]		12		1	2	6	9	12
K		0		()		0	
p	0.01	0.1	1	1	1		1	
Layout		3T		1T	3T		1T	
Link		DL		D	L		UL	
PC		No		Y	es		No	
p_{χ}	0.304	0.977	1	1	1		1	
$a_I =$	-4.665	-2.965	0.455	-1.542	-1.293	0.624	0.680	0.686
$a_I \approx$	-4.665	-2.965	0.455	-1.542	-1.293	0.686	0.686	0.686
$b_I =$	3.308	3.132	2.744	4.007	3.976	1.306	1.995	2.703
$b_I \approx$	3.308	3.132	2.744	3.962	3.931	1.258	1.992	2.703
$m_I =$	NA	-4.183	-4.026	-4.224	-5.915	-1.666	-1.837	-1.937
$m_I \approx$	NA	-4.092	-3.942	-4.224	-5.866	-1.666	-1.837	-1.937
s_I	NA	3.0	94	4.2	225	1.837	2.400	3.017
$t_I =$	NA	3.960	39.60	6	37.83		6	
$t_I \approx$	NA	3.6	36	6	36		6	

Table 4-7 Simulation and analytical parameters for BLN and BPLN methods.

Figure	4-14		4-	-15		4-16		
β	3.5	4					4	
σ [dB]	10		1	12			12	
K	0			0		0	10	∞
p	1	0.02	0.1	0.5	1		1	
Layout	<i>R7</i>		I	R <i>3</i>			1T	
Link	UL		Ι	DL			DL	
PC	No		Y	res		Yes		
p_{χ}	1	0.215	0.718	0.99976	1		1	
$a_I =$	-2.991	-7.171	-6.563	-4.646	-3.607	-1.542	-1.276	-1.196
$a_I \approx$	-2.986	-7.171	-6.563	-4.646	-3.607	-1.542	-1.276	-1.196
$b_I =$	2.082	4.168	4.120	3.961	3.873	4.007	3.940	3.919
$b_I \approx$	2.080	4.126	4.076	3.916	3.827	3.962	3.894	3.874
$m_I =$	-5.525	NA	-6.752	-7.076	-7.068	-4.224	-3.956	-3.876
$m_I \approx$	-5.525	NA	-6.744	-7.068	-7.061	-4.224	-3.956	-3.876
s_I	2.473	NA		4.119		4.225	4.161	4.142
$t_I =$	6	NA	1.209	6.043	12.09		6	
$t_I \approx$	6	NA	1.2	6	12		6	

Table 4-7 (continued)

Figure	4-	17		4-18			4-19
β	3	4		3.5			
σ [dB]	12	6		8			
K	∞	0	0			5	
p		1	0.02	0.1	0.5	1	1
Layout	1	T		I	R <i>3</i>		RD
Link	U	L		U	JL		DL
PC	Y	es	Yes				PC
p_{χ}		1	0.215	0.718	0.99975	1	1
$a_I =$	-0.394	-1.438	-6.431	-5.822	-3.906	-2.867	-2.224
$a_I \approx$	-0.373	-1.542	-6.404	-5.795	-3.879	-2.839	-2.224
$b_I =$	3.891	2.253	2.983	2.914	2.686	2.554	4.022
$b_I \approx$	3.806	2.061	2.895	2.824	2.588	2.450	3.978
$m_I =$	-3.073	-4.007	NA	-6.041	-6.300	-6.244	-5.517
$m_I \approx$	-3.073	-4.007	NA	-6.024	-6.284	-6.228	-5.469
s_I	4.115	2.621	NA	2.922	2.922	2.922	4.174
$t_I =$	(5	NA	1.221	6.107	12.21	14.70
$t_I \approx$	(6	NA	1.2	6	12	14

Table 4-7 (continued)

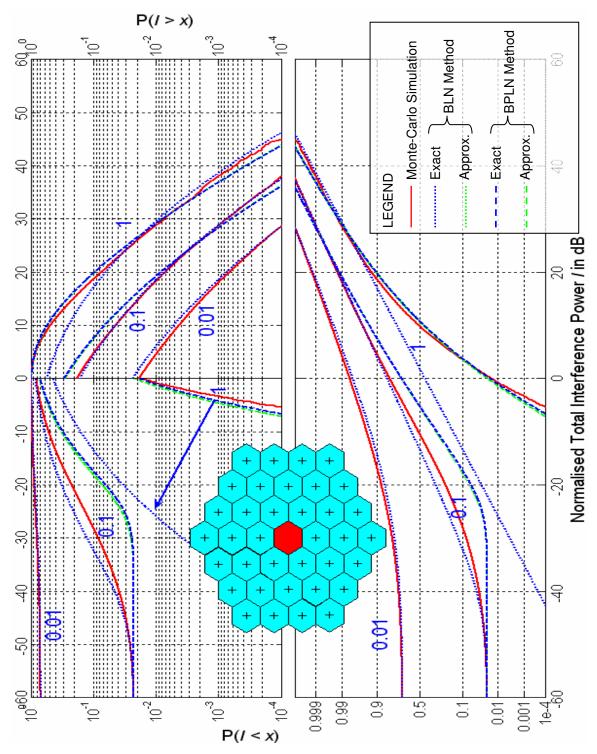


Figure 4-11 Interference in three-tier cellular network with variable activity in the downlink without power control.

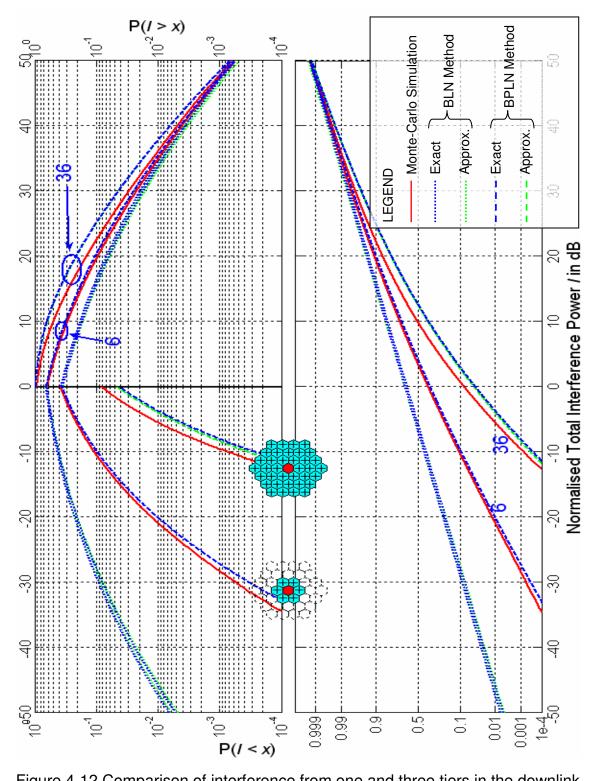


Figure 4-12 Comparison of interference from one and three tiers in the downlink with power control.

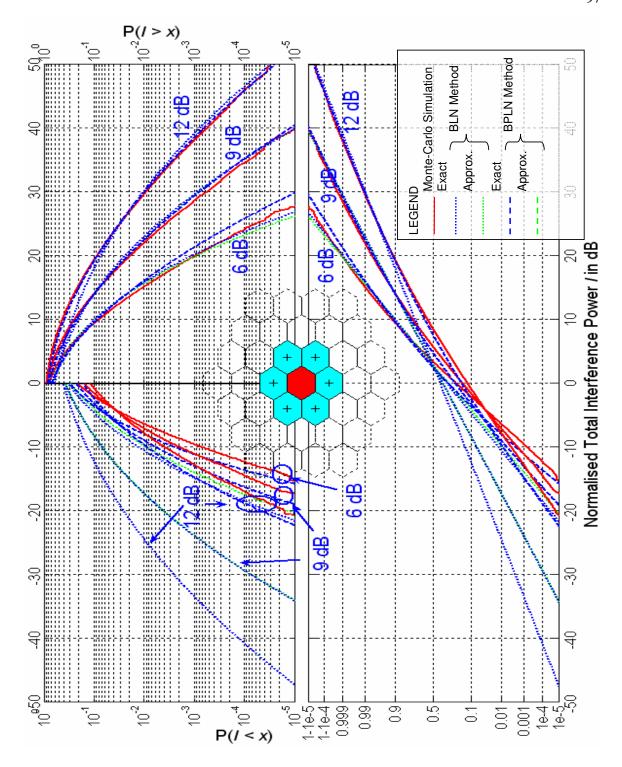
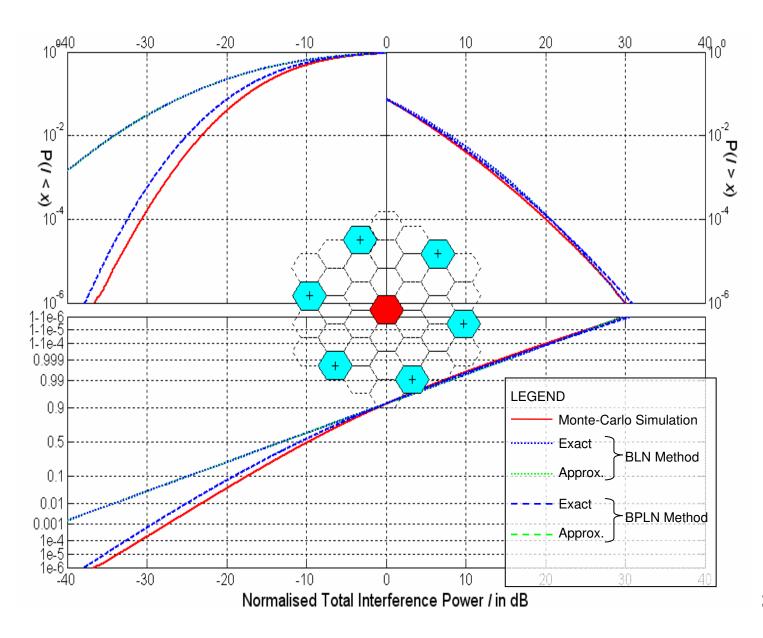


Figure 4-13 Effect of varying the shadowing parameter on interference from first tier in the uplink without power control.

control. Figure 4-14 Interference in a reuse 7 cellular network in the uplink without power



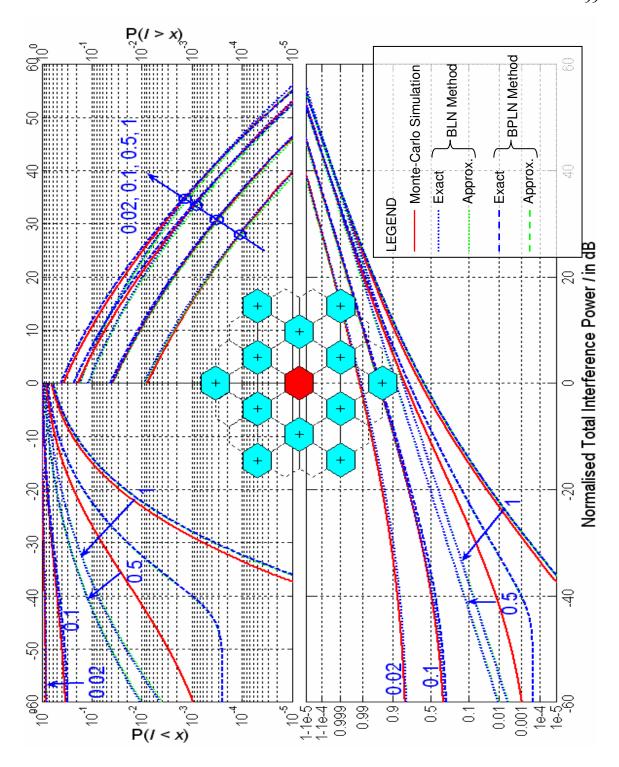


Figure 4-15 Interference in reuse 3 cellular network with variable activity level in the downlink with power control.

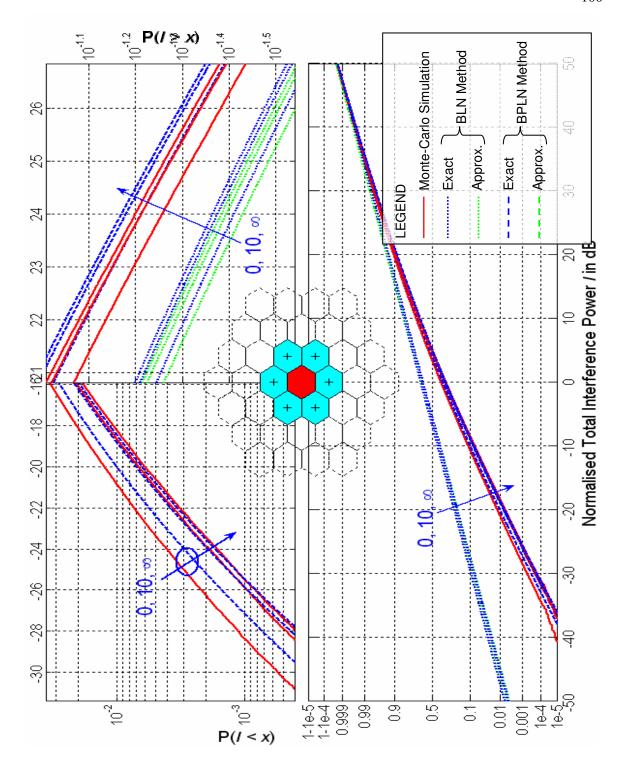


Figure 4-16 Interference from first tier with Rayleigh, Rician and no fading in the downlink with power control.

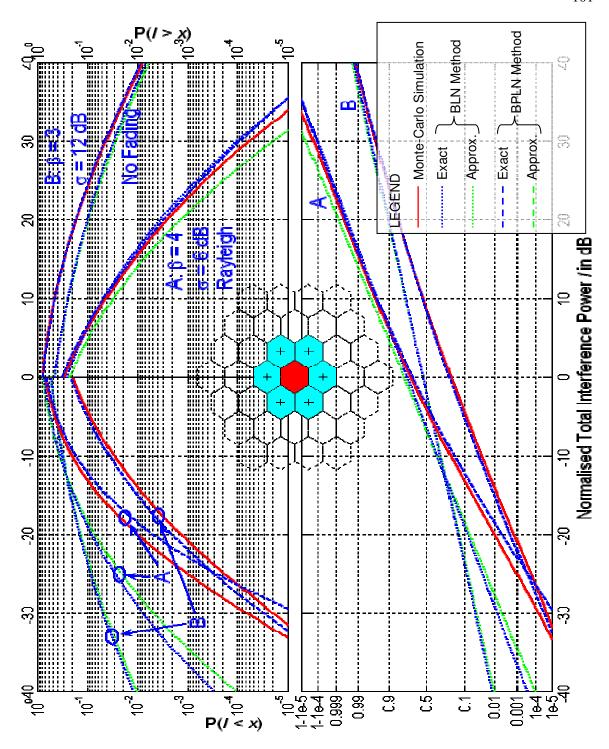


Figure 4-17 Interference from first tier while varying propagation exponent, shadowing and fading in the uplink with power control.

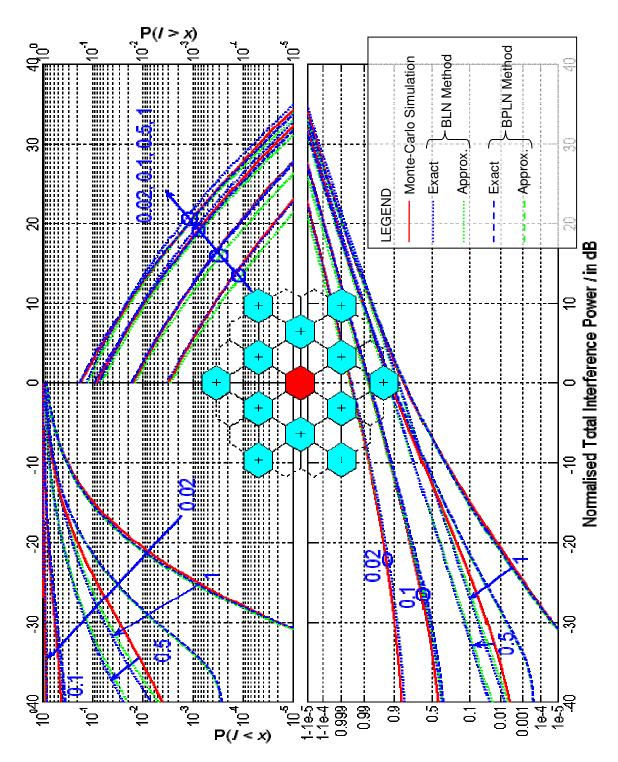
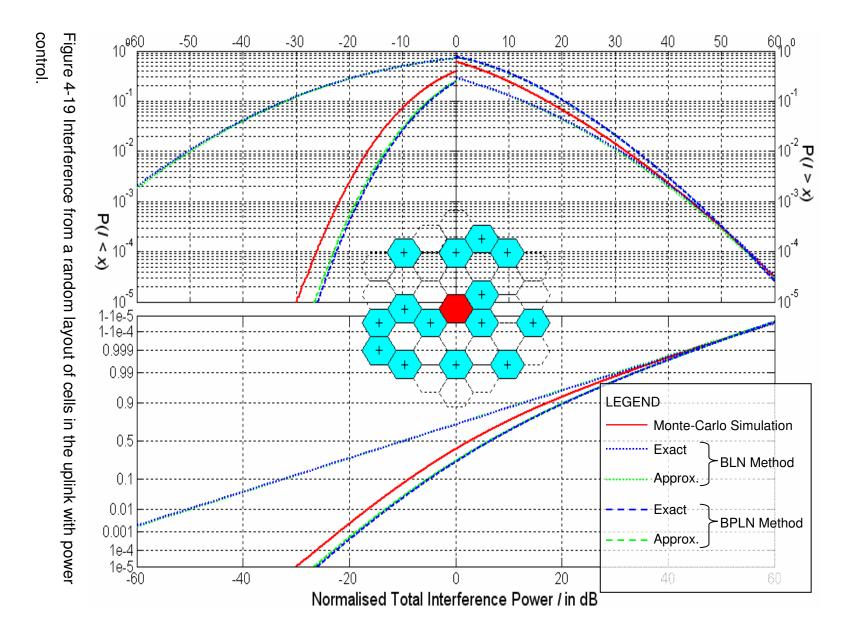


Figure 4-18 Interference in reuse 3 cellular network with variable activity level in the uplink with power control.



4.6.3 Discussion on Observed Results

Based on Figures 4-11-4-19, we are able to come to several important conclusions:

- 1. In most cases, the difference between the approximate (green) and the corresponding exact (blue) curves are negligible compared with the error to the simulated curve. This leads us to the conclusion that the approximations that we have made are very good. Since they are significantly simpler than the exact forms, this is a very valuable finding. We note however, in Figures 4-17 and 4-18, where we considered UL PC, that the exact BLN curves perform much better in the upper tail than their approximate counterparts. This is because, in the approximate expression, we assumed R_j and L_j to be independent. We conclude that this independence assumption should not be made, and that the modified method for the UL PC case is the proper way to proceed.
- 2. For all cases where there is full activity (*p* = 1), we observe that the BPLN outperforms the BLN method, particularly in the lower tail. This is to be expected: moment-matching tends to match the higher values of the distribution [43], but the lower tail cannot be simultaneously matched, because the SLN *cdf* bends on LN paper, whereas the BLN method cannot. However, since the BPLN was designed to match the slope of the lower tail of the SLN *cdf*, we observe that this indeed happens. When the BPLN does not match the lower tail well, it nevertheless runs parallel to it. We would need more degrees of freedom in the approximating *cdf* to always match the tail, not only in slope, but also in offset.
- 3. For low activity ($p \le 0.1$), we observe that the BLN method performs very well, while the BPLN method does not fit the lower tail properly. In the case when pN < 1, the BPLN curve is not even defined. Incidentally, notice that in the lower tail, the probability does

not go to 0, but tapers off to some value. This value is $P(I=0)=1-p_{\chi}$. We can observe that multiplying our LN and PLN models by a Bernoulli RV was a judicious move.

- 4. The conclusion is to use the BPLN method for *p* close to 1, and use the BLN method for *p* for values around 0.1 and lower. For moderate values of *p*, the problem remains open, as neither method works well in the lower tail. It would seem that in this case more work is needed, analysing the problem of a sum of Bernoulli-times-lognormal RV's separately from the SLN problem. In the mean time, it would seem that the BLN and BPLN methods give two bounds for the lower tail. Perhaps a weighed average of the two approximating *cdf* s, with weights that are functions of *p*, would give good results.
- 5. In general, leaving *p* aside, the method works well for a wide variety of shadowing, fading and propagation constants, in both UL and DL, with and without PC. It works for many practical cellular layout scenarios, and seems to perform better for regular layouts than irregular (Figure 4-19) ones. It is not clear why this is the case.
- Observing Figures 4-16 and 4-17, we see that the performance in the lower tail degrades as the fading model becomes more Rayleigh ($K \rightarrow 0$). This could be anticipated from the results in Section 4.2.5.4, where we found that fading causes the SLN *cdf* to bend significantly in the lower tail. However, we are pleased that the effect does not seem to be cumulative when we add several such RV. In fact, even with the unpromising results for one interferer, the approximation to the sum of multiple interferers' powers is quite good.

In conclusion, for low loading we suggest the BLN model, and for p = 1, the BPLN model. Notice that the BPLN model dramatically outperforms the classical moment-matching approach for constant activity and is perhaps the best known quasi-closed-form solution to the problem.

Chapter 5

Conclusion and Future Work

We conclude with recalling the contributions that were made in this thesis, and propose several ideas for improving our method, and applying it to more complicated problems.

5.1 Achievements of This Thesis

While we did not meet all the ambitious requirements that we set up for ourselves in Section 1.1.3, we were nevertheless able to develop a set of quite simple equations that approximate the SLN distribution and the interference distribution with better precision than other known methods of comparable complexity.

5.1.1 Tails of Sum of Independent Lognormals

We showed that the lower tail of the SLN distribution in the independent case can be approximated by a LN distribution with parameters given by closed-form expressions. To our knowledge, there has to date been no analytical study of the lower tail. We defined the *best lognormal fit* in order to formally study whether the SLN *cdf* has LN tails, and also whether proposed approximating *cdf* s have this property. We propose the *blf* to be a design criterion for future approximating *cdf* s. This will ensure that the approximation remains good in the far tails, and also that moment-matching can be used.

5.1.2 Simple Expressions for Moments of Total Interference Power

We gave exact closed-form expressions for the first two statistical moments of the total interference power in a cellular system, not only for UL, but also for the correlated DL case. These expressions contain geometrical coefficients that have been tabulated against the

propagation constant and can be used directly in our expressions. We also showed that the expression for the second moment can be significantly simplified without loss of precision.

5.1.3 Approximate Interference Distributions: Two Methods

We have found that applying the FW method to the interference problem using a BLN resulted in simple closed-form expressions and gave excellent results for low probabilities of user activity (p). For fully loaded systems with continuous interference (p = 1), we have found that using a PLN function to approximate the interference distribution also gave good results, particularly in the lower tail, which has often been neglected.

5.2 Ideas for Continuation of this Work

While working on the material presented in this thesis, we have encountered many other ideas that follow the same general line of reasoning, and that could potentially give an even better solution to the problems of the SLN and interference distributions. We also show that our approach can be modified to solve for a more detailed interference model without altering the fundamentals of our method.

Time did not permit us to follow all these ideas, so we only give a sketch of them here, with the hope that we, or someone else, will develop them more fully.

5.2.1 Improving the SLN and Interference Approximations

We have developed two methods for analysing the interference distribution: a moment-matching method using a BLN distribution, and a novel method using a BPLN distribution. We found that the second method performs very well when p = 1, i.e., the distribution is PLN. The fit in the upper tail and body is always excellent, but tends to somewhat diverge in the lower tail. There are a few reasons for this:

- 1. The distribution of the interference I_i coming from just one cell is not quite well approximated in the lower tail, as seen in Figures 4-5 4-10.
- 2. The PLN distribution does not have a *blf* in the lower tail, while the SLN distribution does. Only the slopes are matched in the lower tail, as can be seen in Figures 4-11 4-19.
- The lower tail was calculated based on the assumption that all summands are independent. However, in the DL, the individual interference components are slightly correlated.

The second reason could be remedied by developing a better SLN method, one that matches both tails' *blf*, while also fitting the body. We have proposed such a method in Sections 3.2.2.2 and 3.3.2. However the initial results show a very poor match in the body of the distribution in Figures 3-5 – 3-8. We believe that by introducing additional free parameters to the *cdf* it would be possible to make the curve fit in the body while retaining the tail properties. The additional parameters could be solved by moment-matching. The moments of the SLN are known in closed form. However, the moments of the new approximating distribution will probably need to be found by numerical integration. However, this integration can probably be performed simply by using a Riemann sum on the transformed integral:

$$E[X^{n}] = \int_{0}^{+\infty} x^{n} f(x) dx = \int_{-\infty}^{+\infty} e^{x(n+1)} f(e^{x}) dx.$$
 (5.1)

Since f(x) is approximately LN in each tail, $f(e^x)$ will be approximately Gaussian in each tail, i.e., it will decay as e^{-x^2} , and the integral can be truncated at moderate values in both limits.

5.2.2 Correlated Shadowing

We have assumed that all wireless paths suffer from independent shadowing. However, the shadowing paths are in fact correlated, with correlation coefficients depending on distance and angular separation [61]. We will briefly show how our moment-matching method can be expanded to include the correlated shadowing case.

We begin with a set of correlated LN RV's $\{Y_i\}_{i=1}^N$ where $\{\ln Y_i\}_{i=1}^N$ form a set of jointly Gaussian RV's. Then we may then write:

$$\begin{bmatrix} \ln Y_1 \\ \ln Y_2 \\ \vdots \\ \ln Y_N \end{bmatrix} = \mathbf{M} \begin{bmatrix} Z_1 \\ Z_2 \\ \vdots \\ Z_N \end{bmatrix}, \tag{5.2}$$

where **M** is a matrix and $\{Z_i\}_{i=1}^N$ are *i.i.d.* standard Gaussian RV's. The first moment of $X = \sum_{i=1}^N Y_i$ is not affected by the correlation and can be found from (2.17). The second moment can be calculated according to (2.18), where the correlation coefficient $\rho_{i,j} \neq 0$.

Now let us examine how we can incorporate this into our analysis of the interference moments. For simplicity, let us consider a system with two interferers, no fast fading $(K = \infty)$, constant activity (p = 1), and no PC $(P_i = 1)$. This can be modeled by a simplified version of (1.1):

$$I = R_1 \zeta_1 + R_2 \zeta_2 \,. \tag{5.3}$$

The second moment can then be calculated as

$$E[I^{2}] = (E[R_{1}^{2}] + E[R_{2}^{2}])E[\zeta_{1}^{2}] + 2E[R_{1}R_{2}\zeta_{1}\zeta_{2}]$$

$$= A_{2}e^{2\sigma^{2}} + 2E[R_{1}R_{2}e^{\ln\zeta_{1}+\ln\zeta_{2}}]$$

$$= A_{2}e^{2\sigma^{2}} + 2E[R_{1}R_{2}e^{(M_{11}Z_{1}+M_{12}Z_{2})+(M_{21}Z_{1}+M_{22}Z_{2})}]$$

$$= A_{2}e^{2\sigma^{2}} + 2E[R_{1}R_{2}e^{(M_{11}+M_{21})Z_{1}+(M_{12}+M_{22})Z_{2}}]$$

$$= A_{2}e^{2\sigma^{2}} + 2E[R_{1}R_{2}e^{(M_{11}+M_{21})Z_{1}+(M_{12}+M_{22})Z_{2}}]$$

$$= A_{2}e^{2\sigma^{2}} + 2E[R_{1}R_{2}e^{(M_{11}+M_{12})(Z_{1}+Z_{2})}]$$

$$= A_{2}e^{2\sigma^{2}} + 2E[R_{1}R_{2}e^{\sqrt{2}(M_{11}+M_{12})Z_{12}}]$$

$$= A_{2}e^{2\sigma^{2}} + 2E[R_{1}R_{2}e^{\sqrt{2}M_{12}Z_{12}}]E[e^{\sqrt{2}M_{11}Z_{12}}].$$
(5.4)

Notice that, in this case, the matrix term M12 = M21 is itself a RV, correlated with R1 and R2, while M11 = M22 is related to the shadowing variance, and Z12 is an independent standard Gaussian RV.

Now in the case of a sum of correlated LN RV's, the behaviour in the upper tail will be the same as in the independent case [21], and so will the first moment. It is in the body and lower tail that there might be a substantial difference.

5.2.3 More Realistic User Distribution

We have assumed in this work that the MU connects to the FS that is physically nearest. This is not always a realistic assumption, as in most cellular systems the MU connects to the FS of which it receives the strongest long-term power. In this case, while the distribution of the location of the MU's remains spatially uniform, the distribution of the location of an MU (conditioned upon being) connected to a particular FS is no longer uniform within that cell's area. Because of this, the distribution functions of R_i and L_i are significantly different [62], and should be modified in our analysis. It is important to note that, in that case, both the pathloss and the shadowing combine together to decide to which FS an MU will connect, and the problem becomes more challenging because the distributions of R_i and L_i are a function of the shadowing variance σ .

5.2.4 Calculating the SIR, SINR and Outage Probability

Our analysis has not included noise. This is because we have considered flat-fading systems, which use narrowband channels (OFDMA, FH-SS). Narrowband systems are generally more affected by interference than wideband systems [2]. Much of our analysis has focused on studying the lower tail of the interference *cdf*. In fact, this lower tail would be affected by adding a noise term. One could say that the analysis of the lower tail of the *cdf* is then of little interest. However, the amount of noise added will vary from system to system, and it is not evident how much of the lower tail is really needed. This is why a good knowledge of the lower tail is still valuable. The calculation of the Signal-to-Interference Ratio (SIR) or Signal-to-Interference-and-Noise Ratio (SINR) is a problem in itself:

- In the case of DL, no PC, the signal pathloss affecting signal strength will be correlated
 with the interference because the position of the MU affects both signal and interference
 powers. New geometrical coefficients need to be devised.
- 2. We can model the signal strength as LN, due to the shadowing that the signal suffers, and the other effects (fading, pathloss) can be incorporated into this LN RV similarly to what was done in Section 4.1.
- If the interference is modeled as a LN as in Section 4.4, the ratio of two LN RV's is simply a LN RV with log-variances added and log-means subtracted. The Bernoulli RV does not affect this calculation.
- 4. If the interference is modeled as a PLN as in section 4.5, the ratio of a LN and a PLN RV is not known. The integral involved in calculating the *cdf* of that ratio is:

$$\int_{-\infty}^{+\infty} e^{By - \frac{1}{2}Ay^2} \Phi^{t-1}(y) dy, \tag{5.5}$$

where *A* and *B* are two constants. It is almost identical to the integral (3.50), which gives the moments of the PLN distribution. This makes the integral even more important to solve.

The outage probability in a system is simply found by calculating the probability that the SIR or SINR is above a certain threshold, predetermined by the modulation and coding scheme. The SIR/SINR is studied in [1], [3], [7]–[9], [12], [16], [59], [63].

5.2.5 Application to Fixed Relay Networks

We consider a two-hop fixed relaying scenario [2]. We define an augmented cell as the combination of a central sub-cell containing a BS and its surrounding six sub-cells containing

RS's. The six RS's each have a good orthogonal wireless link with their BS. There is no communication between augmented cells, and we can assume that the interference from all sub-cells within the augmented cell has been adequately managed [2] and is negligible, while all other sub-cells are considered as interferers (see Figure 5-1).

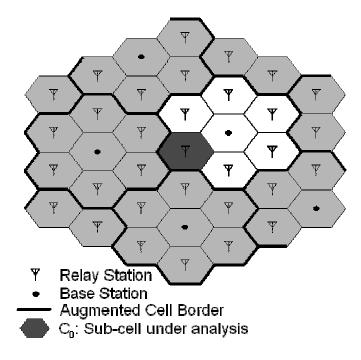


Figure 5-1 Interfering cells (in gray) in a two-hop fixed relay network.

5.2.5.1 Irregular Geometry

As a first approximation to the problem, we can describe the set of interfering sub-cells as having the layout described by Figure 5-1, with geometry parameters given in Table 5-1.

β	A_1	A_2	$A_{ m DL}$	\widetilde{A}_1	$ ilde{ ilde{A}}_2$
3	1.310	0.5184	2.044	0.4506	0.1611
3.2	1.160	0.5027	1.678	0.3874	0.1570
3.4	1.039,	0.4933	1,417	0.3384	0.1550
3.5	0.9868	+ 0.4905	1.316	0.3180	0.1546
3.6	0.9400	0.4888	1.230	0.3000	0.1546
3.8	0.8596	+ 0.4886	1.093	0.2697	0.1557
4	0.7935	0,4921,	0.9932	0.2455 +	0.1580
4.2	0.7389	+ 0.4989	+0.9208	,0.2260	0.1614
4.4	0.6937	0.5088+	0.8 688 ₊	0.2103 ₊	0.1658
4.5	0.6739	+ 0.5148	+0.8490—	0.2036/	0.1685
4.6	0.6559	0.5216⊦	0.8 327+	0.1975	0.1714
4.8	0.6244	0.5374	+0.8091	0.1870	0.1779
5	0.598	0.5561	0.7957	0.1784	0.1855

Table 5-1 Geometry parameters for three tiers of sub-cells in a two-hop fixed relay network.

5.2.5.2 Lower Probability of Activity per Sub-cell

It can be argued that the probability of activity within any sub-cell will be substantially lower than 1. Indeed, [2] defines loading for such a network to be 1 when there is one user per channel per augmented cell. This corresponds to $p \approx 0.14$. A more aggressive reuse scheme can be used, where each channel is used once per sub-cell, which corresponds to a loading of 7 and p = 1.

In the approach proposed in this section, we have assumed that there is no interference coming from within the augmented cell. This corresponds to a scenario where the loading is 1 or less, and

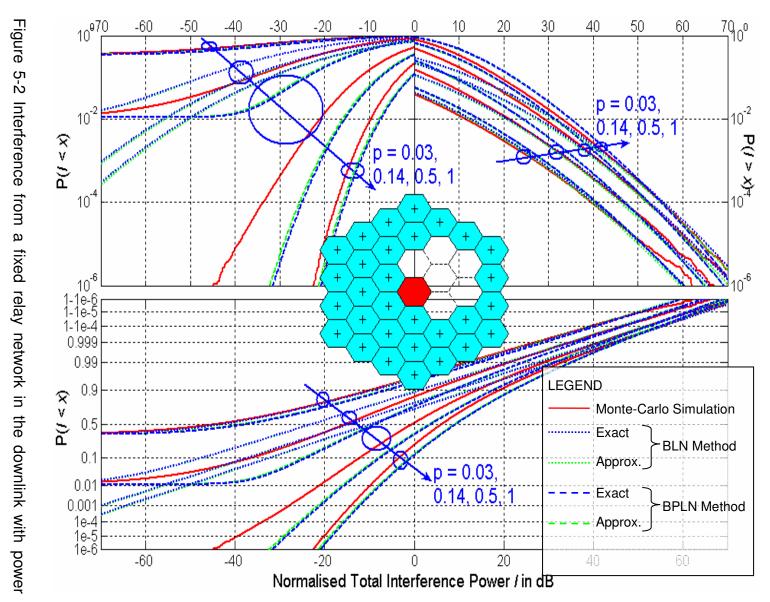
is only approximately true for higher loading. Thus we will be primarily interested in values of p in the range [0, 0.14], where our methods perform well, as will be seen in the following section.

5.2.5.3 Some Simulation Results

We simulate the interference in the fixed relay network in the same way as we did in Section 4.6, except that the layout is different. The simulation parameters are summarized in Table 5-2, and the results are shown in Figure 5-2.

Figure	5-2						
β	4						
σ [dB]			12				
K			0				
p	0.035	0.14	0.5	1			
Layout		3-Tier Two	o-Hop Relay				
Link		1	DL				
PC		Yes (o	$\tau_e = 1 dB$)				
p_{χ}	0.6566	0.9892	$1-9.3\cdot10^{-10}$	1			
$a_I =$	-6.601	-5.136	-3.243	-2.203			
$a_I \approx$	-6.601	-5.136	-3.243	-2.203			
$b_I =$	4.394	4.281	4.131	4.046			
$b_I \approx$	4.353	4.239	4.088	4.002			
$m_I =$	-5.976	-6.38098	-6.38056	-6.369			
$m_I \approx$	-5.921	-6.32710	-6.32711	-6.316			
s_I	4.225						
$t_I =$	1.108	4.434	15.83	31.67			
$t_I \approx$	1.05	4.2	15	30			

Table 5-2 Simulation and analytical parameters for a two-hop fixed relay network.



1

5.2.6 Evolution of Interference in Time

Throughout this work, we have investigated what is the distribution of the total interference power at any given instant. However, in order to get useful knowledge about such performance statistics as word or packet error rates, quality of service, etc., it is important to understand how the total interference power evolves in time, which would require studying the interference as a random process. This has already been done in various ways in literature [11]–[13]. We hope to expand our method to include the evolution of interference in time, so as to obtain very precise long-term performance statistics.

5.2.7 CDMA Systems

It would be interesting to apply our method to broadband CDMA networks; however there are several necessary considerations:

- 1. The channel is no longer flat-fading.
- 2. The number of users in each cell is no longer one, but can be quite large. This has at least two important consequences:
- 3. The number of summands in the SLN method needs to be large. Large values of N have not typically been analysed. We are only aware of the LL method proving robust for as much as 1000 summands (Figure 6 in [27]). Incidentally, our PLN method does not perform very well for N = 1000 (we do not show the simulations here).
- 4. The correlation between shadowing paths needs to be studied. In the UL, many users will be located within the same narrow angular sector, which implies high shadowing correlation [61]. The independence assumption will no longer be acceptable.

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Papers that have resulted from our thesis work:

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These last two papers can be obtained from our website: www.sce.carleton.ca/~sz.